

Dynamical models: from neural microcircuits to cortical regions

Steve Coombes



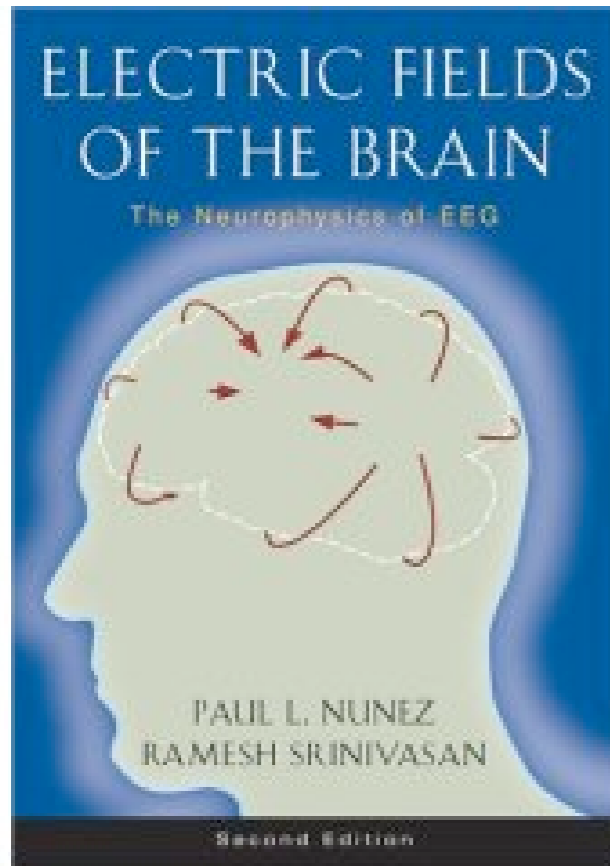
The University of
Nottingham

S Coombes 2005 Waves, bumps, and patterns in neural field theories, *Biological Cybernetics*, Vol 93, 91-108.

<http://eprints.nottingham.ac.uk/153/>

http://www.scholarpedia.org/article/Neural_fields

How do we get to a brain-wave equation?

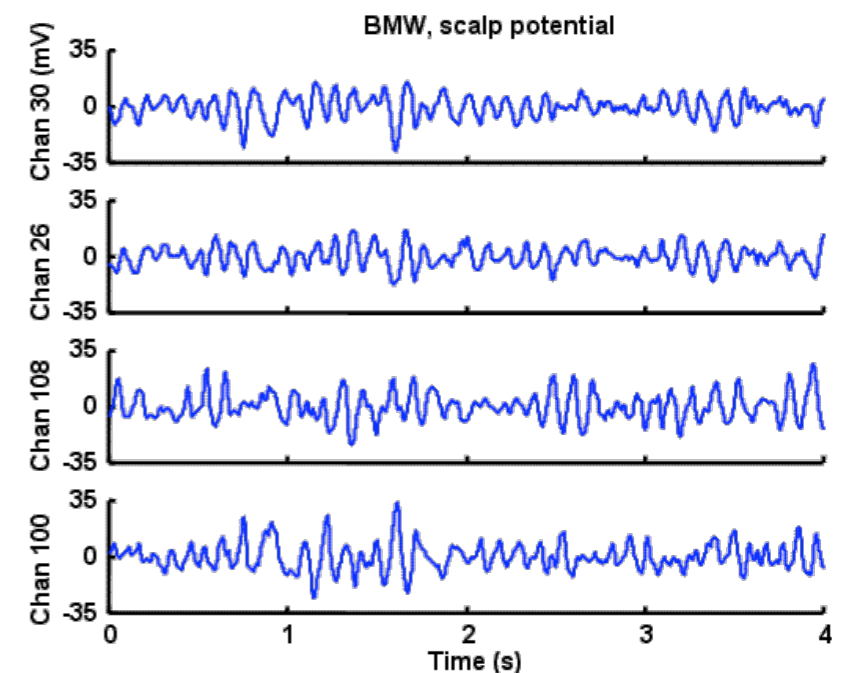


Electric Fields of the Brain: The Neurophysics of EEG

Paul L. Nunez and Ramesh Srinivasan, OUP 2006

$$\left(1 + \frac{1}{\alpha} \frac{\partial}{\partial t}\right) u = \psi$$

EEG Rhythms

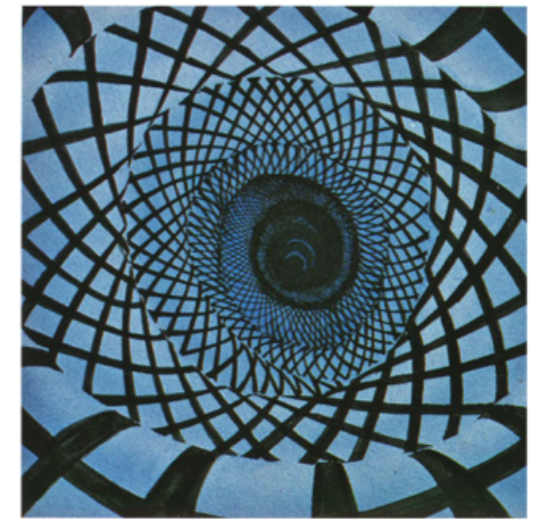
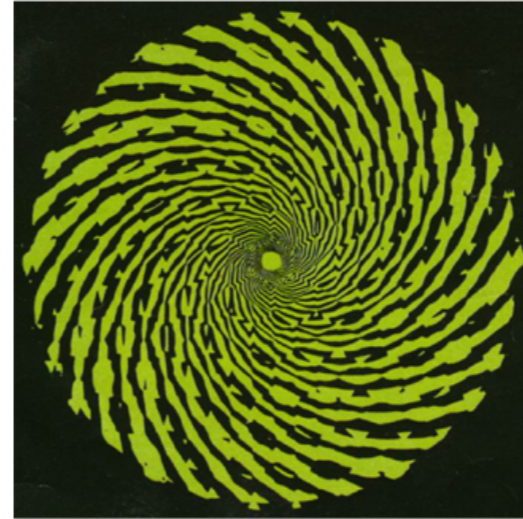


$$[(v + \partial_t)^2 - v^2 \partial_{xx}] \psi(x, t) = [v^2 + v \partial_t] f \circ u(x, t)$$

V.K. Jirsa, H. Haken. A Derivation of a Macroscopic Field Theory of the Brain from the Quasi microscopic Neural Dynamics, Physica D 99, 503-526 (1997)

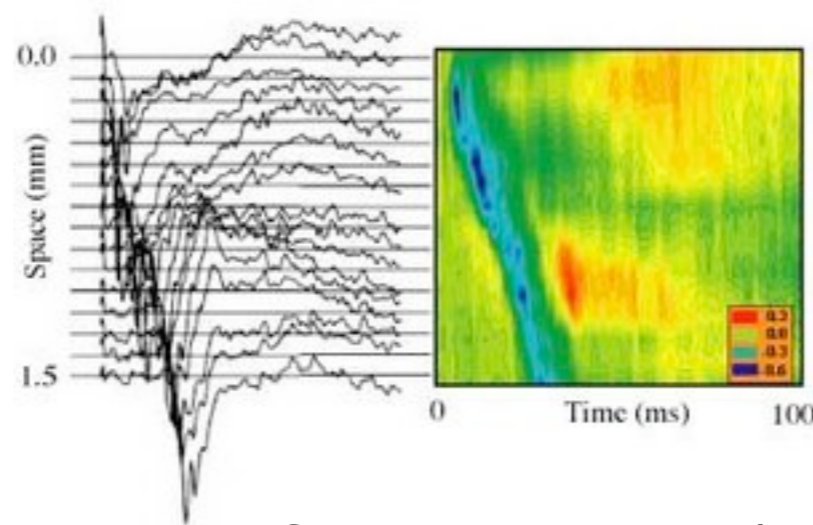
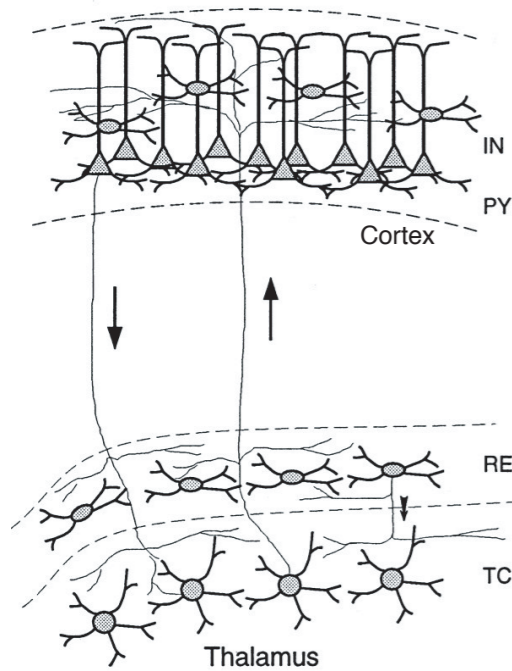
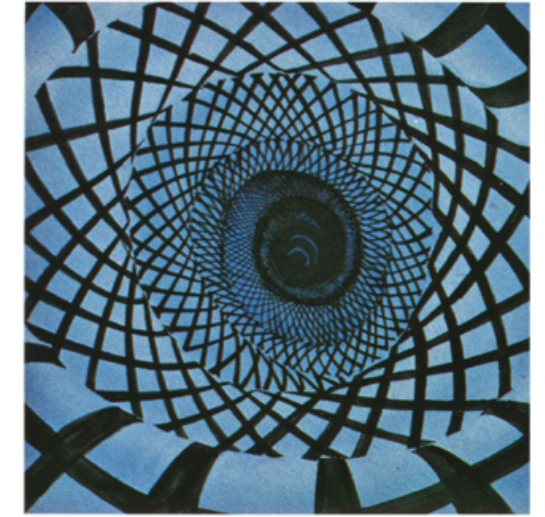
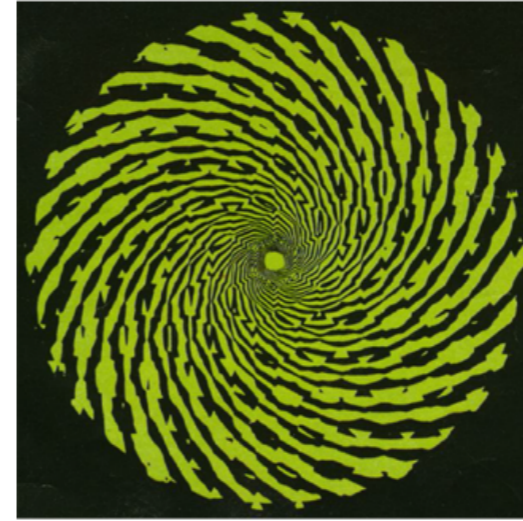
What else is it good for?

Visual hallucinations
(Ermentrout & Cowan)

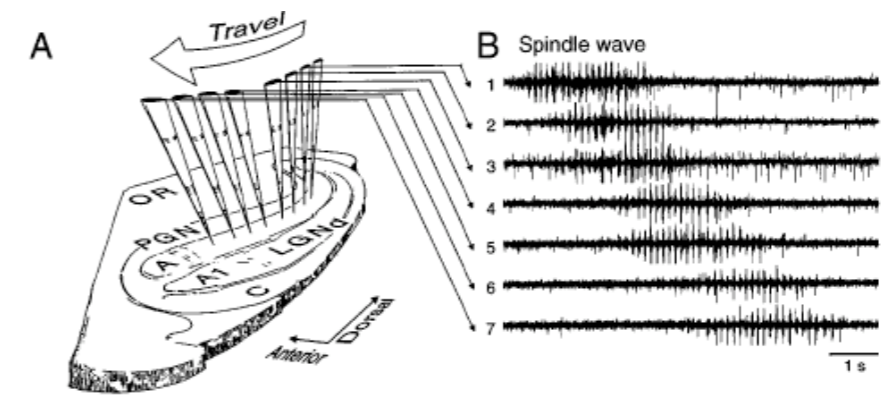


What else is it good for?

Visual hallucinations (Ermentrout & Cowan)



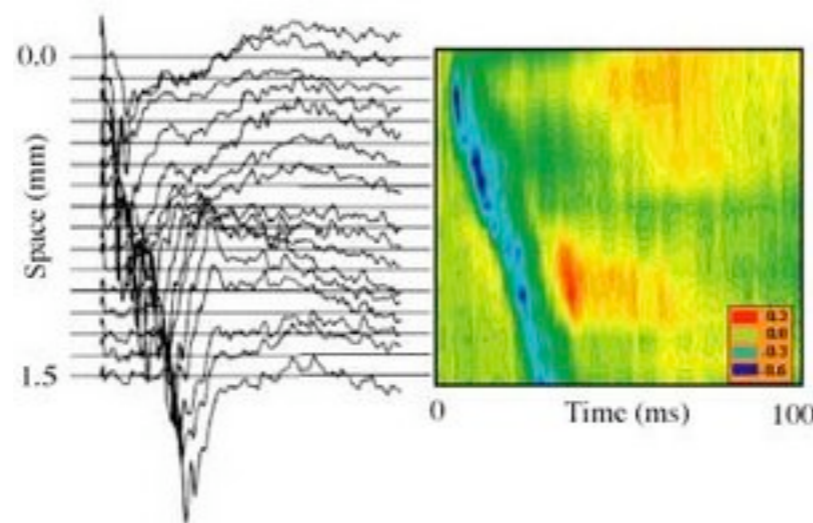
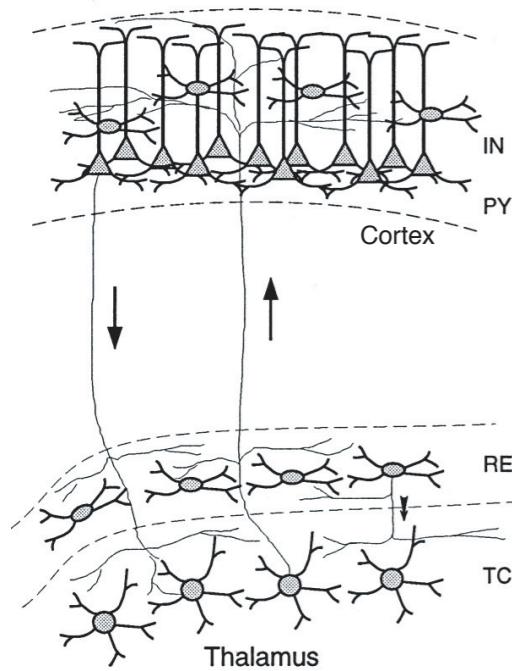
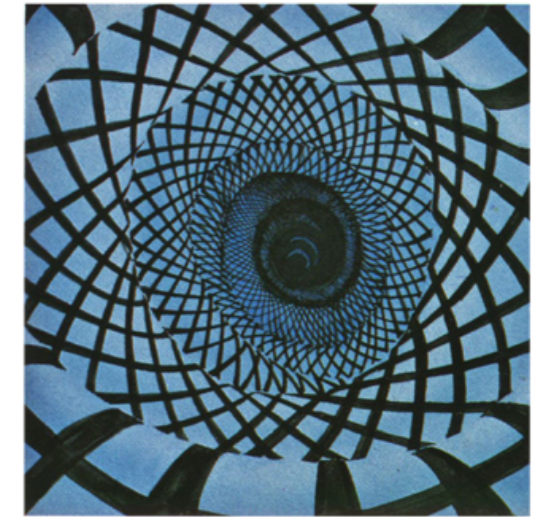
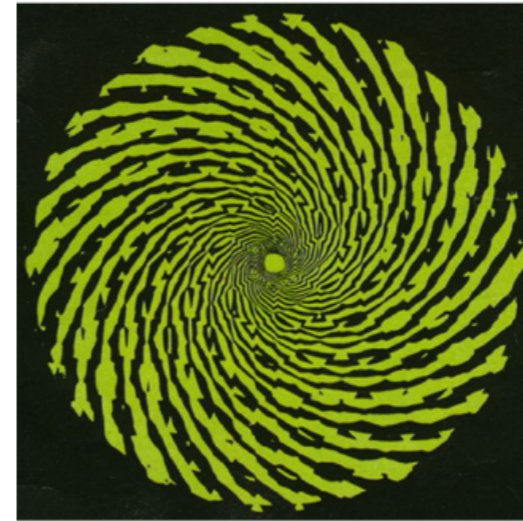
Cortical waves cm/s



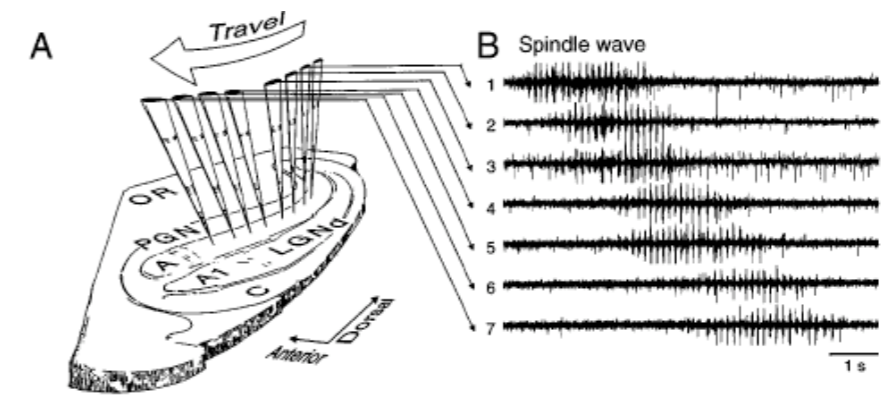
Spindle waves in thalamic slices mm/s

What else is it good for?

Visual hallucinations (Ermentrout & Cowan)

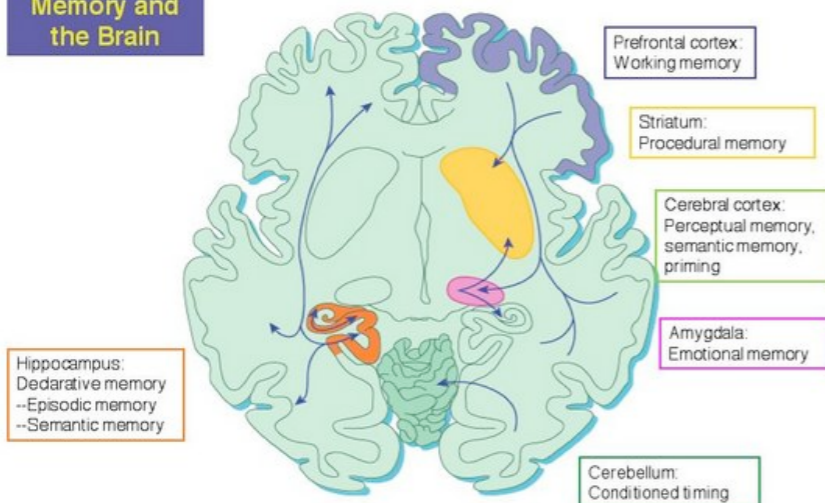


Cortical waves cm/s

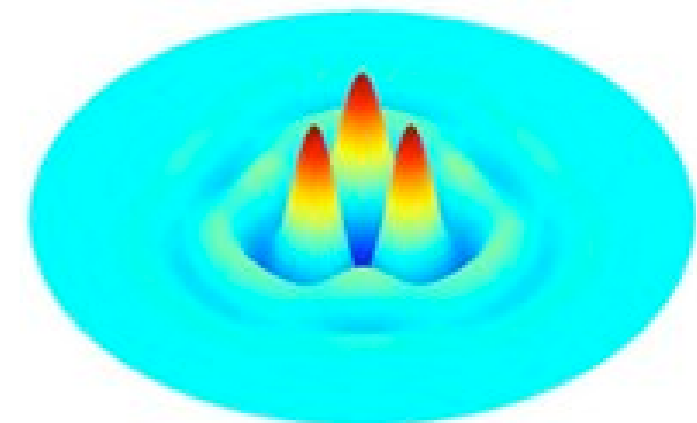


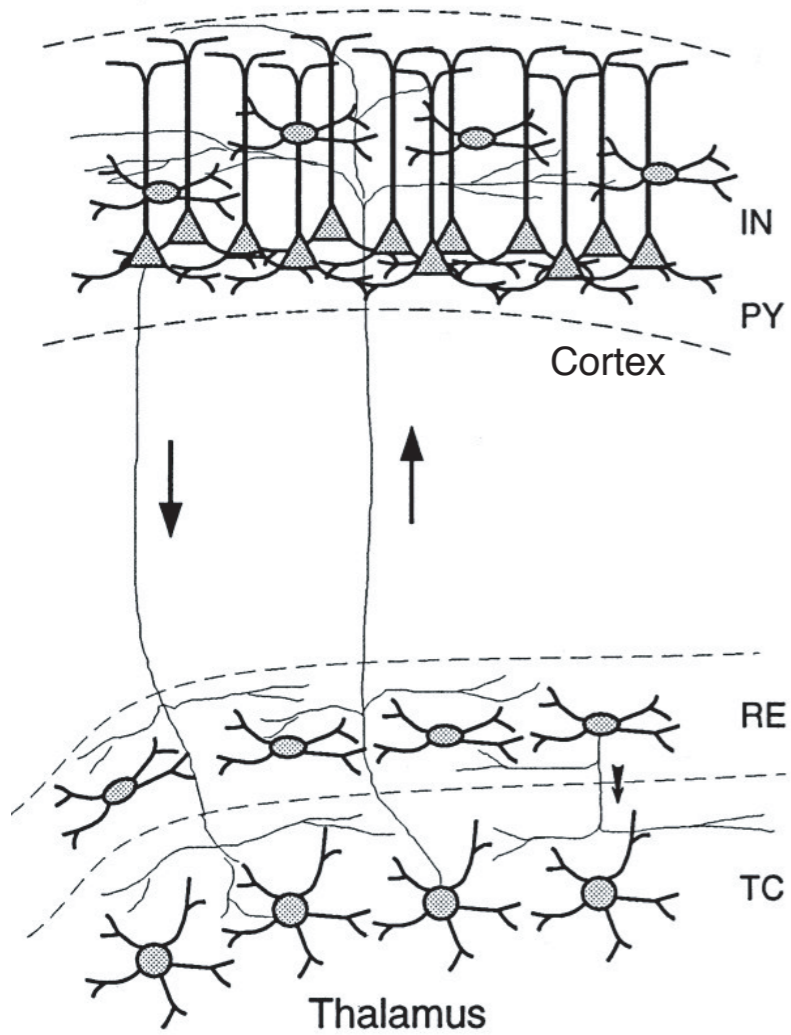
Spindle waves in thalamic slices mm/s

Memory and the Brain



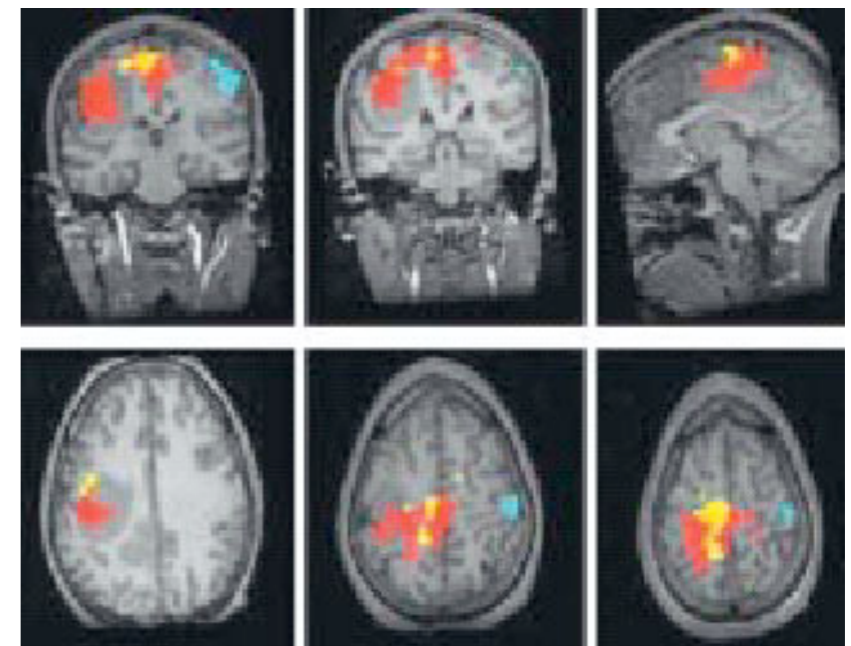
mechanisms for short term memory

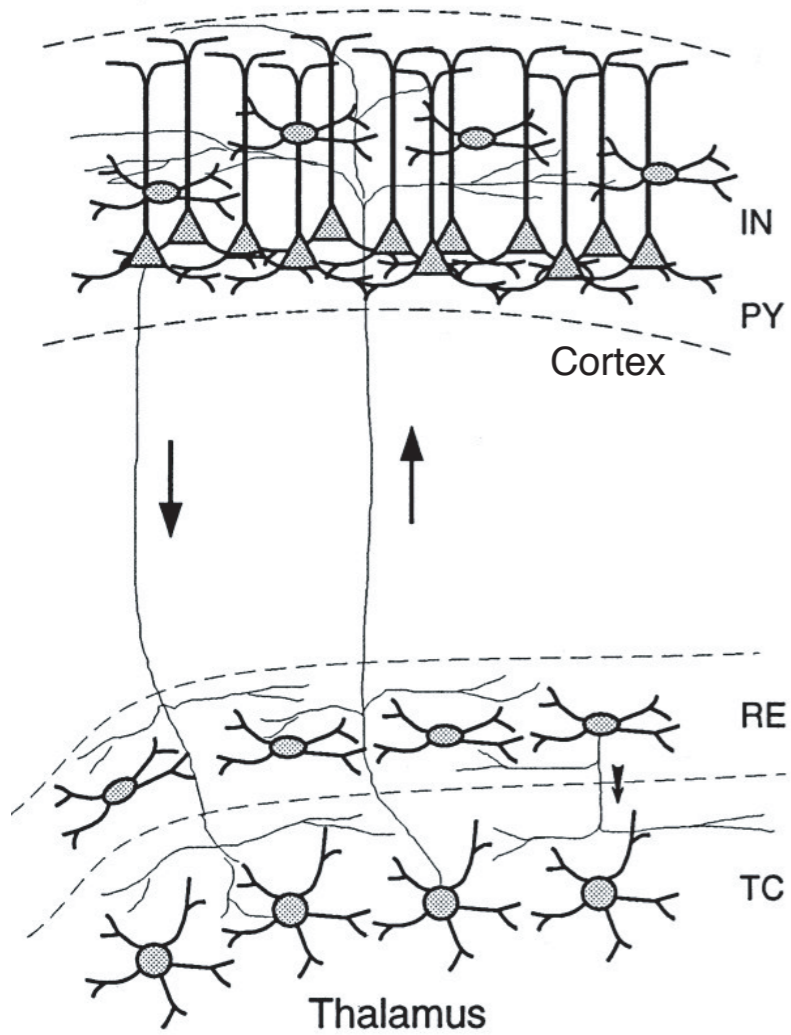




need a spiking single neuron model

HH, FHN, ML, WB, PR, ..., IF

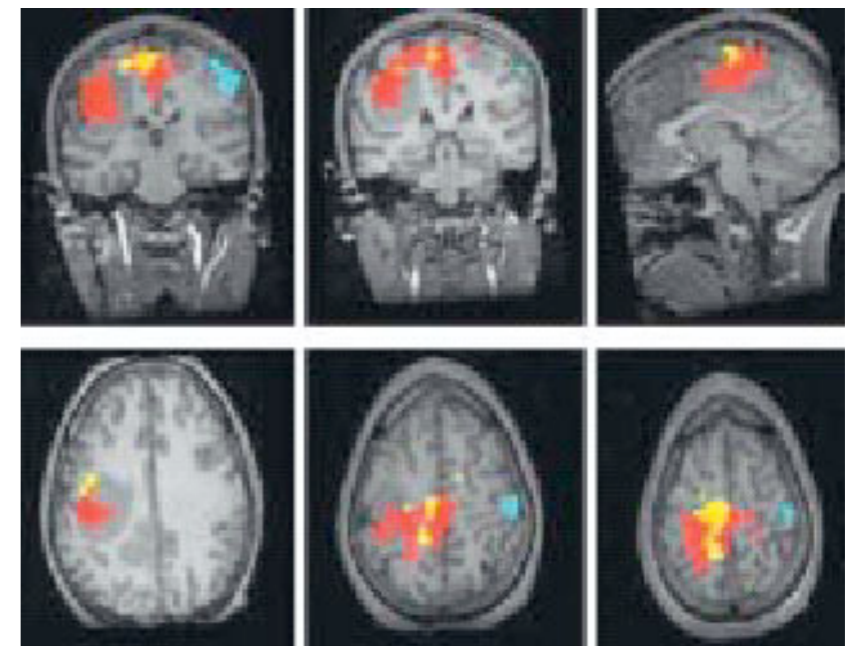


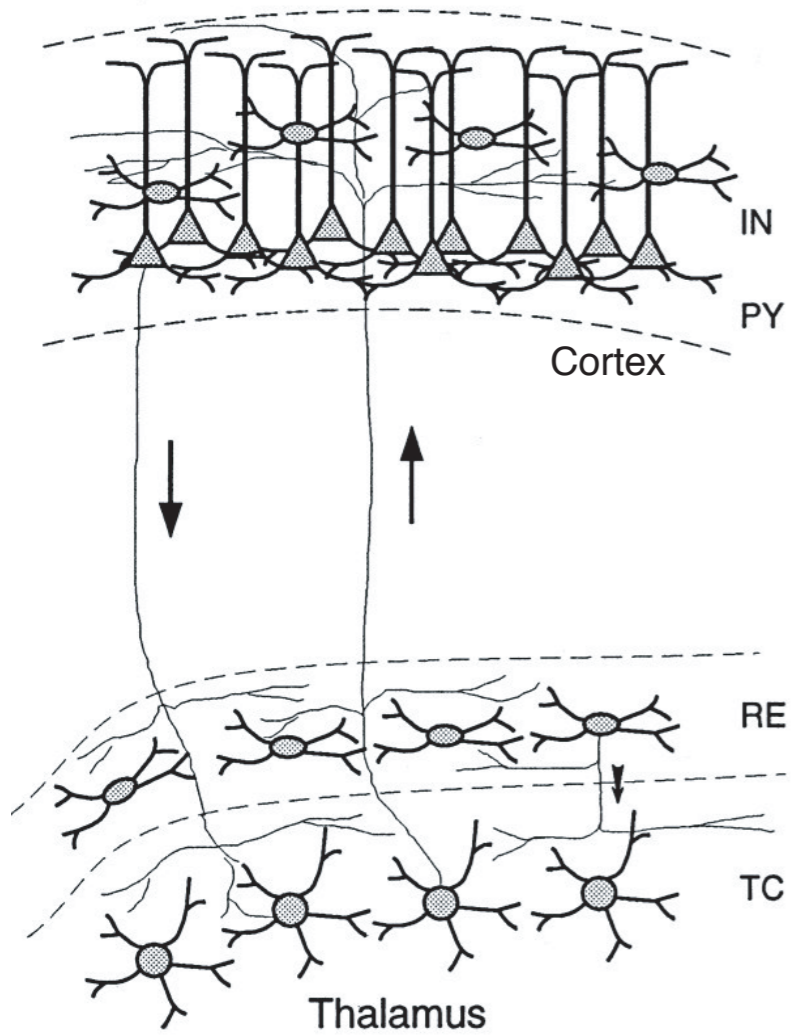


need a spiking single neuron model

HH, FHN, ML, WB, PR, ..., IF

$$C \frac{dv}{dt} = - \sum_k g_k m_k^{p_k} h_k^{q_k} (v - v_k) + u$$

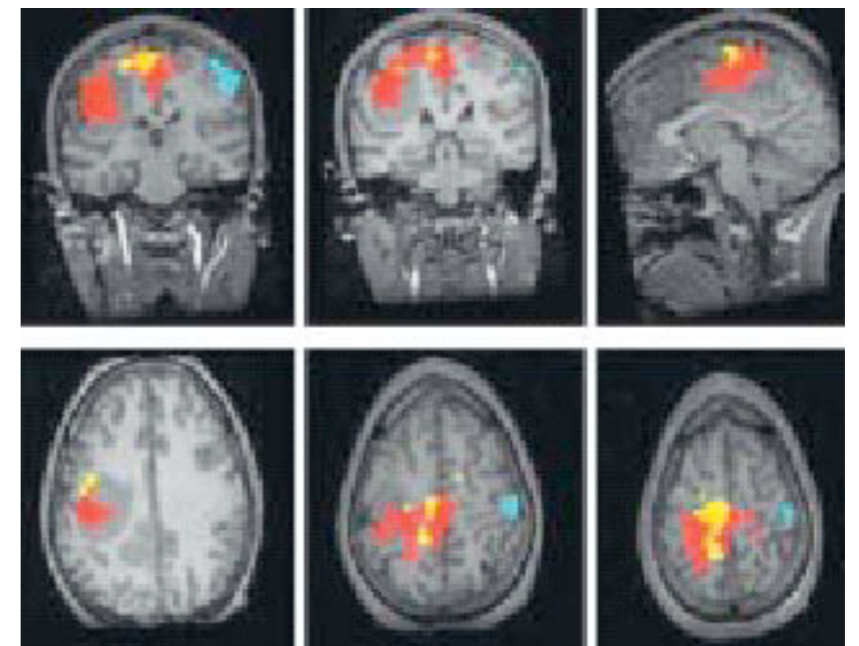
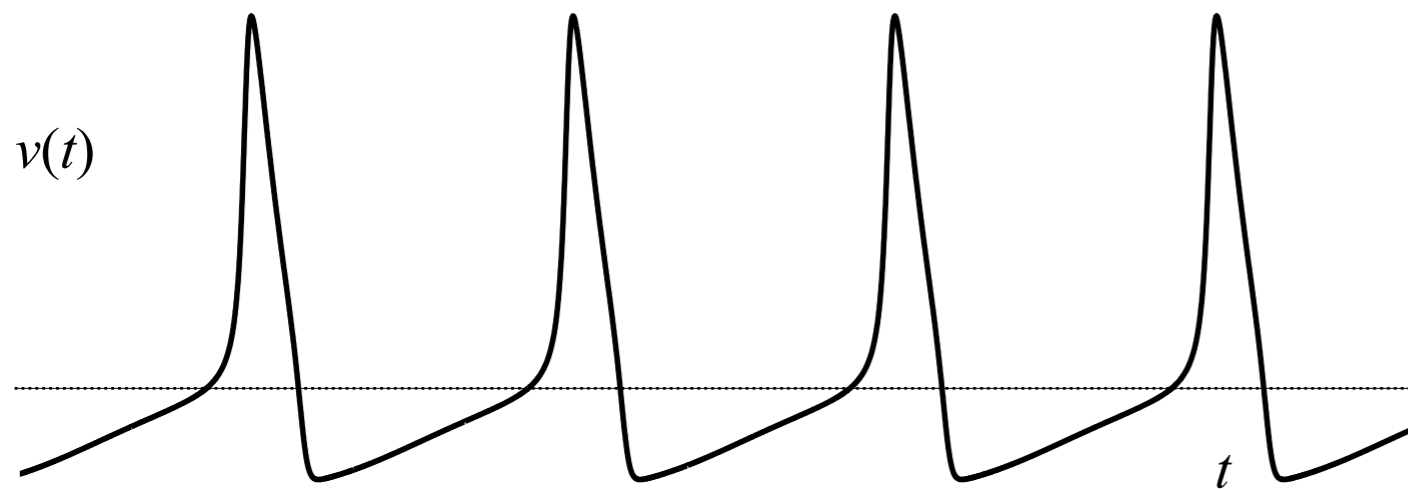




need a spiking single neuron model

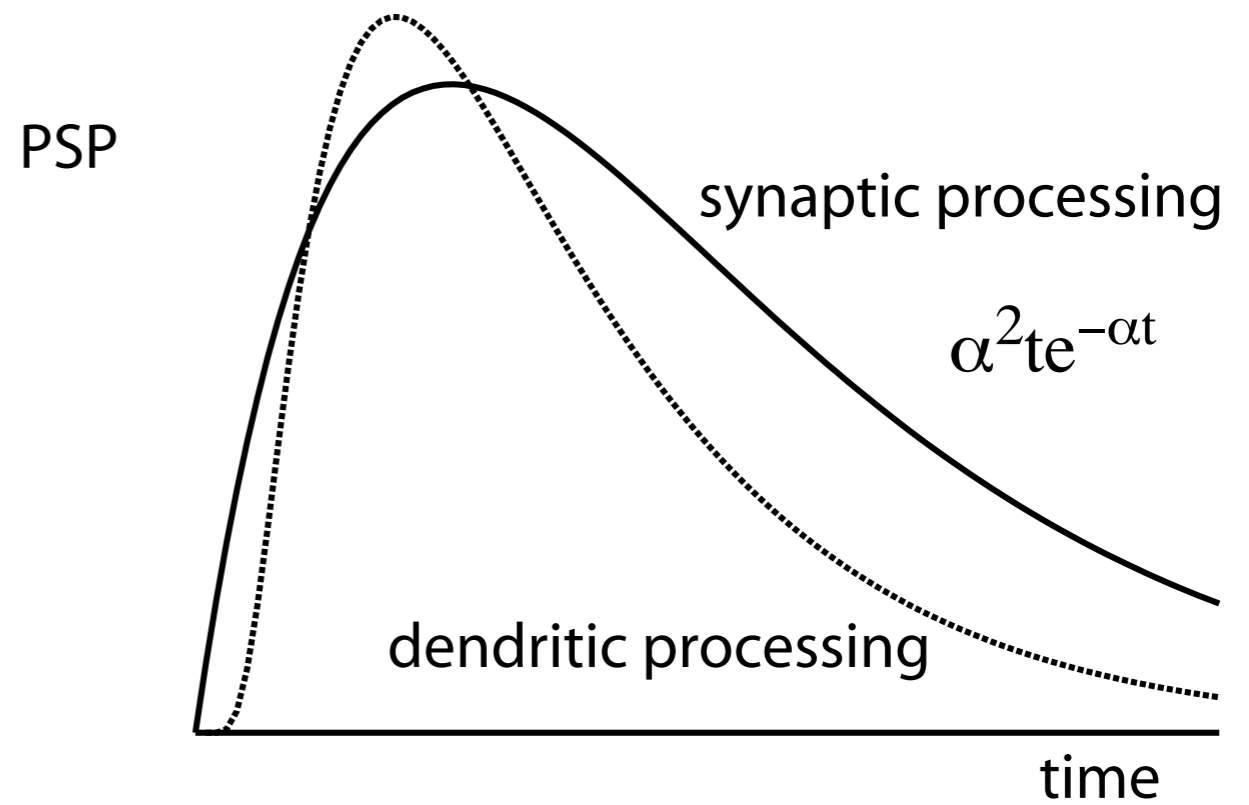
HH, FHN, ML, WB, PR, ..., IF

$$C \frac{dv}{dt} = - \sum_k g_k m_k^{p_k} h_k^{q_k} (v - v_k) + u$$



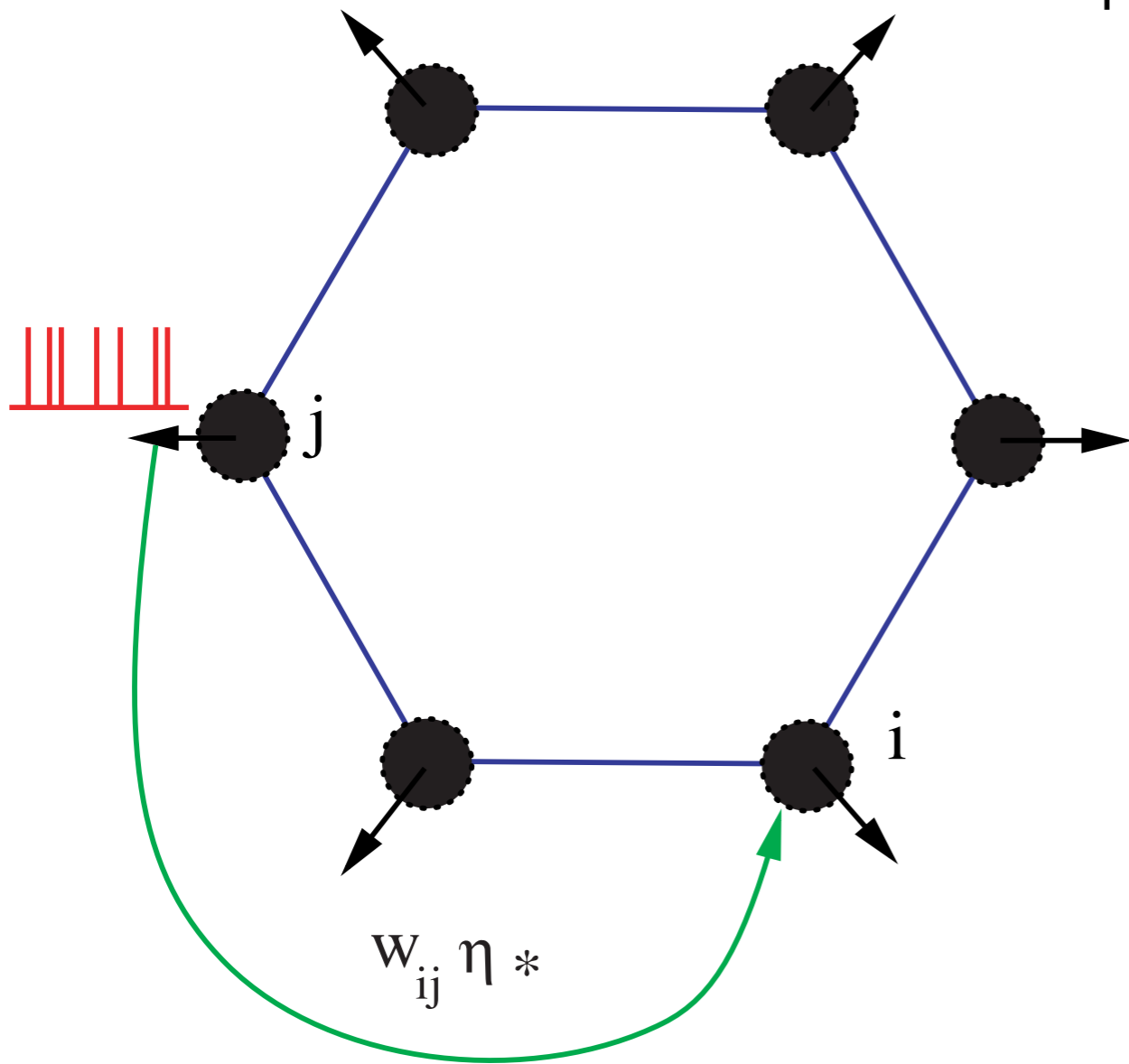
$$u(t) = \sum_m \eta(t - T^m)$$

also need a synapse model

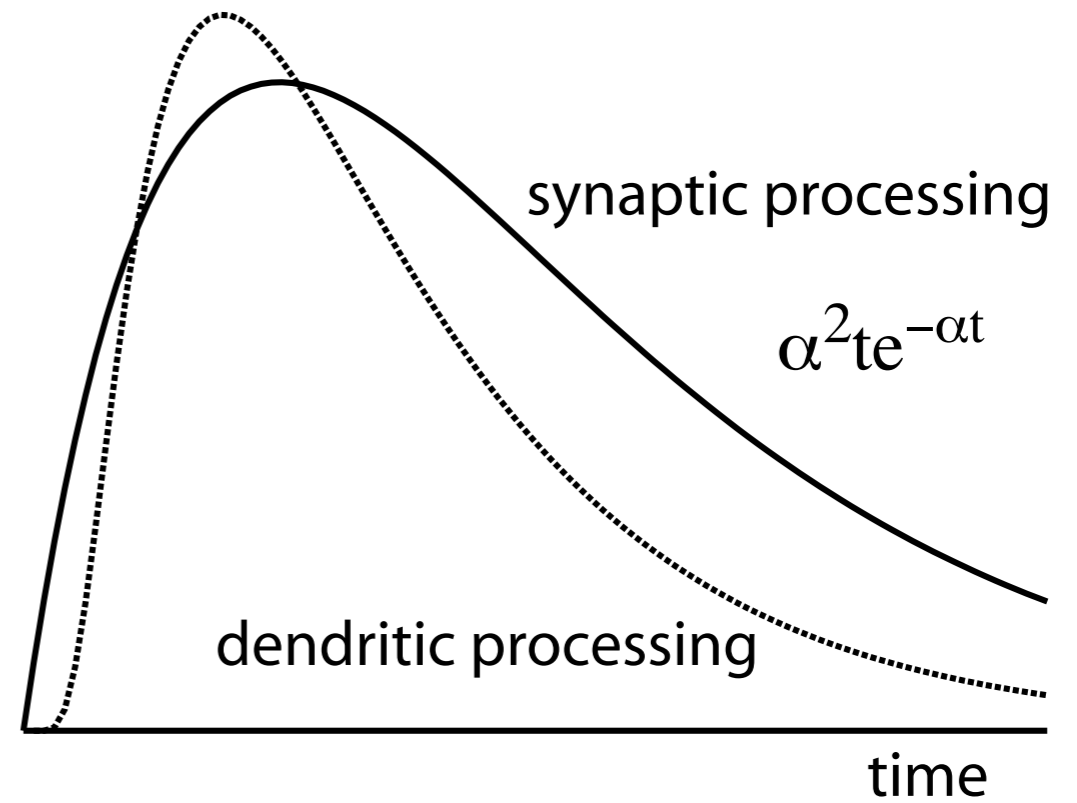


$$u(t) = \sum_m \eta(t - T^m)$$

also need a synapse model

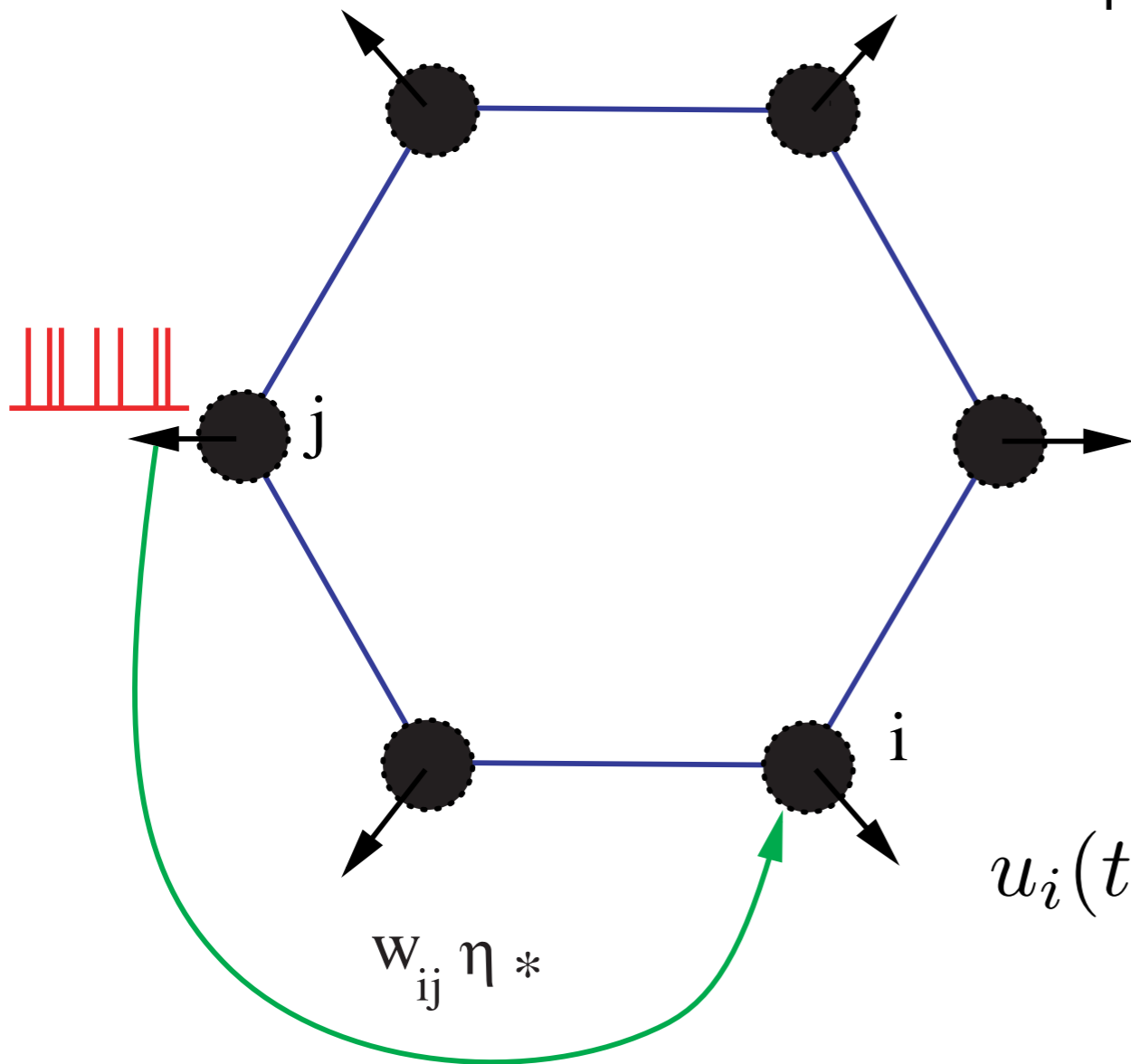


PSP

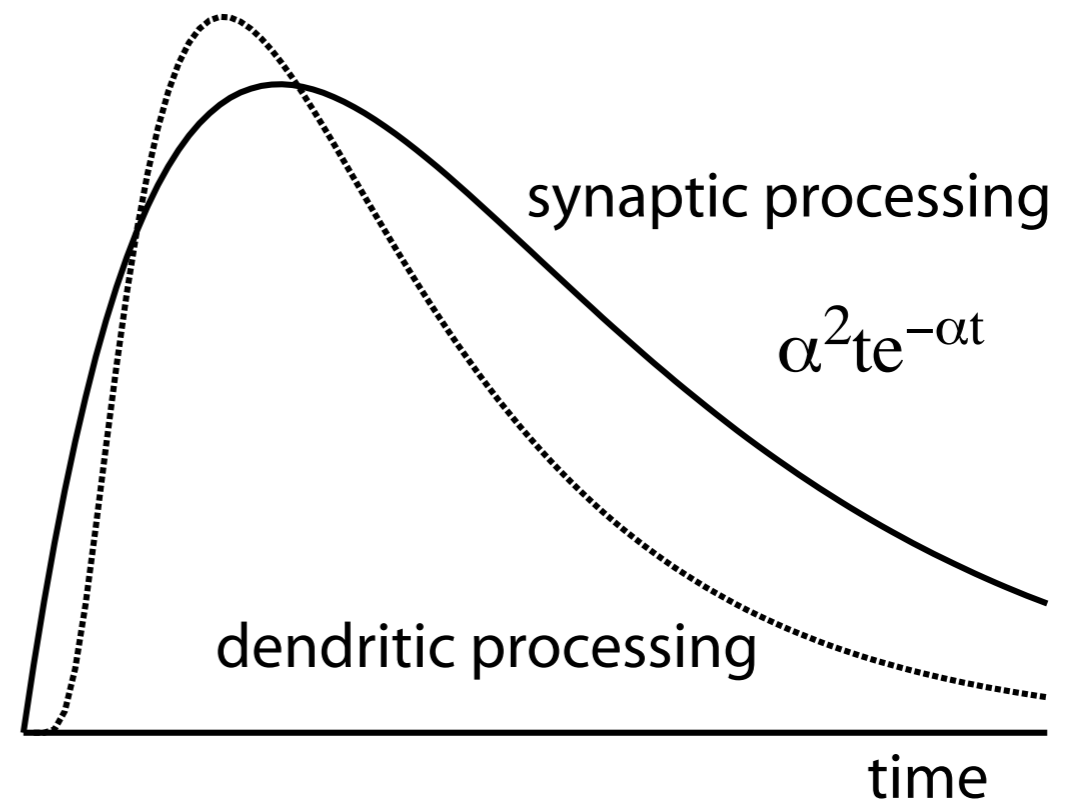


$$u(t) = \sum_m \eta(t - T^m)$$

also need a synapse model



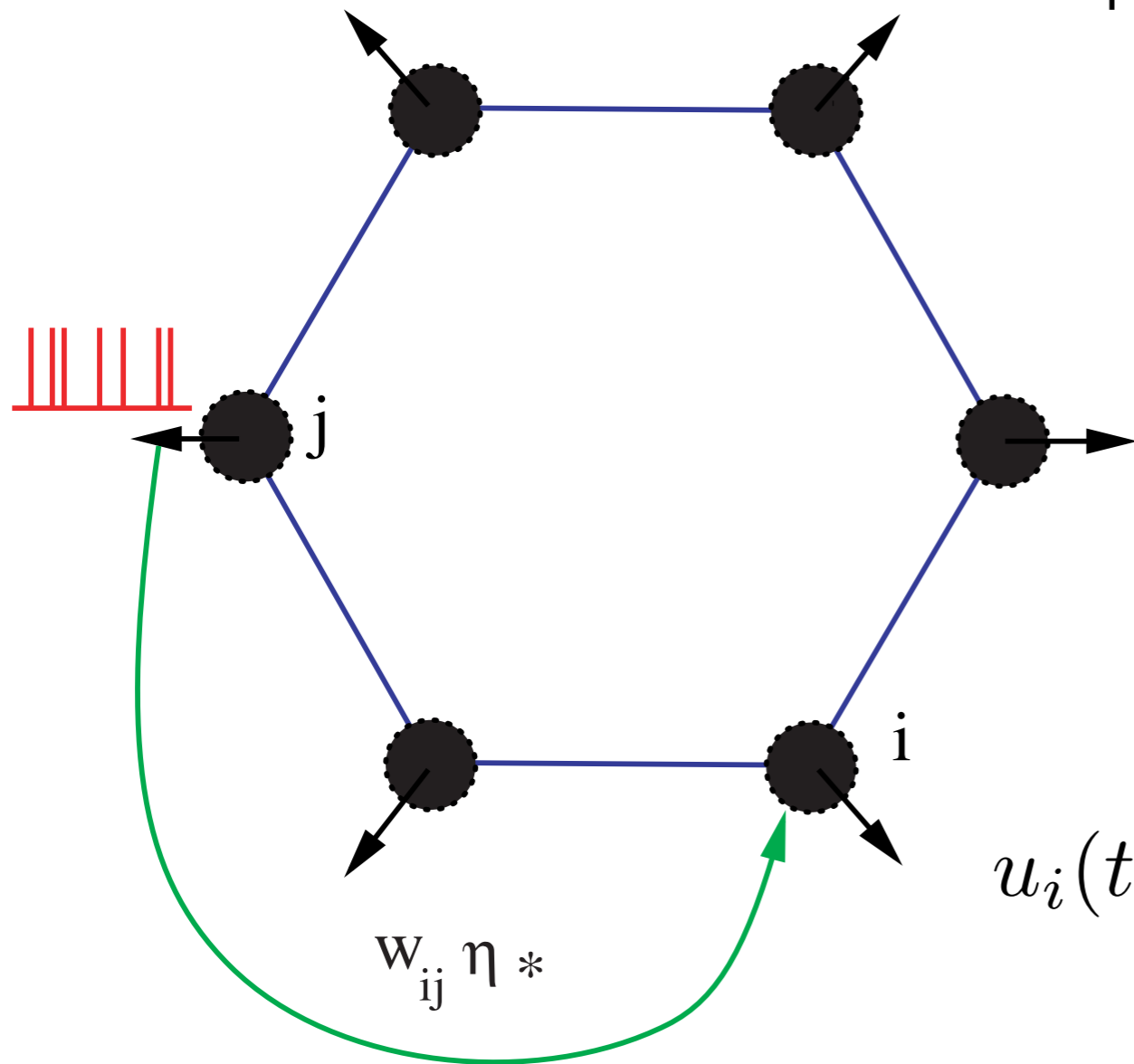
PSP



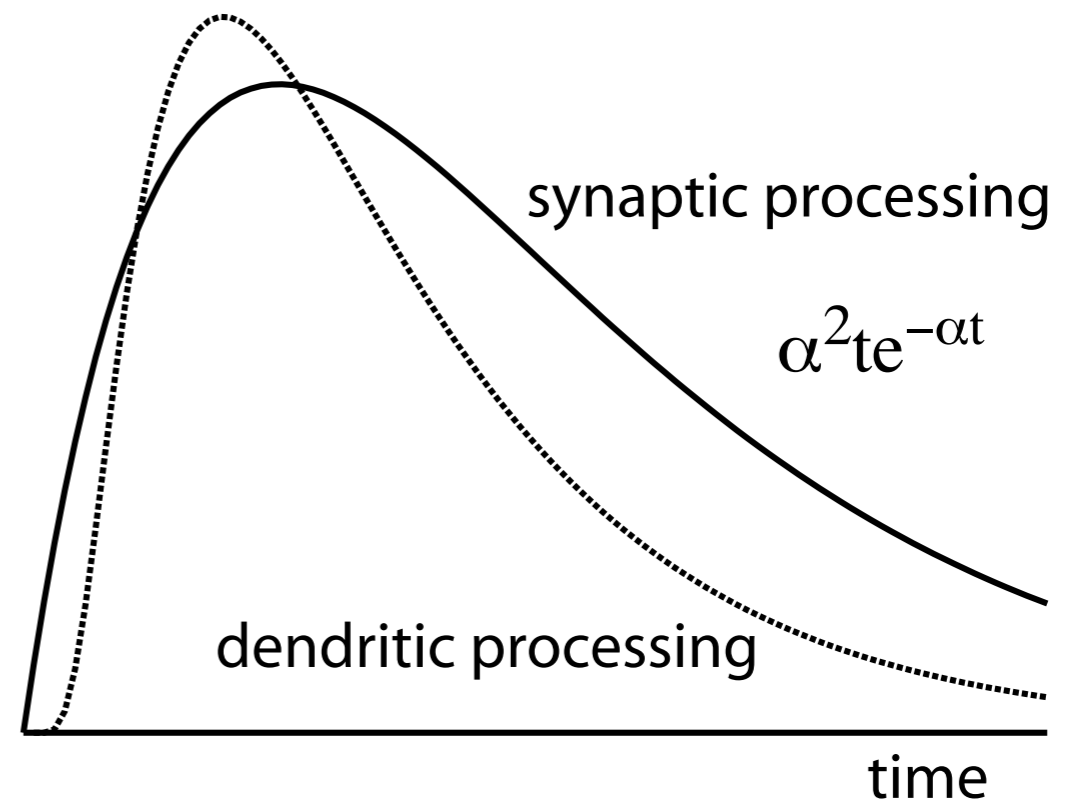
$$u_i(t) = \sum_j w_{ij} \sum_m \eta(t - T_j^m)$$

$$u(t) = \sum_m \eta(t - T^m)$$

also need a synapse model



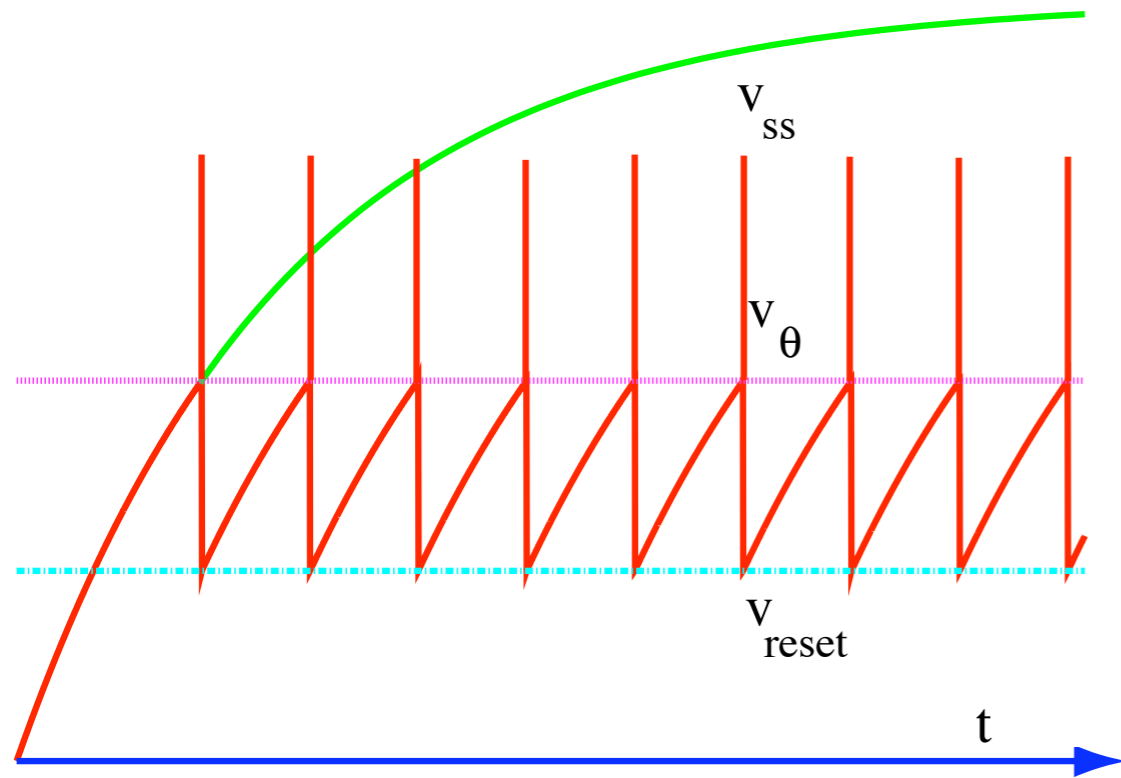
PSP



$$u_i(t) = \sum_j w_{ij} \sum_m \eta(t - T_j^m)$$

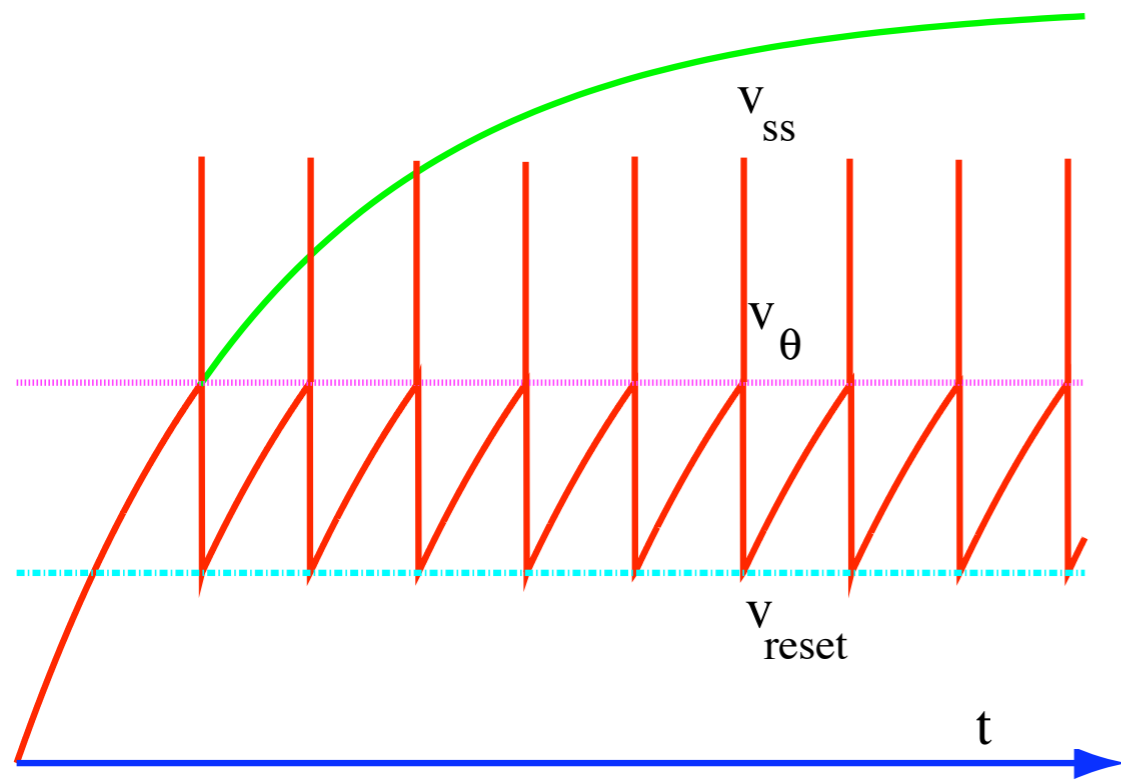
$$= \sum_j w_{ij} \int_0^\infty \eta(s) \sum_m \delta(s - t + T_j^m) ds$$

Slow synapses: spike train \rightarrow firing rate f



Firing rate for single neuron model in response to constant current injection is known

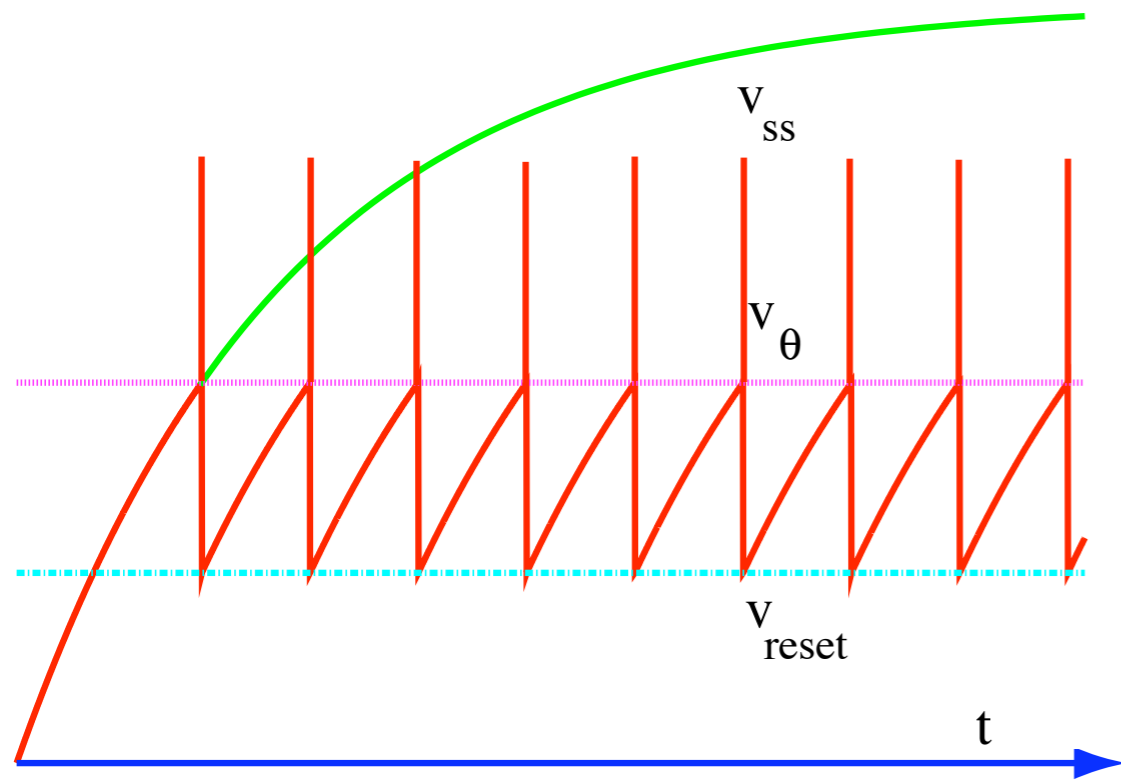
Slow synapses: spike train \rightarrow firing rate f



Firing rate for single neuron model in response to constant current injection is known

$$f = f(v_{\text{ss}}) \quad v_{\text{ss}} = v_{\text{ss}}(u)$$

Slow synapses: spike train \rightarrow firing rate f

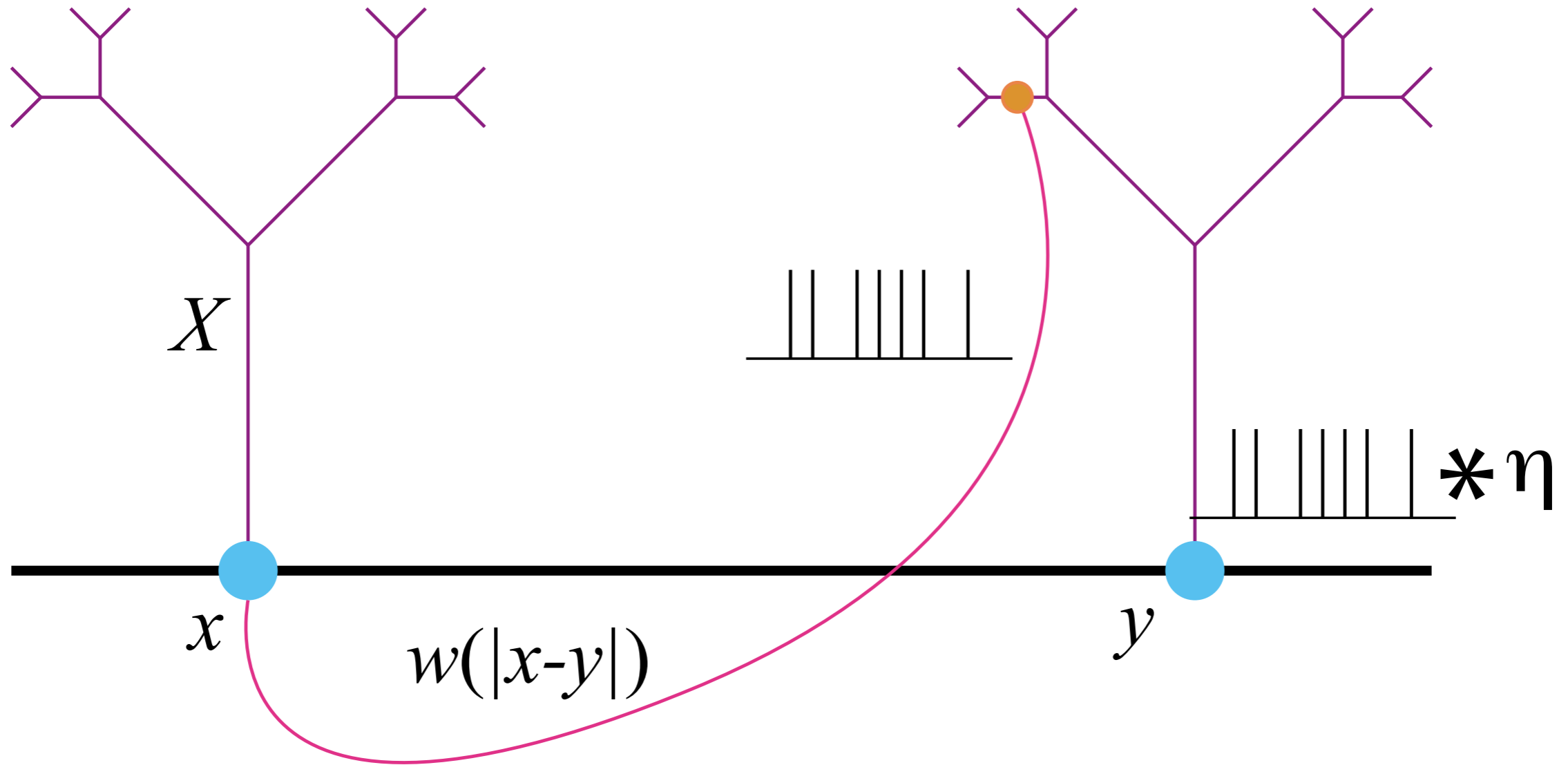


Firing rate for single neuron model in response to constant current injection is known

$$f = f(v_{ss}) \quad v_{ss} = v_{ss}(u)$$

$$u_i(t) = \sum_j w_{ij} \int_0^{\infty} \eta(s) f(v_j(t-s)) ds$$

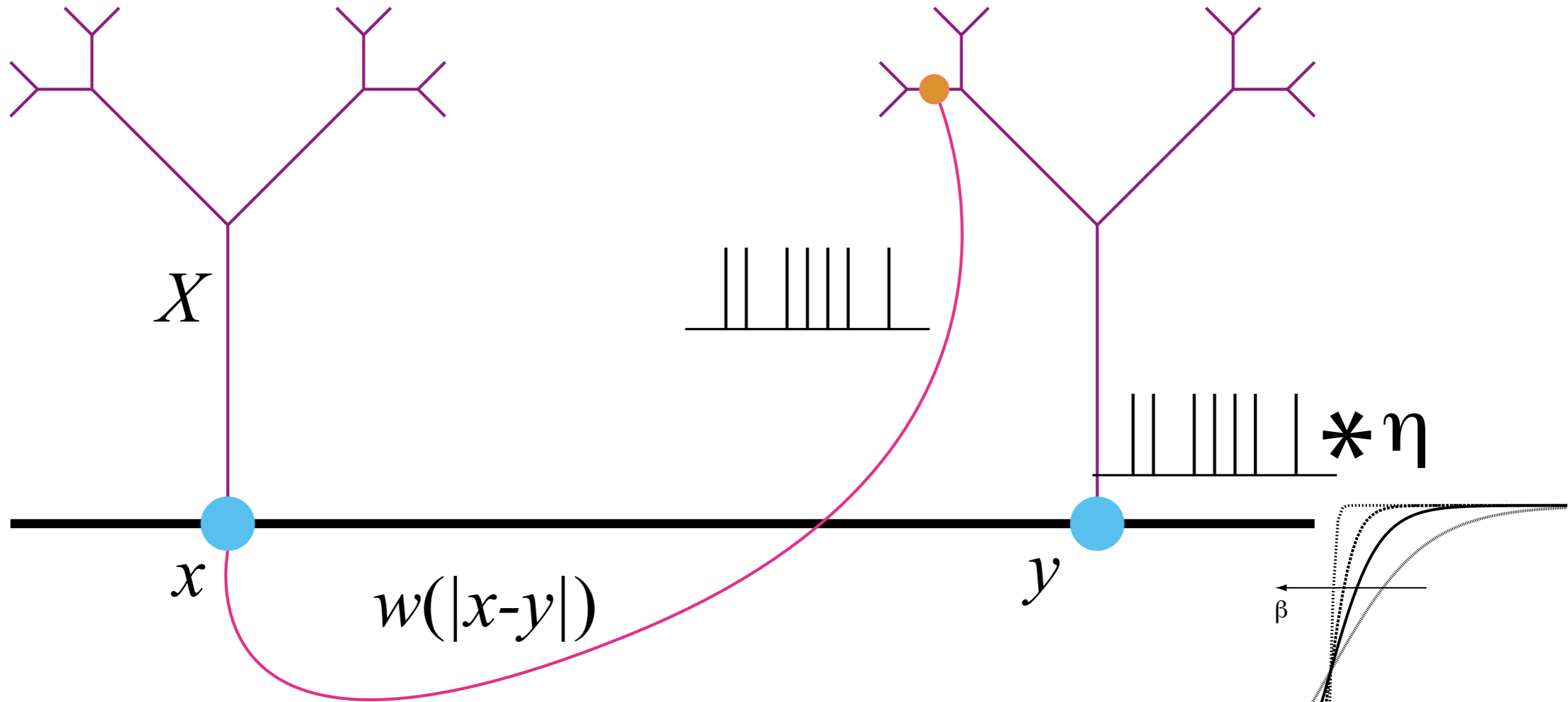
Neural Field Model



Wilson and Cowan (1972, 1973), Amari (1977)

$$u(x, t) = \int_{-\infty}^{\infty} dy w(y) \int_0^{\infty} ds \eta(s) f(u(x - y, t - s - |y|/v))$$

Neural Field Model

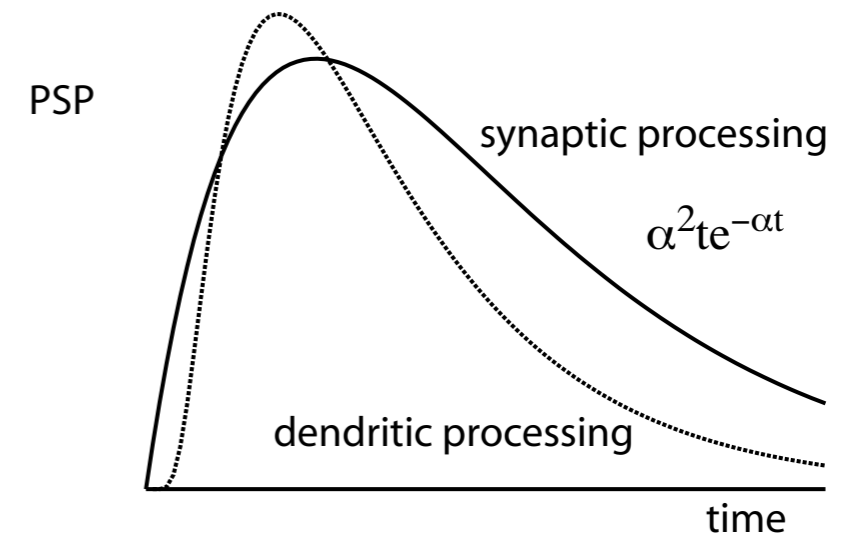


Wilson and Cowan (1972, 1973), Amari (1977)

$$u(x, t) = \int_{-\infty}^{\infty} dy w(y) \int_0^{\infty} ds \eta(s) f(u(x - y, t - s - |y|/v))$$

Exponential synapse

$$\eta(t) = \alpha e^{-\alpha t}, \quad t > 0$$

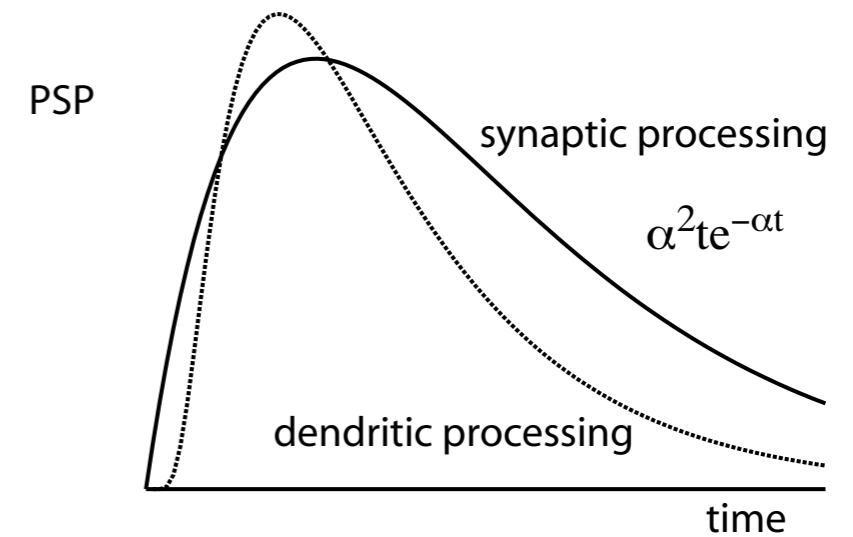


Exponential synapse

$$\eta(t) = \alpha e^{-\alpha t}, \quad t > 0$$

$$\left(1 + \frac{1}{\alpha} \frac{\partial}{\partial t}\right) u = \psi$$

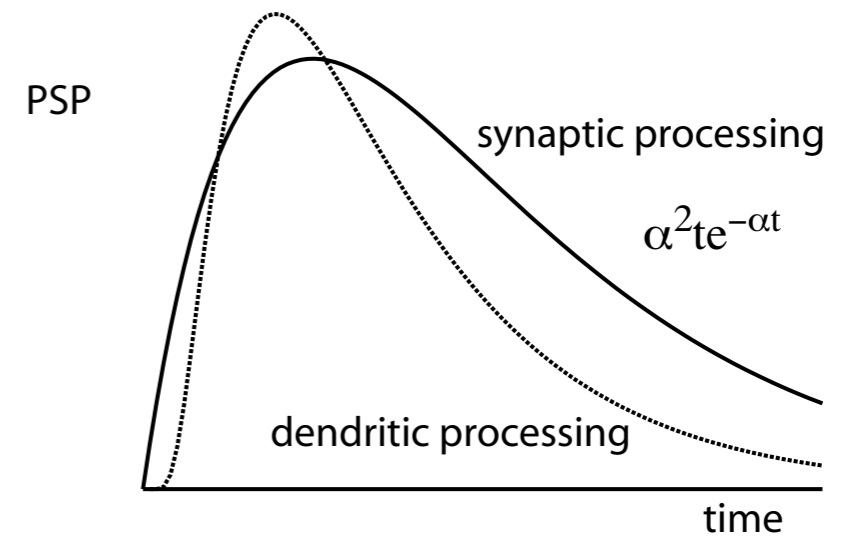
$$\psi(x, t) = \int_{-\infty}^{\infty} w(x - y) f(u(y, t - |x - y|/v)) dy$$



Exponential synapse

$$\eta(t) = \alpha e^{-\alpha t}, \quad t > 0$$

$$\left(1 + \frac{1}{\alpha} \frac{\partial}{\partial t}\right) u = \psi$$



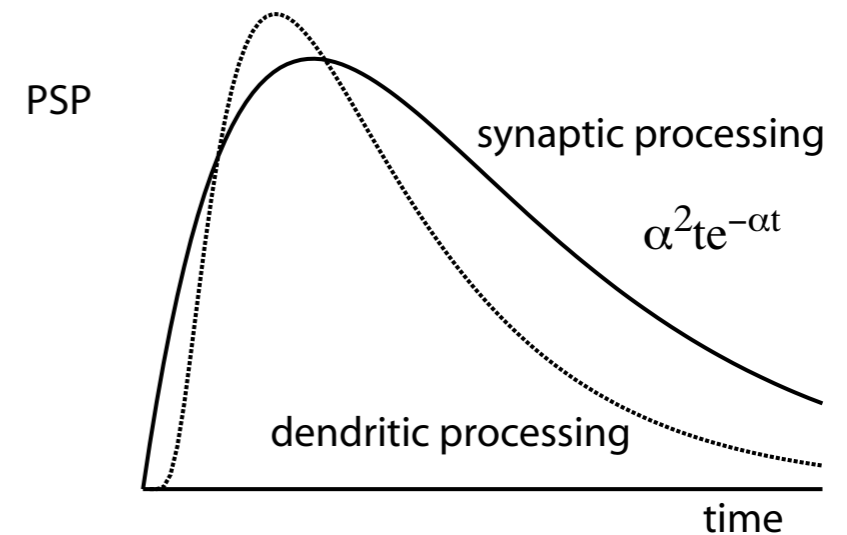
$$\begin{aligned} \psi(x, t) &= \int_{-\infty}^{\infty} w(x - y) f(u(y, t - |x - y|/v)) dy \\ &= \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} ds G(x - y, t - s) \rho(y, s) \end{aligned}$$

$$G(x, t) = w(x) \delta(t - |x|/v) \quad \rho = f \circ u$$

Exponential synapse

$$\eta(t) = \alpha e^{-\alpha t}, \quad t > 0$$

$$\left(1 + \frac{1}{\alpha} \frac{\partial}{\partial t}\right) u = \psi$$



$$\psi(x, t) = \int_{-\infty}^{\infty} w(x - y) f(u(y, t - |x - y|/v)) dy$$

$$= \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} ds G(x - y, t - s) \rho(y, s)$$

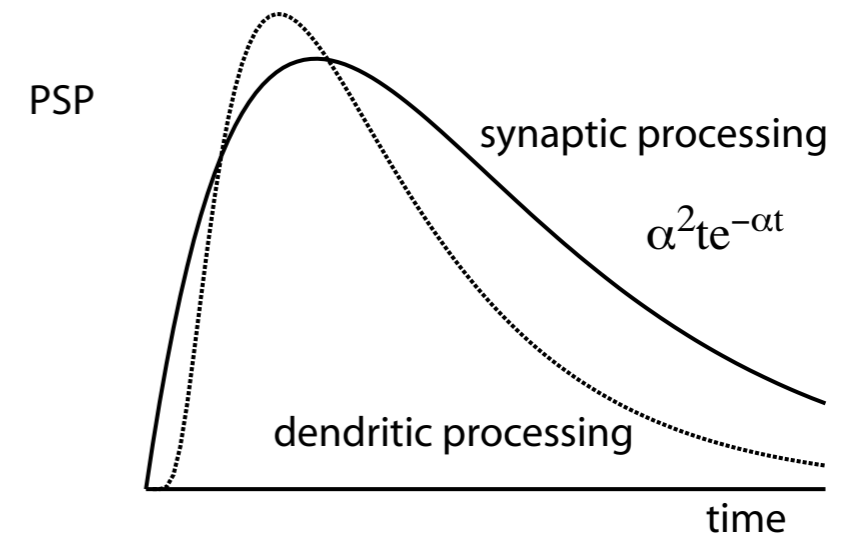
$$G(x, t) = w(x) \delta(t - |x|/v) \quad \rho = f \circ u$$

exploit
convolution
structure

Exponential synapse

$$\eta(t) = \alpha e^{-\alpha t}, \quad t > 0$$

$$\left(1 + \frac{1}{\alpha} \frac{\partial}{\partial t}\right) u = \psi$$



$$\begin{aligned} \psi(x, t) &= \int_{-\infty}^{\infty} w(x - y) f(u(y, t - |x - y|/v)) dy \\ &= \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} ds G(x - y, t - s) \rho(y, s) \end{aligned}$$

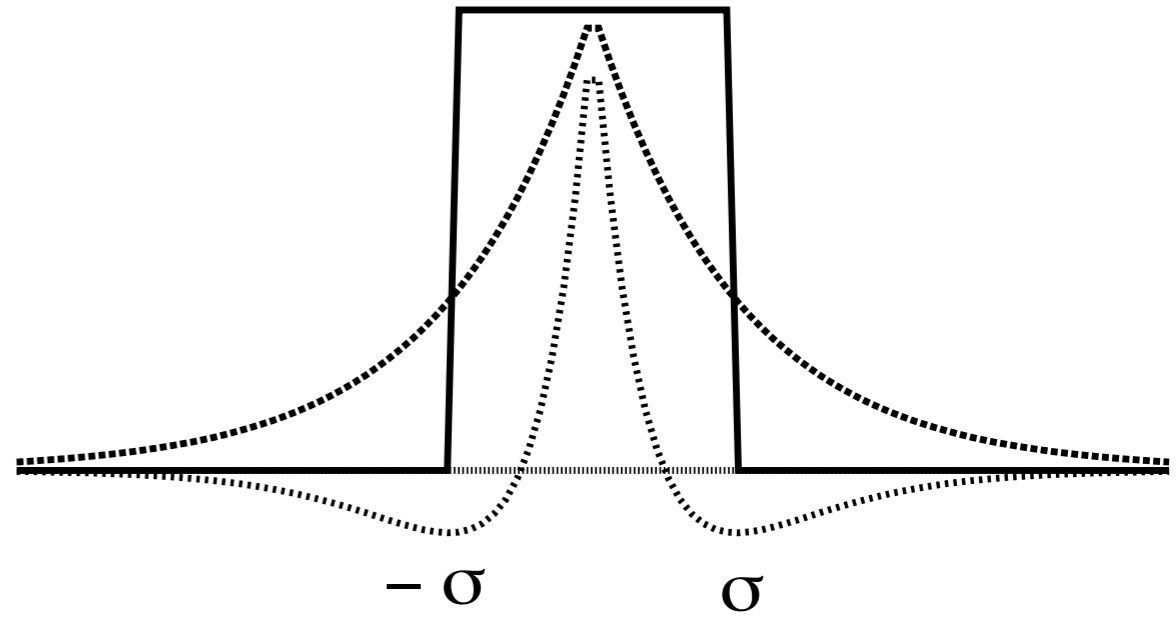
$$G(x, t) = w(x) \delta(t - |x|/v) \quad \rho = f \circ u$$

exploit
convolution
structure

$$\psi(k, \omega) = G(k, \omega) \rho(k, \omega)$$

Exponential anatomy

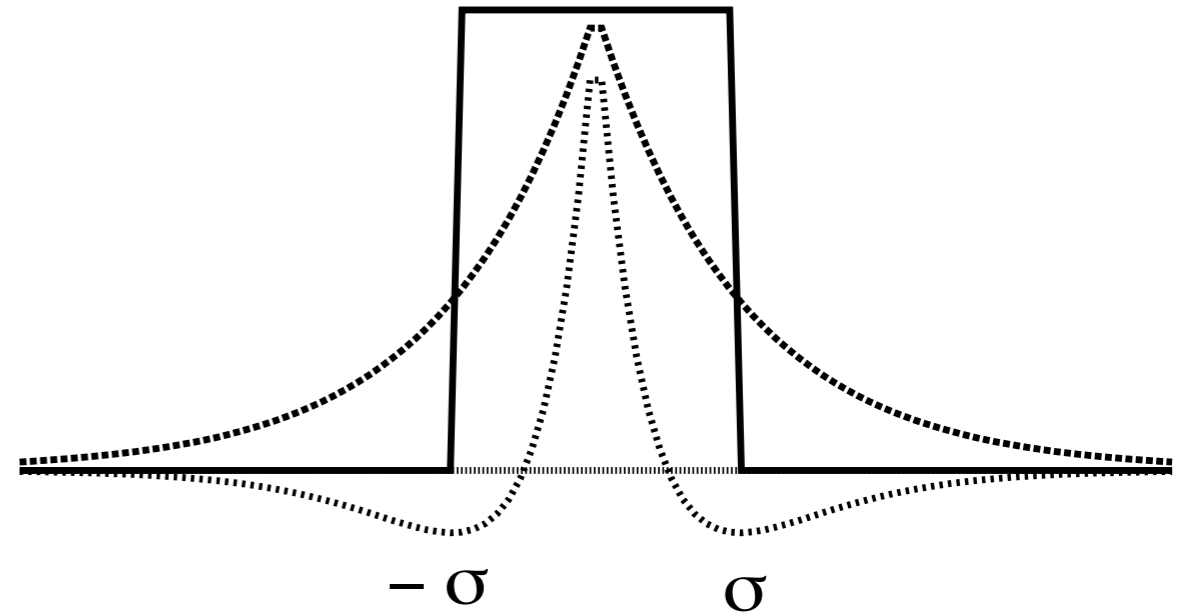
$$w(x) = e^{-|x|}/2$$



Exponential anatomy

$$w(x) = e^{-|x|}/2$$

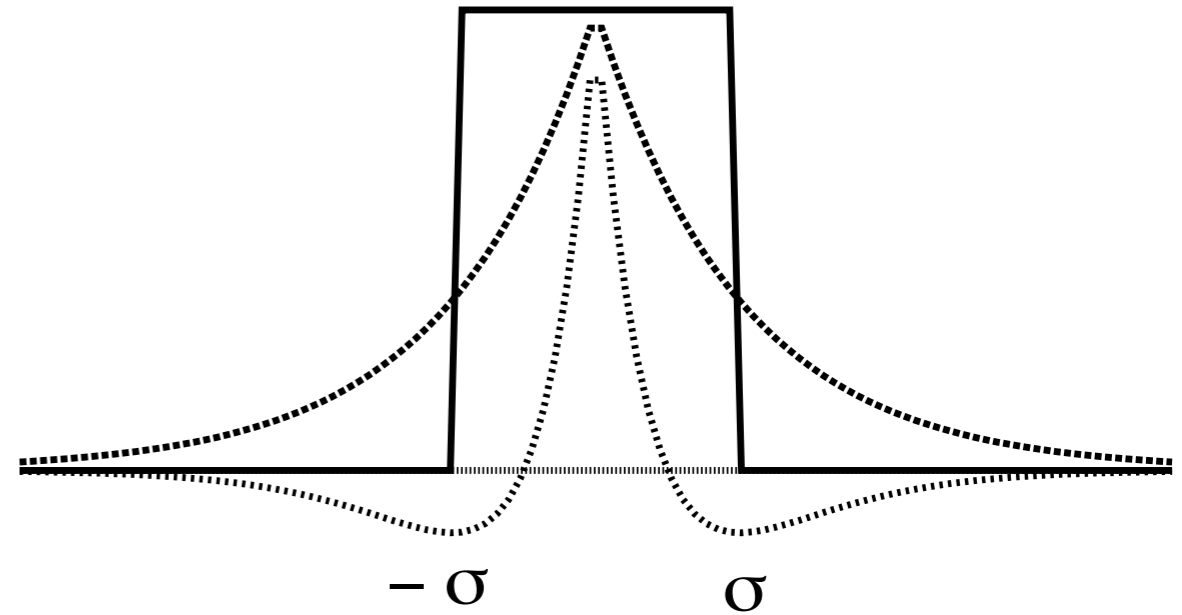
$$G(k, \omega) = \frac{1 + i\frac{\omega}{v}}{\left(1 + i\frac{\omega}{v}\right)^2 + k^2}$$



Exponential anatomy

$$w(x) = e^{-|x|}/2$$

$$G(k, \omega) = \frac{1 + i\frac{\omega}{v}}{\left(1 + i\frac{\omega}{v}\right)^2 + k^2}$$

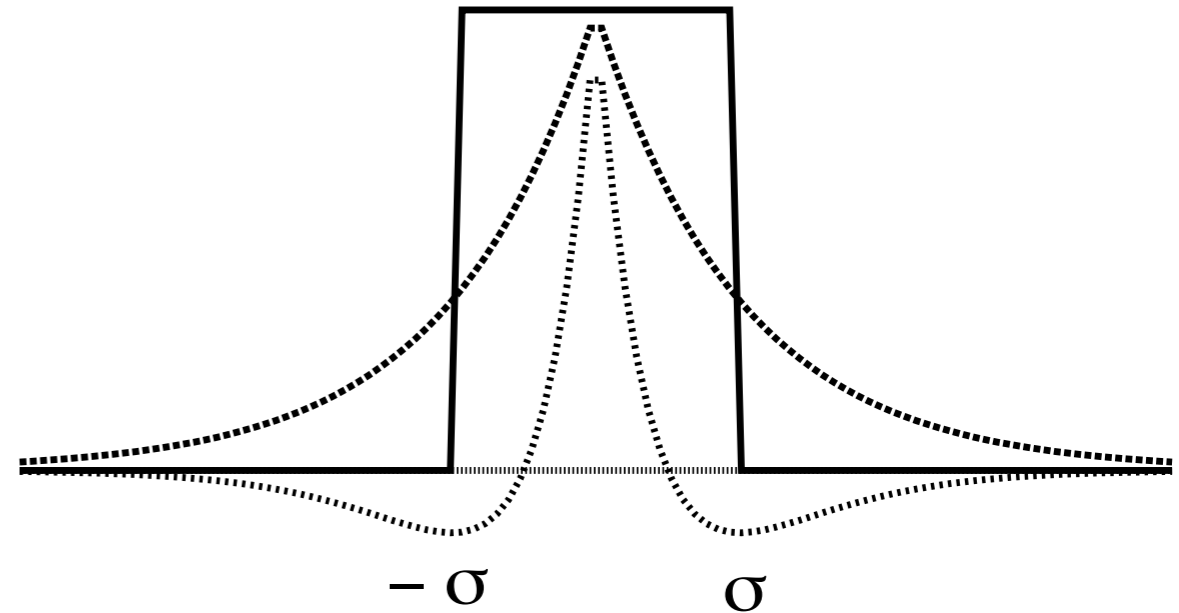


$$\left[\left(1 + i\frac{\omega}{v}\right)^2 + k^2 \right] \psi(k, \omega) = \left[1 + i\frac{\omega}{v} \right] \rho(k, \omega)$$

Exponential anatomy

$$w(x) = e^{-|x|}/2$$

$$G(k, \omega) = \frac{1 + i\frac{\omega}{v}}{\left(1 + i\frac{\omega}{v}\right)^2 + k^2}$$



$$\left[\left(1 + i\frac{\omega}{v}\right)^2 + k^2 \right] \psi(k, \omega) = \left[1 + i\frac{\omega}{v} \right] \rho(k, \omega)$$

$$k^2 \leftrightarrow -\nabla^2 \quad \text{and} \quad i\omega \leftrightarrow \partial_t$$

PDE in terms of the local operators ∇^2, ∂_t

$$[(v + \partial_t)^2 - v^2 \partial_{xx}] \psi(x, t) = [v^2 + v \partial_t] f \circ u(x, t)$$

Damped inhomogeneous brain-wave equation

2D Brain-wave equation

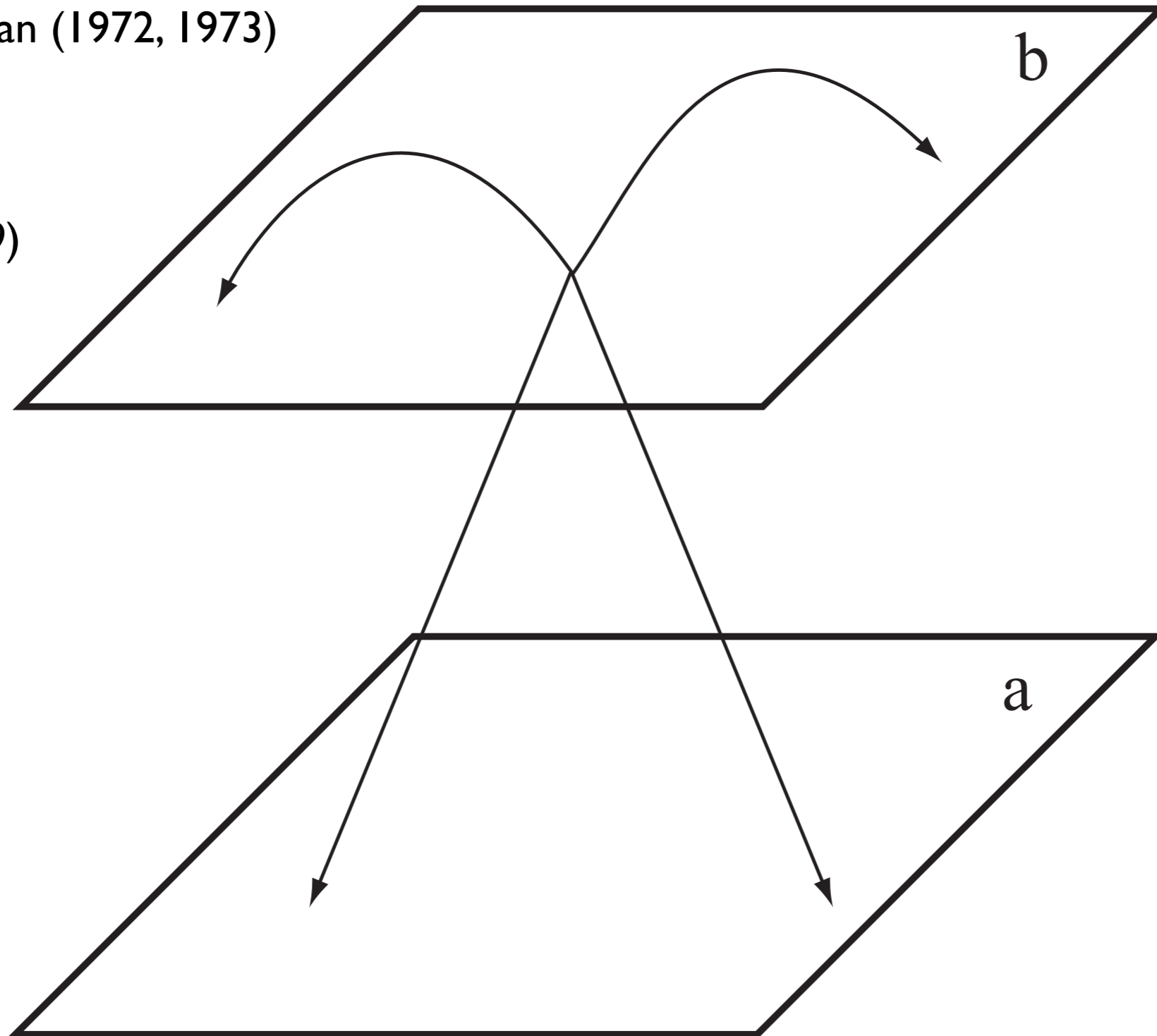
Wilson and Cowan (1972, 1973)

Nunez (1974)

Amari (1977)

Liley, Cadusch

and Wright (1999)



2D Brain-wave equation

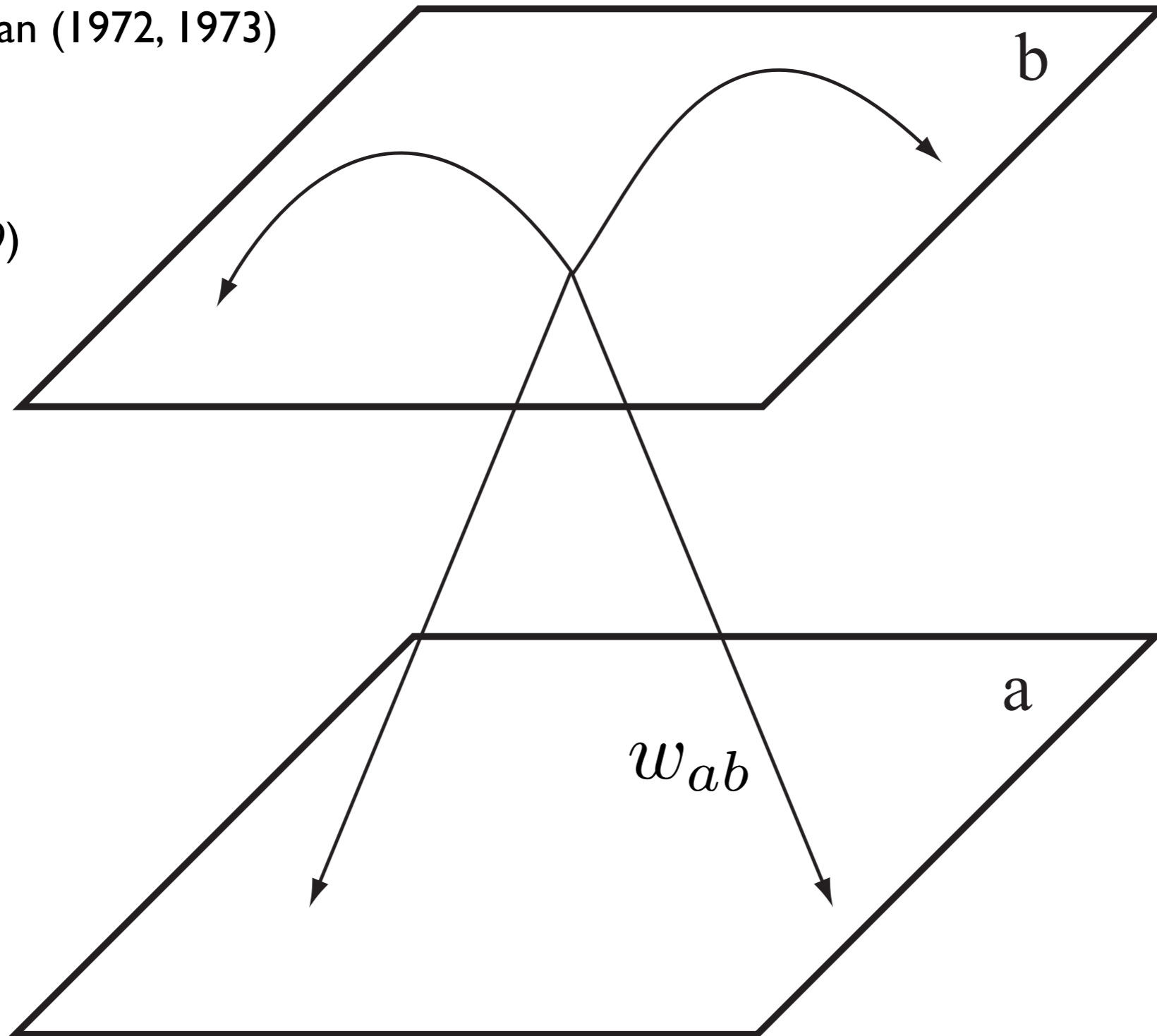
Wilson and Cowan (1972, 1973)

Nunez (1974)

Amari (1977)

Liley, Cadusch

and Wright (1999)



2D Brain-wave equation

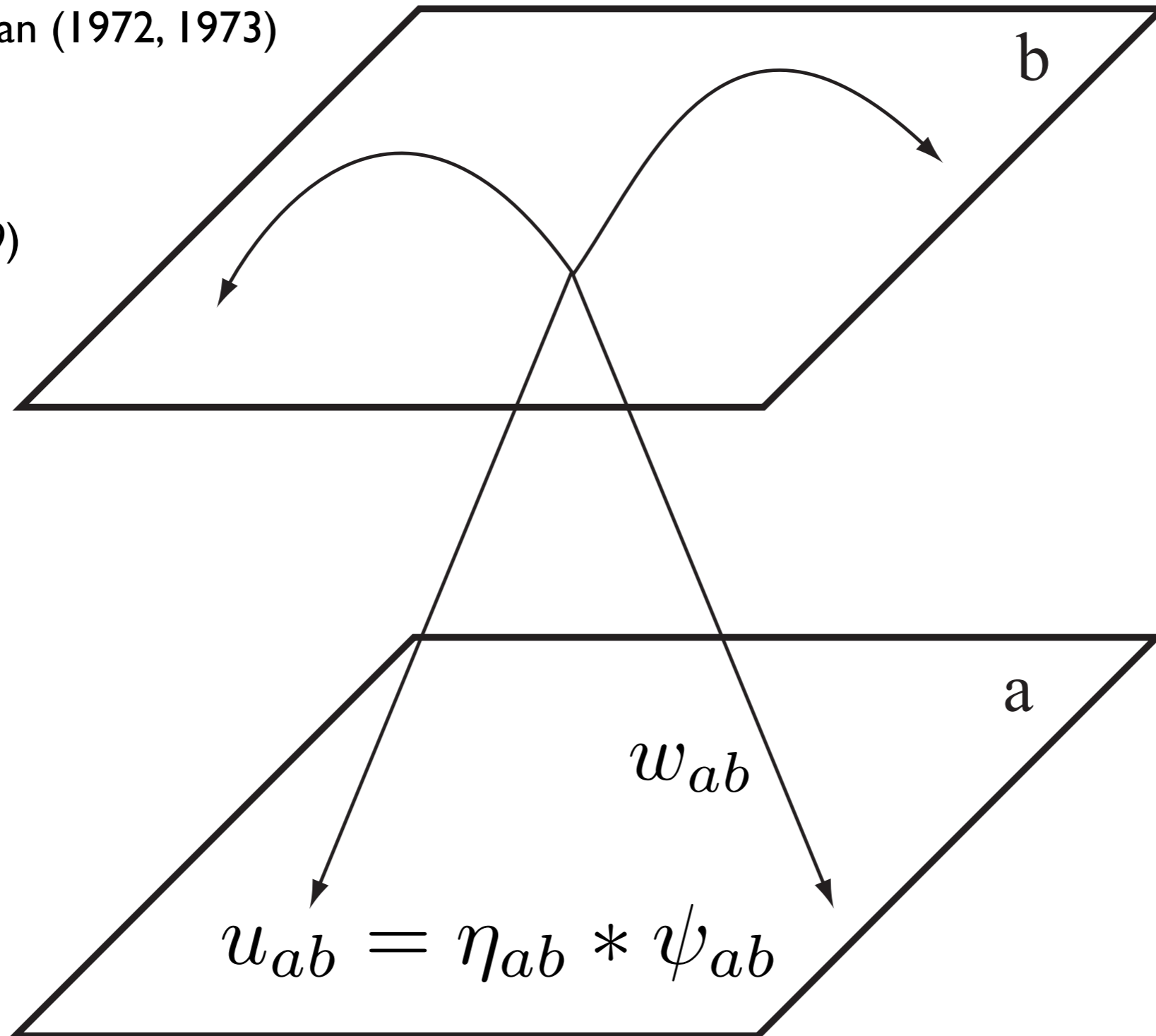
Wilson and Cowan (1972, 1973)

Nunez (1974)

Amari (1977)

Liley, Cadusch

and Wright (1999)



2D Brain-wave equation

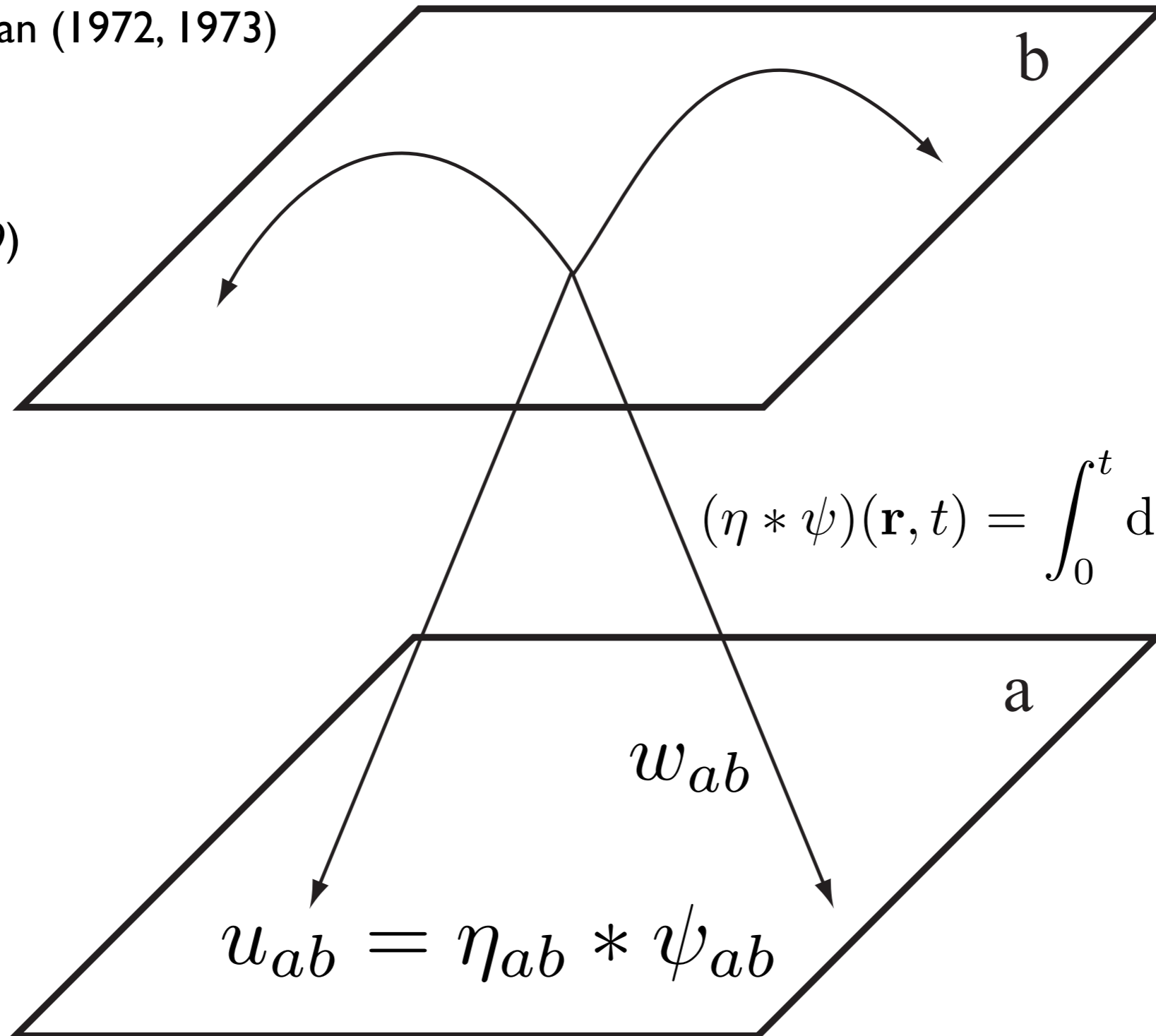
Wilson and Cowan (1972, 1973)

Nunez (1974)

Amari (1977)

Liley, Cadusch

and Wright (1999)



$$(\eta * \psi)(\mathbf{r}, t) = \int_0^t ds \eta(s) \psi(\mathbf{r}, t - s)$$

$$u_{ab} = \eta_{ab} * \psi_{ab}$$

2D Brain-wave equation

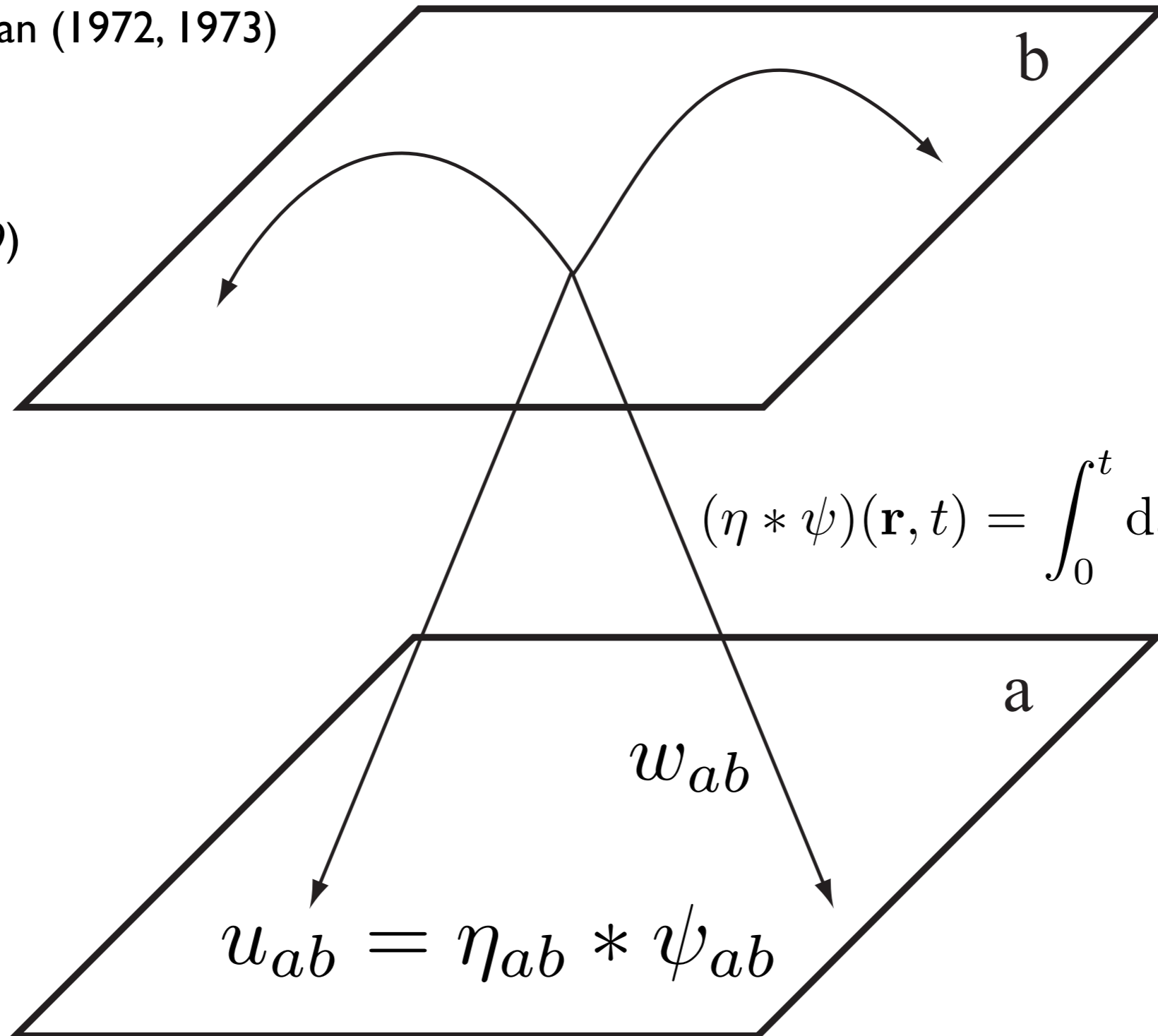
Wilson and Cowan (1972, 1973)

Nunez (1974)

Amari (1977)

Liley, Cadusch

and Wright (1999)



$$(\eta * \psi)(\mathbf{r}, t) = \int_0^t ds \eta(s) \psi(\mathbf{r}, t - s)$$

$$\psi_{ab}(\mathbf{r}, t) = \int_{\mathbb{R}^2} d\mathbf{r}' w_{ab}(\mathbf{r}, \mathbf{r}') f_b \circ h_b(\mathbf{r}', t - |\mathbf{r} - \mathbf{r}'|/v_{ab})$$

2D Brain-wave equation

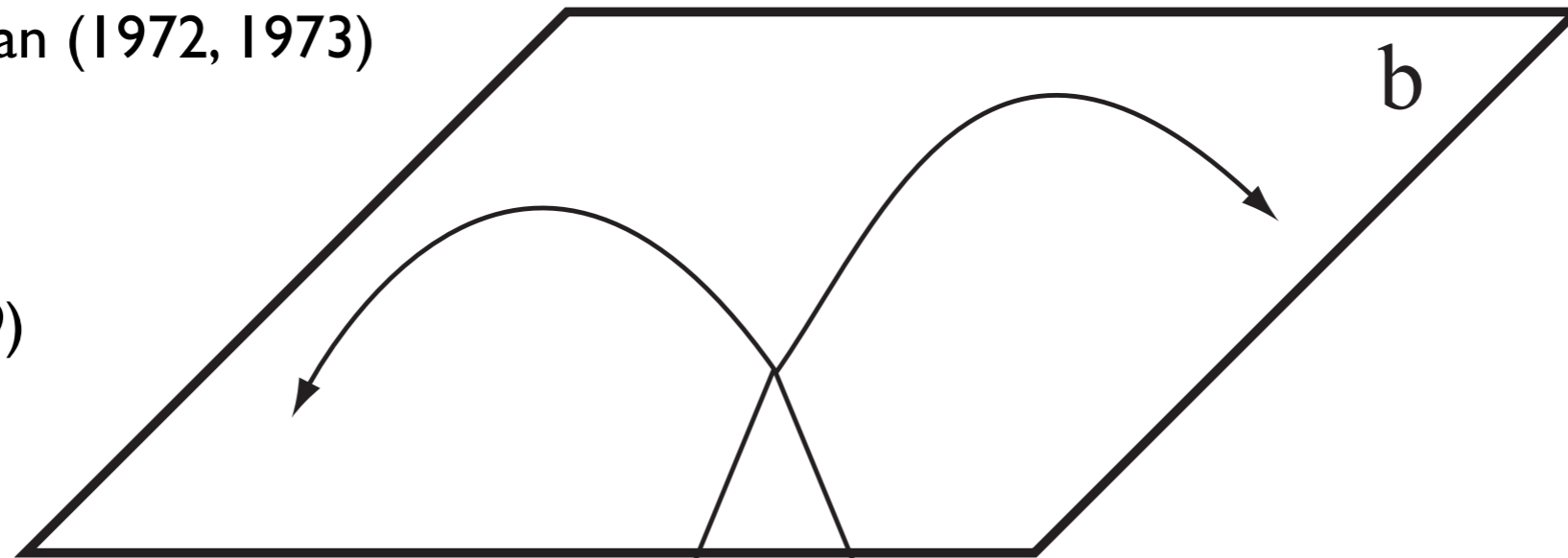
Wilson and Cowan (1972, 1973)

Nunez (1974)

Amari (1977)

Liley, Cadusch

and Wright (1999)



$$h_a = \sum_b u_{ab}$$

$$(\eta * \psi)(\mathbf{r}, t) = \int_0^t ds \eta(s) \psi(\mathbf{r}, t - s)$$

$$w_{ab}$$

$$u_{ab} = \eta_{ab} * \psi_{ab}$$

$$\psi_{ab}(\mathbf{r}, t) = \int_{\mathbb{R}^2} d\mathbf{r}' w_{ab}(\mathbf{r}, \mathbf{r}') f_b \circ h_b(\mathbf{r}', t - |\mathbf{r} - \mathbf{r}'|/v_{ab})$$

No such luck for
the natural choice

$$w_{ab}(r) = w_{ab}^0 e^{-r/\sigma_{ab}} / (2\pi)$$

$$G_{ab}(k, \omega) = w_{ab}^0 \frac{A_{ab}(\omega)}{(A_{ab}^2(\omega) + k^2)^{3/2}}$$

$$A_{ab}(\omega) = \sigma_{ab}^{-1} + i\omega/v_{ab}$$

No such luck for
the natural choice

$$w_{ab}(r) = w_{ab}^0 e^{-r/\sigma_{ab}} / (2\pi)$$

$$G_{ab}(k, \omega) = w_{ab}^0 \frac{A_{ab}(\omega)}{(A_{ab}^2(\omega) + k^2)^{3/2}}$$

$$A_{ab}(\omega) = \sigma_{ab}^{-1} + i\omega/v_{ab}$$

Introducing the operator \mathcal{A}_{ab} :

$$\mathcal{A}_{ab} = \left(\frac{1}{\sigma_{ab}} + \frac{1}{v_{ab}} \partial_t \right)^2$$

then the problem arises as how to interpret $[\mathcal{A}_{ab} - \nabla^2]^{3/2}$

No such luck for the natural choice

$$w_{ab}(r) = w_{ab}^0 e^{-r/\sigma_{ab}} / (2\pi)$$

$$G_{ab}(k, \omega) = w_{ab}^0 \frac{A_{ab}(\omega)}{(A_{ab}^2(\omega) + k^2)^{3/2}}$$

$$A_{ab}(\omega) = \sigma_{ab}^{-1} + i\omega/v_{ab}$$

Introducing the operator \mathcal{A}_{ab} :

$$\mathcal{A}_{ab} = \left(\frac{1}{\sigma_{ab}} + \frac{1}{v_{ab}} \partial_t \right)^2$$

then the problem arises as how to interpret $[\mathcal{A}_{ab} - \nabla^2]^{3/2}$

Long-wavelength approximation - expand around $k = 0$

$$\left(\mathcal{A}_{ab} - \frac{3}{2} \nabla^2 \right) \psi_{ab} = w_{ab}^0 \rho_b$$

No such luck for the natural choice

$$w_{ab}(r) = w_{ab}^0 e^{-r/\sigma_{ab}} / (2\pi)$$

$$G_{ab}(k, \omega) = w_{ab}^0 \frac{A_{ab}(\omega)}{(A_{ab}^2(\omega) + k^2)^{3/2}}$$

$$A_{ab}(\omega) = \sigma_{ab}^{-1} + i\omega/v_{ab}$$

Introducing the operator \mathcal{A}_{ab} :

$$\mathcal{A}_{ab} = \left(\frac{1}{\sigma_{ab}} + \frac{1}{v_{ab}} \partial_t \right)^2$$

then the problem arises as how to interpret $[\mathcal{A}_{ab} - \nabla^2]^{3/2}$

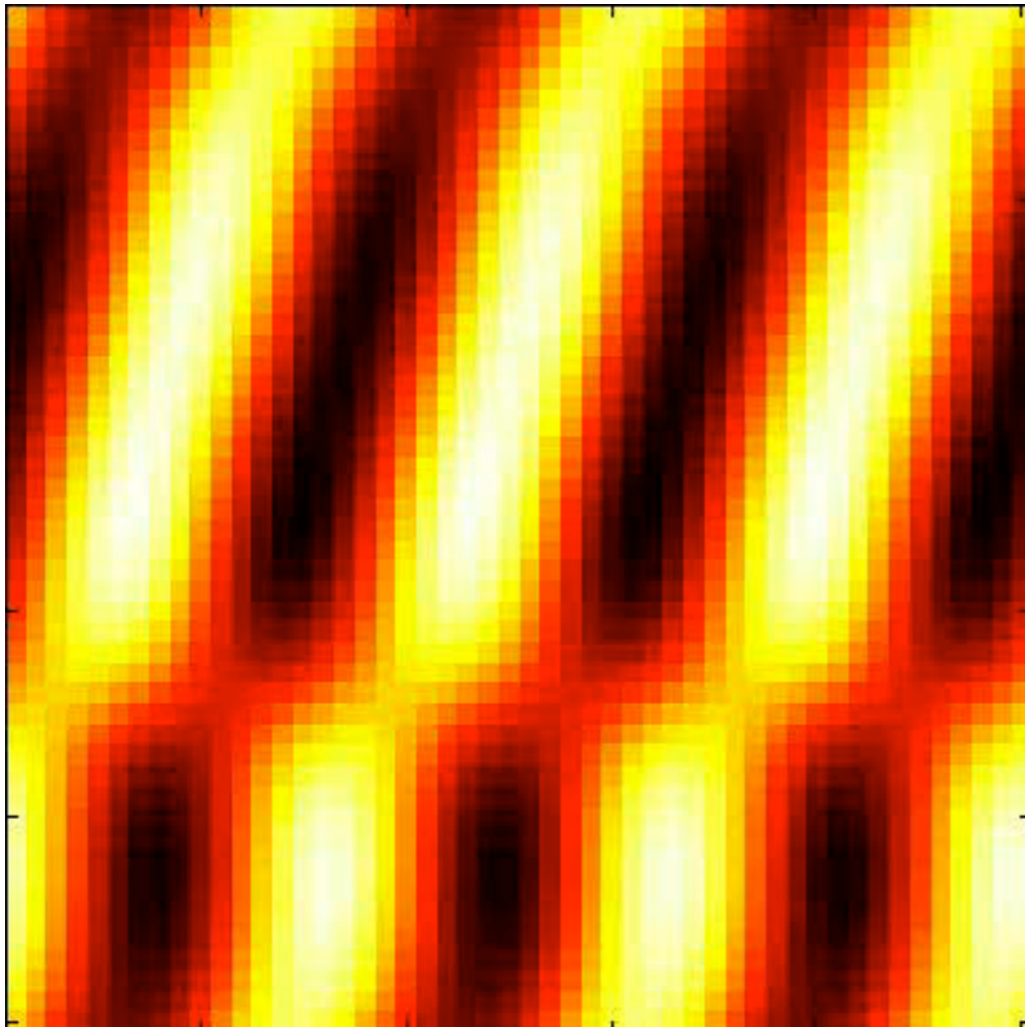
Long-wavelength approximation - expand around $k = 0$

$$\left(\mathcal{A}_{ab} - \frac{3}{2} \nabla^2 \right) \psi_{ab} = w_{ab}^0 \rho_b$$

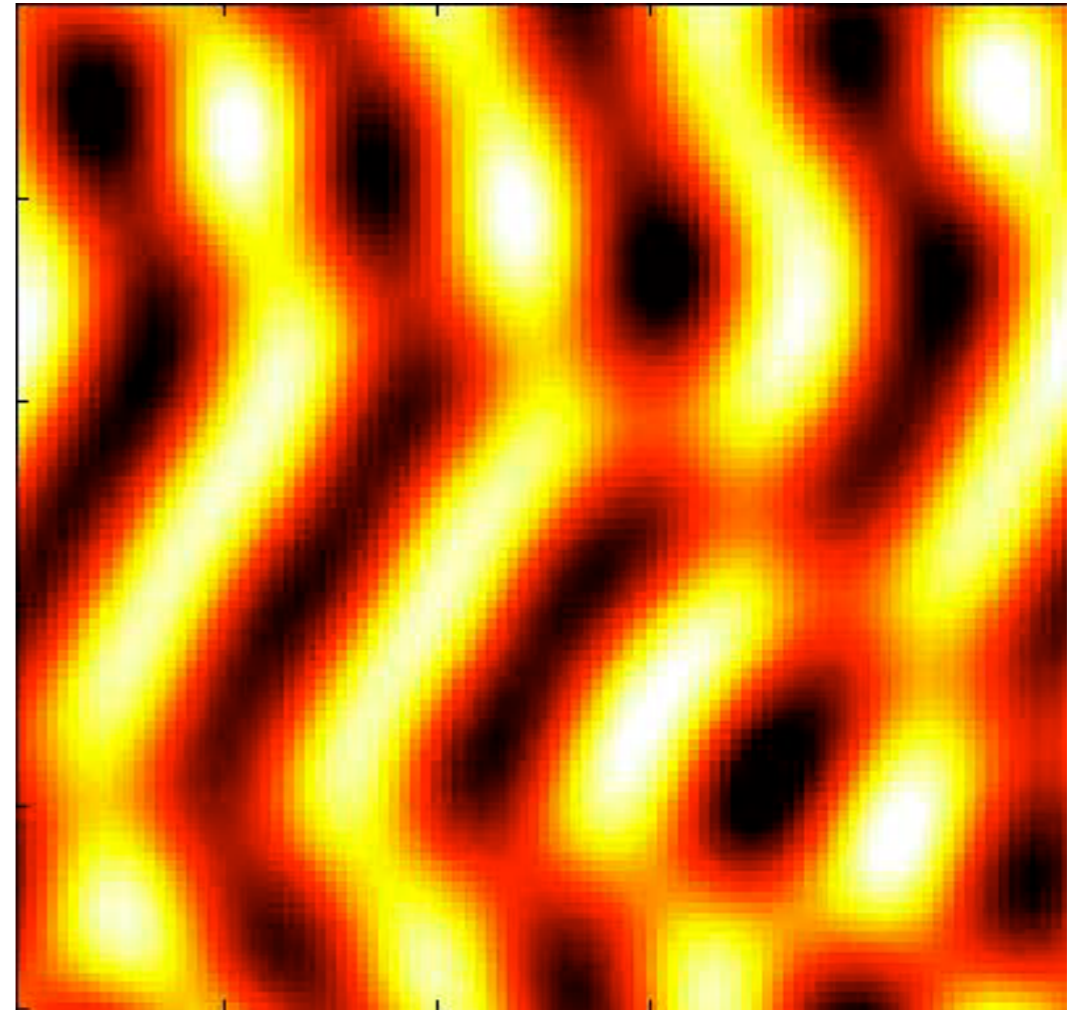
Better *rational* approximation

S Coombes, N A Venkov, L Shiau, I Bojak, D T J Liley and C R Laing 2007 Modeling electrocortical activity through improved local approximations of integral neural field equations, Physical Review E, Vol 76, 051901-8. (Open access).

Eye candy (for a two-layer E-I network)



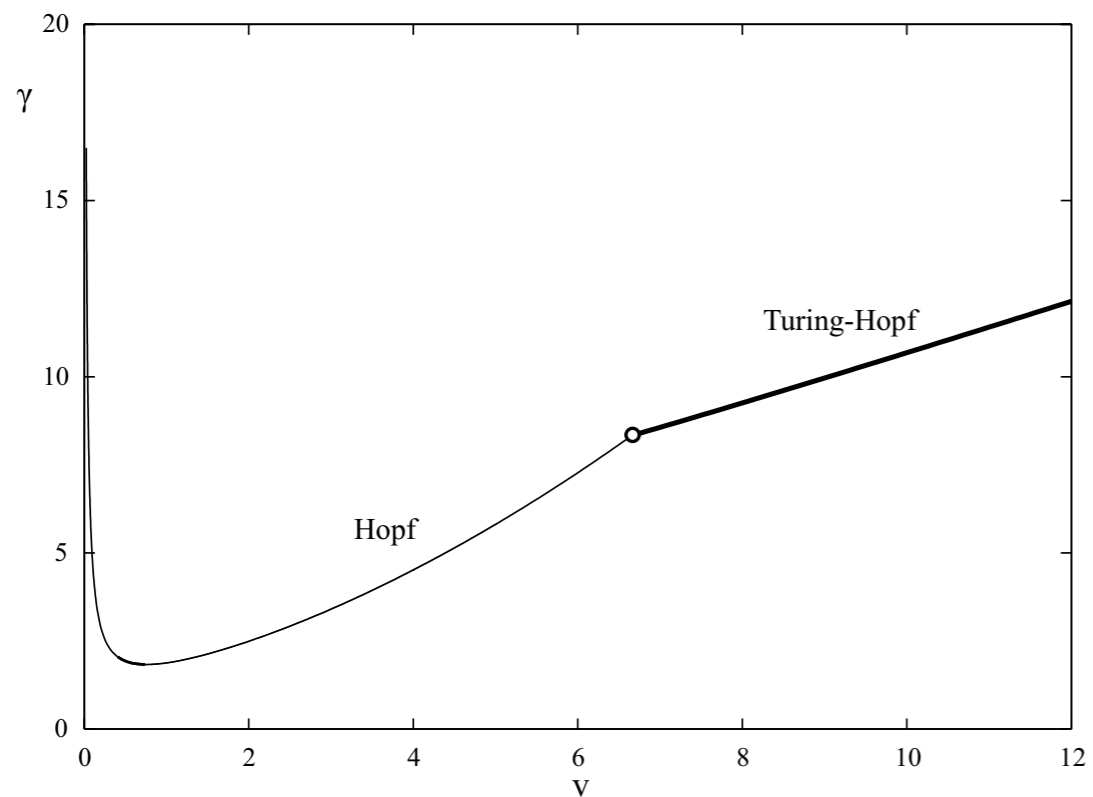
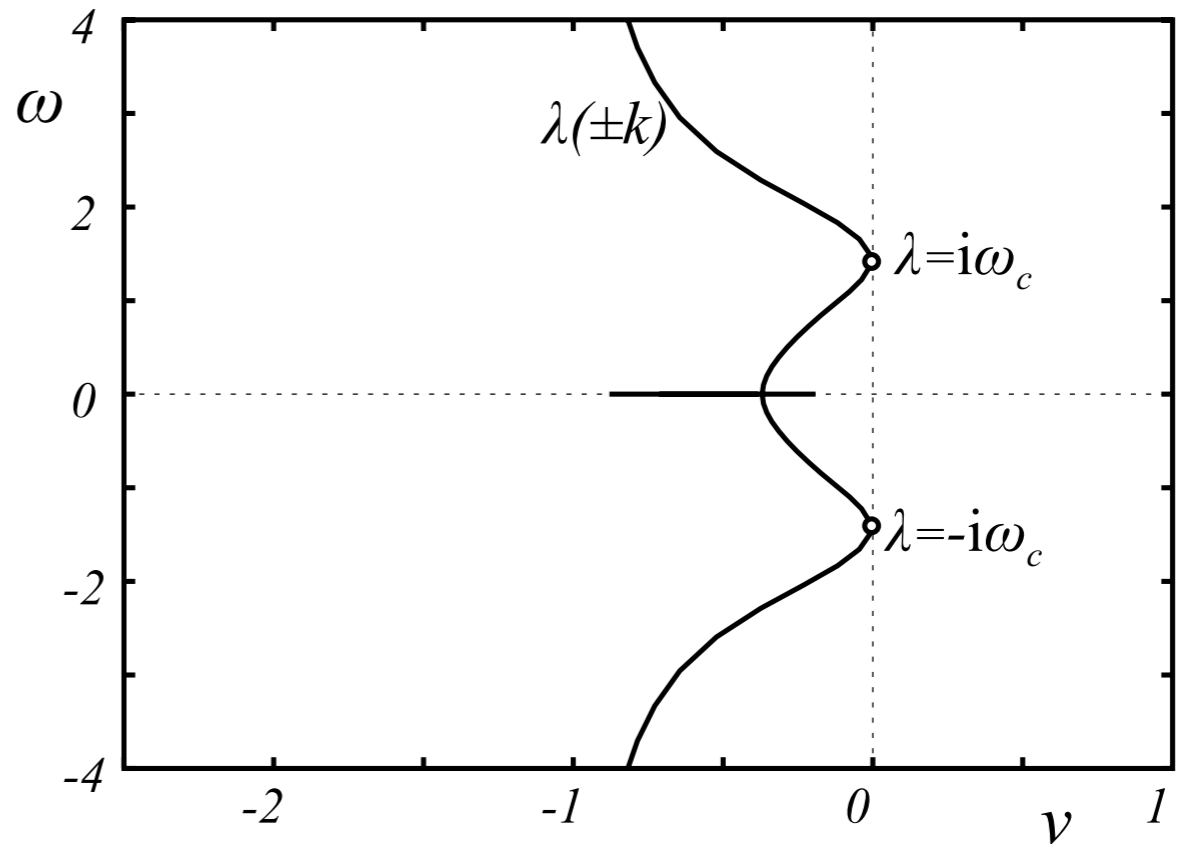
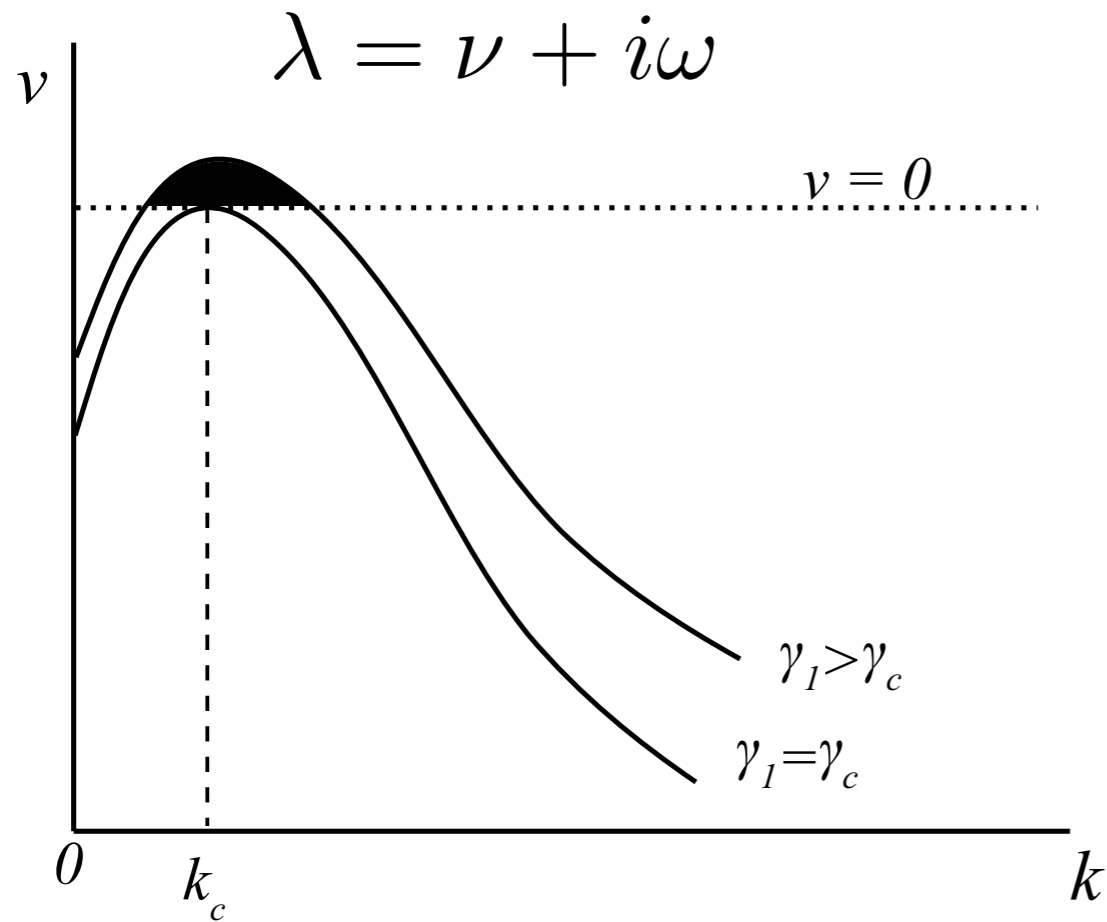
Long-wavelength



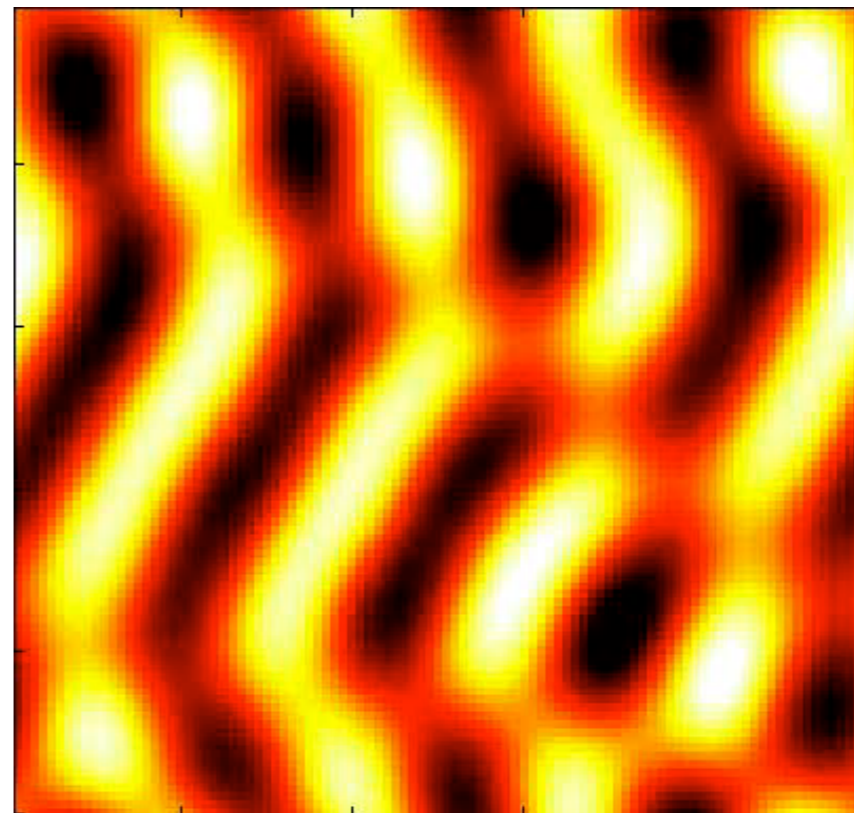
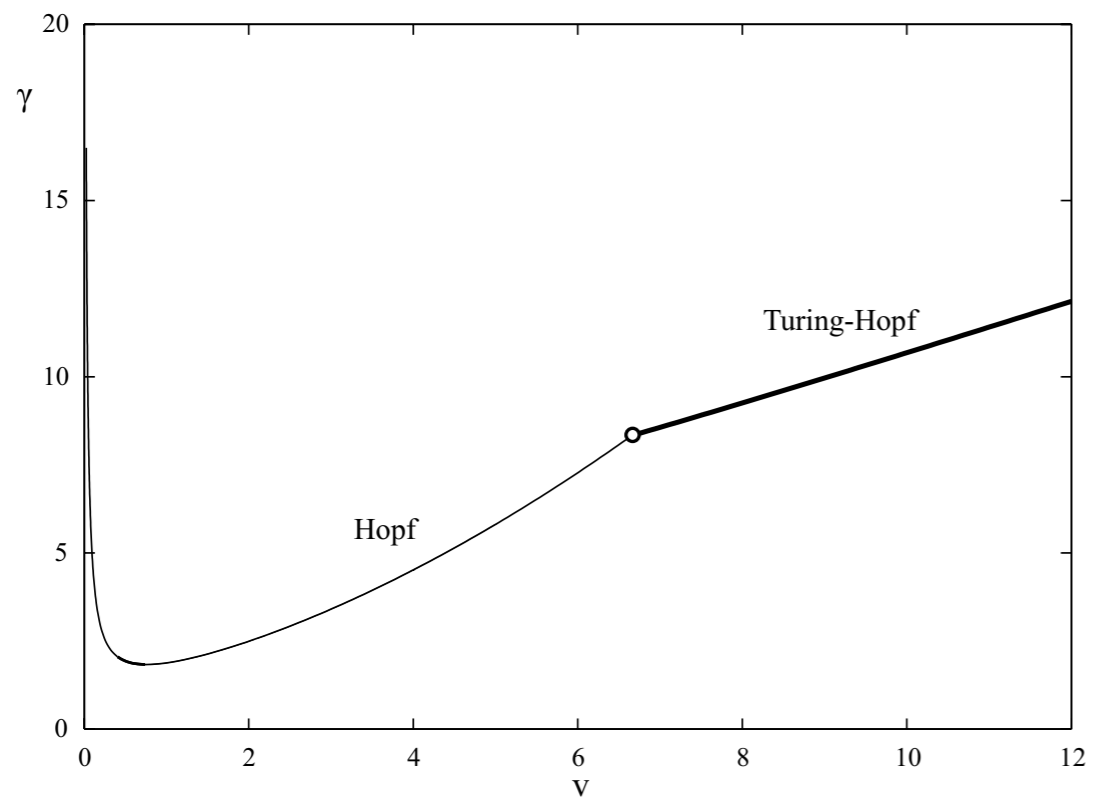
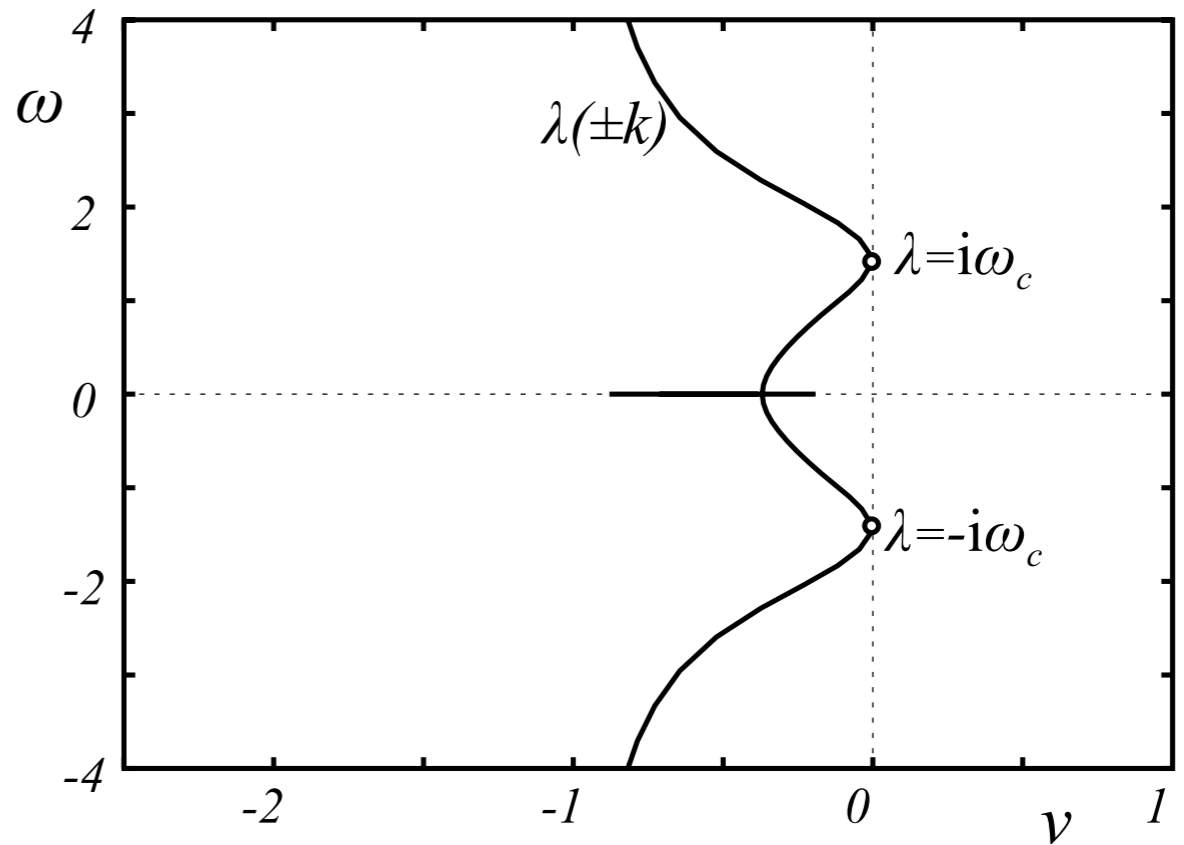
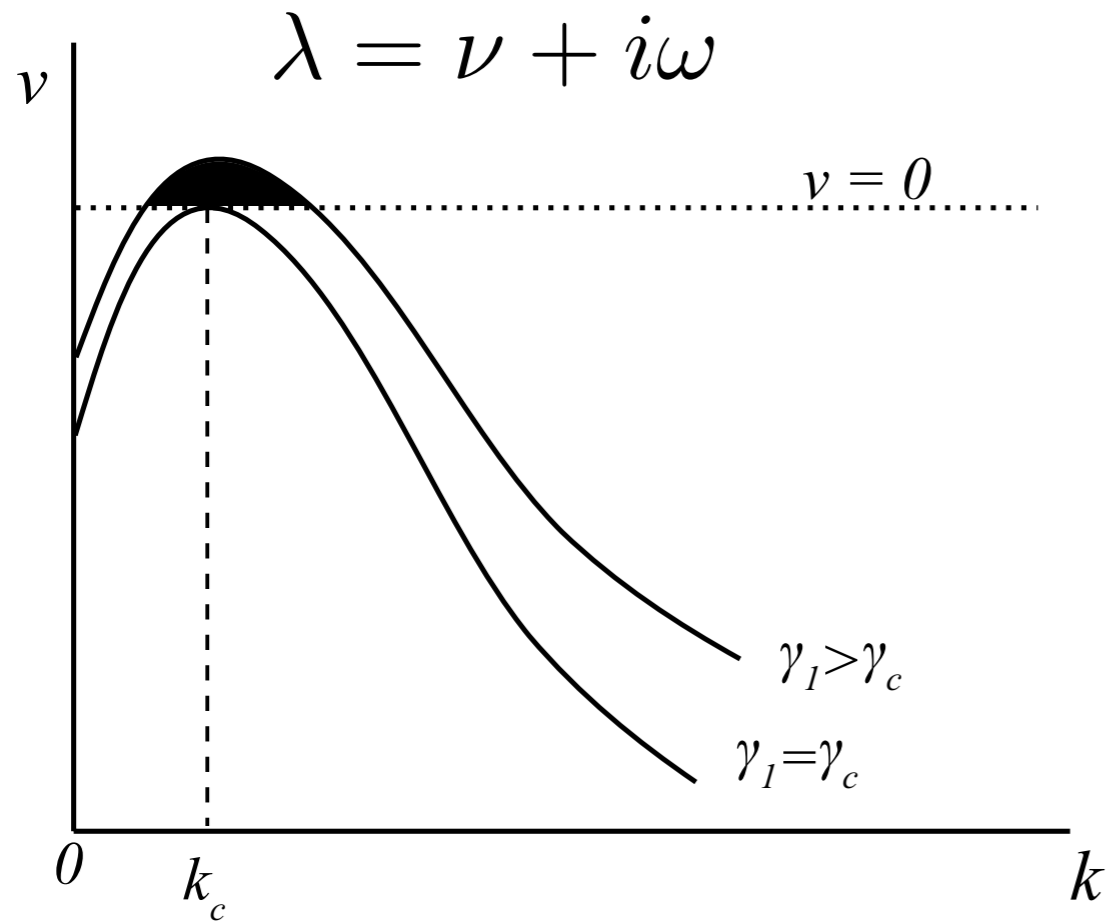
Rational

Dynamic Turing instability

$$e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$$

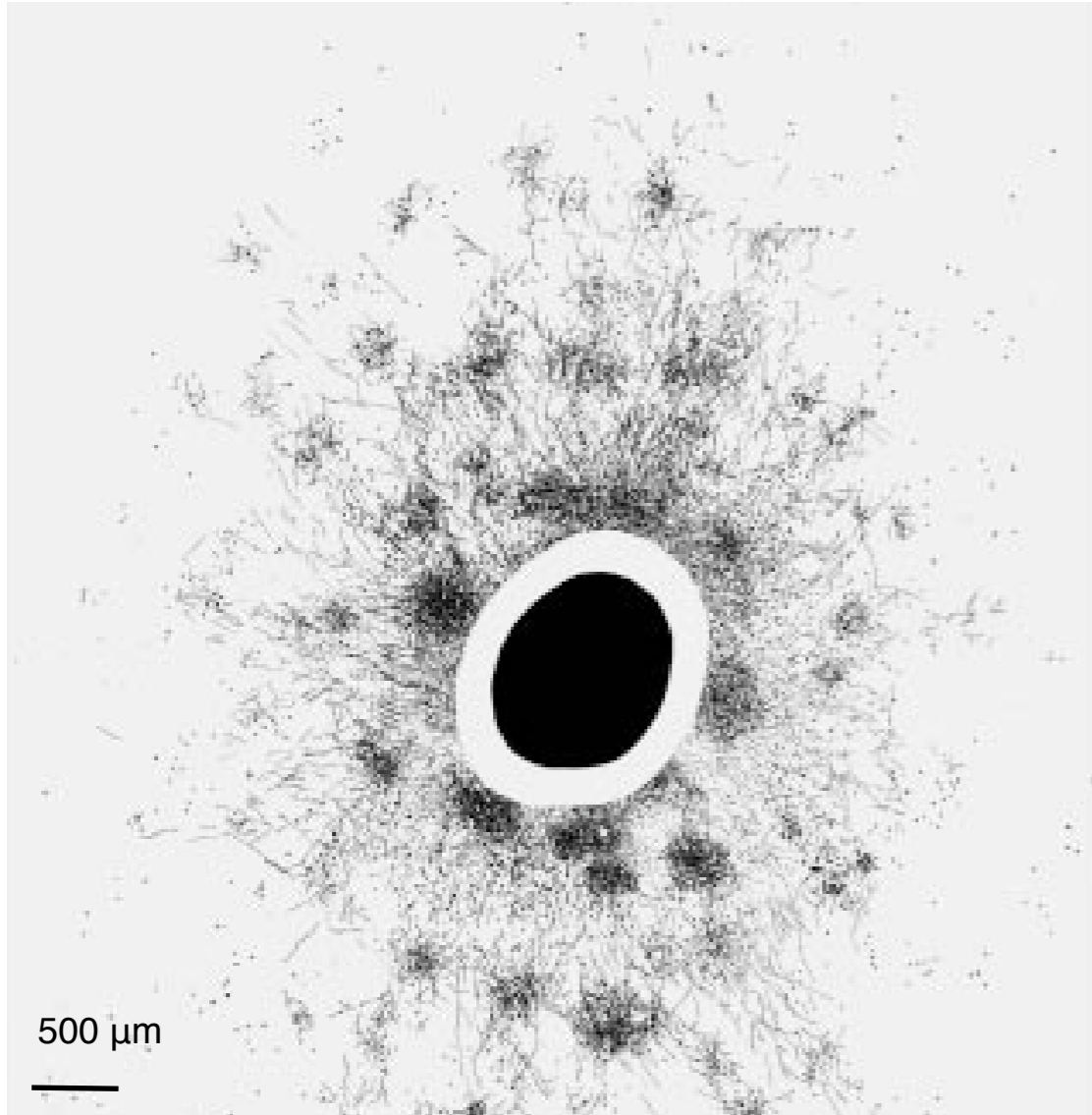


Dynamic Turing instability $e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$



Spatial modulation

Neocortex has a crystalline micro-structure at the mm length scale, so that the assumption of isotropic connectivity has to be revised.

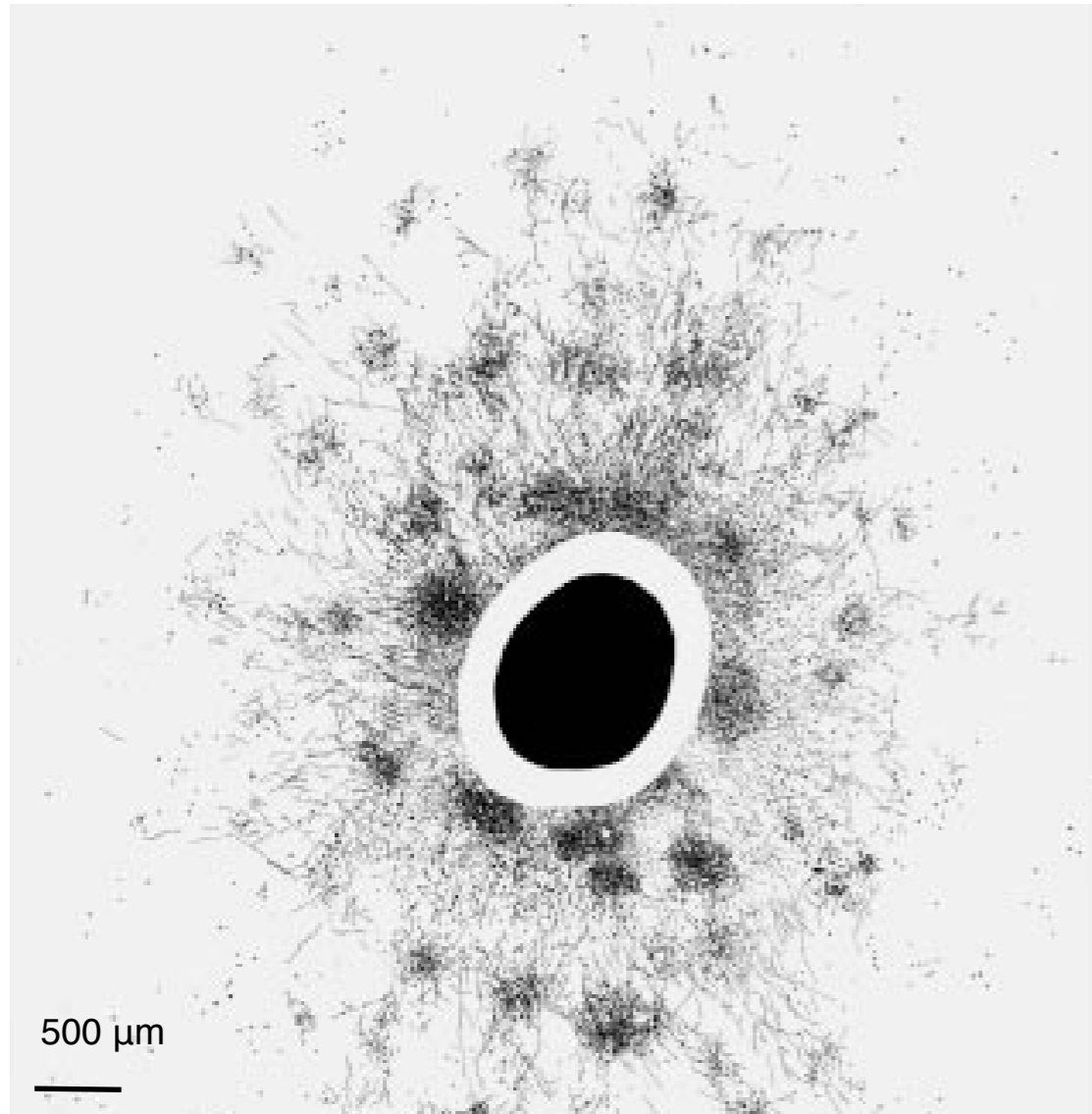


Anisotropy in a patchy connectional field:
reconstruction of a tangential section through layers 2/3 of macaque area V1, showing a cholera toxin B tracer injection site and surrounding transported orthograde and retrograde label.

J. S. Lund, A. Angelucci and P. C. Bressloff, Anatomical substrates for the functional column in macaque primary visual cortex. *Cerebral Cortex* 12:15–24 (2003)

Spatial modulation

Neocortex has a crystalline micro-structure at the mm length scale, so that the assumption of isotropic connectivity has to be revised.



Anisotropy in a patchy connectional field: reconstruction of a tangential section through layers 2/3 of macaque area V1, showing a cholera toxin B tracer injection site and surrounding transported orthograde and retrograde label.

J. S. Lund, A. Angelucci and P. C. Bressloff, Anatomical substrates for the functional column in macaque primary visual cortex. *Cerebral Cortex* 12:15–24 (2003)

The patchy nature of the horizontal connections immediately implies breaking of continuous rotation symmetry but not necessarily continuous translation symmetry.

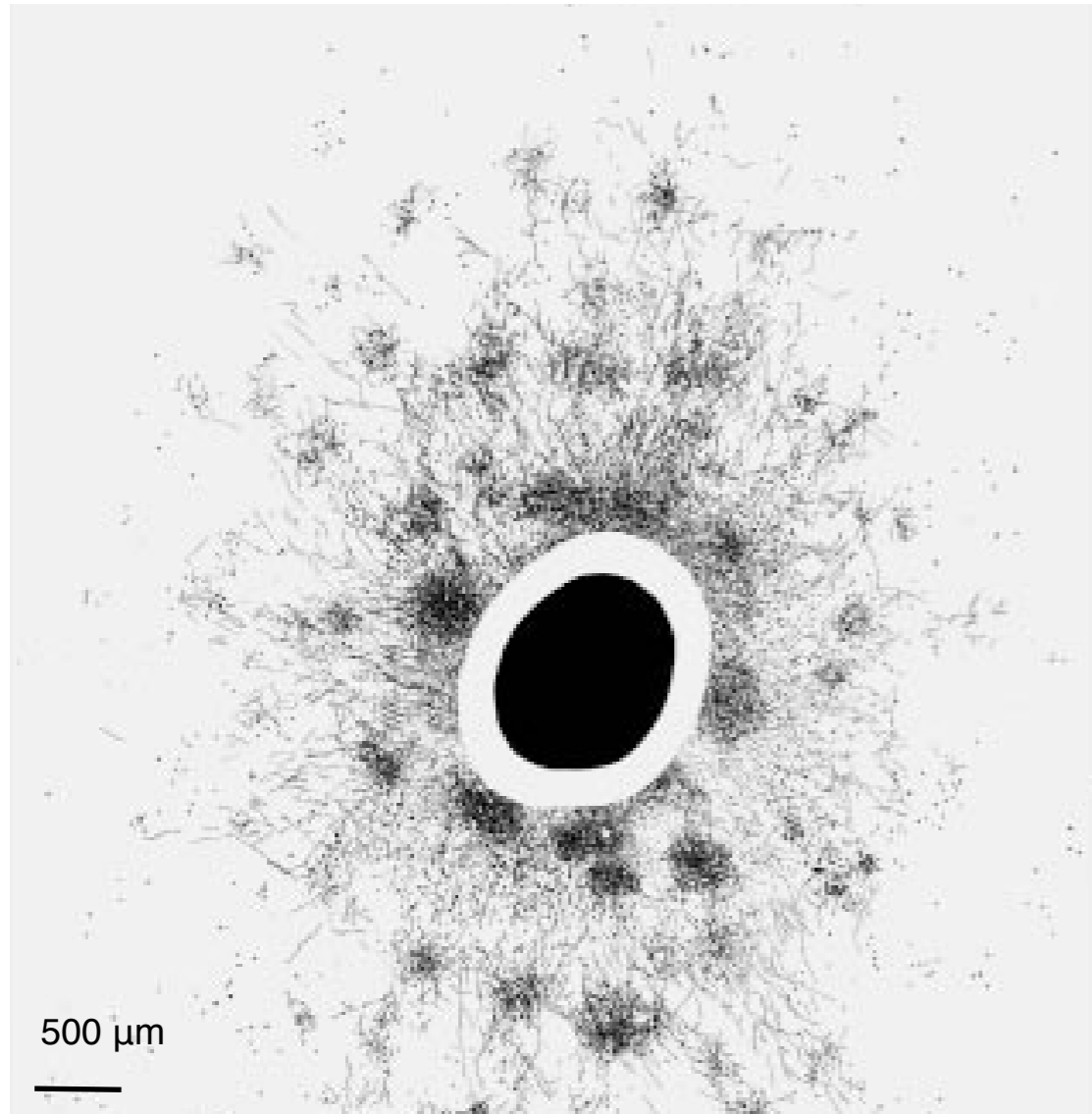
Introduce a periodically modulated spatial kernel of the form

$$w_{ab}^{\text{P}}(\mathbf{r}, \mathbf{r}') = w_{ab}(|\mathbf{r} - \mathbf{r}'|) J_{ab}(\mathbf{r} - \mathbf{r}')$$

where $J_{ab}(\mathbf{r})$ varies periodically with respect to a regular planar lattice \mathcal{L} , with reciprocal lattice vectors \mathbf{q}

Spatial modulation

Neocortex has a crystalline micro-structure at the mm length scale, so that the assumption of isotropic connectivity has to be revised.



Anisotropy in a patchy connectional field: reconstruction of a tangential section through layers 2/3 of macaque area V1, showing a cholera toxin B tracer injection site and surrounding transported orthograde and retrograde label.

J. S. Lund, A. Angelucci and P. C. Bressloff, Anatomical substrates for the functional column in macaque primary visual cortex. *Cerebral Cortex* 12:15–24 (2003)

The patchy nature of the horizontal connections immediately implies breaking of continuous rotation symmetry but not necessarily continuous translation symmetry.

Introduce a periodically modulated spatial kernel of the form

$$w_{ab}^{\text{P}}(\mathbf{r}, \mathbf{r}') = w_{ab}(|\mathbf{r} - \mathbf{r}'|) J_{ab}(\mathbf{r} - \mathbf{r}')$$

where $J_{ab}(\mathbf{r})$ varies periodically with respect to a regular planar lattice \mathcal{L} , with reciprocal lattice vectors \mathbf{q}

Patchy propagators - P A Robinson, *Physical Review E*, 73:041904, 2006.

Linear stability analysis ... $e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$

same formal approach as before – λ now depends on the direction as well as the magnitude of \mathbf{k}

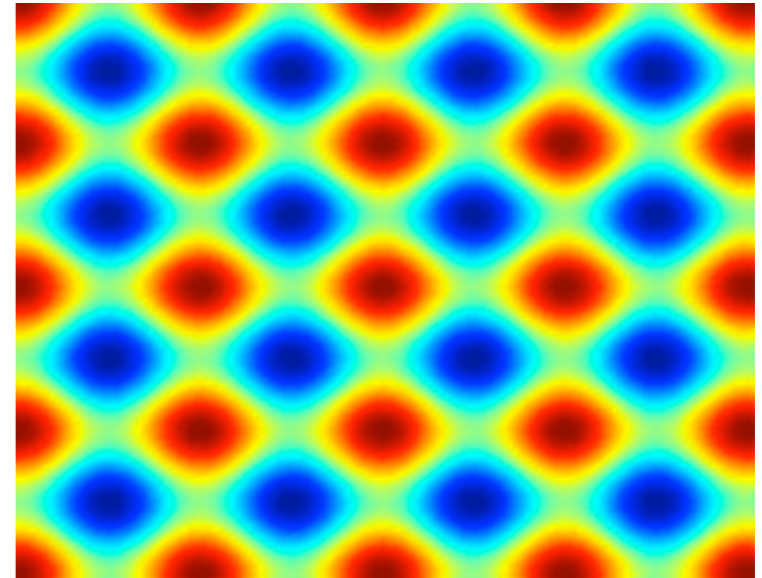
Linear stability analysis ... $e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$

same formal approach as before – λ now depends on the direction as well as the magnitude of \mathbf{k}

e.g. consider a square lattice with spacing d and choose

$$J_{ab}(\mathbf{r}) = [\cos(\mathbf{k}_1 \cdot \mathbf{r}) + \cos(\mathbf{k}_2 \cdot \mathbf{r})]/2$$

$$\mathbf{k}_1 = 2\pi/d(1, 0) \text{ and } \mathbf{k}_2 = 2\pi/d(0, 1)$$



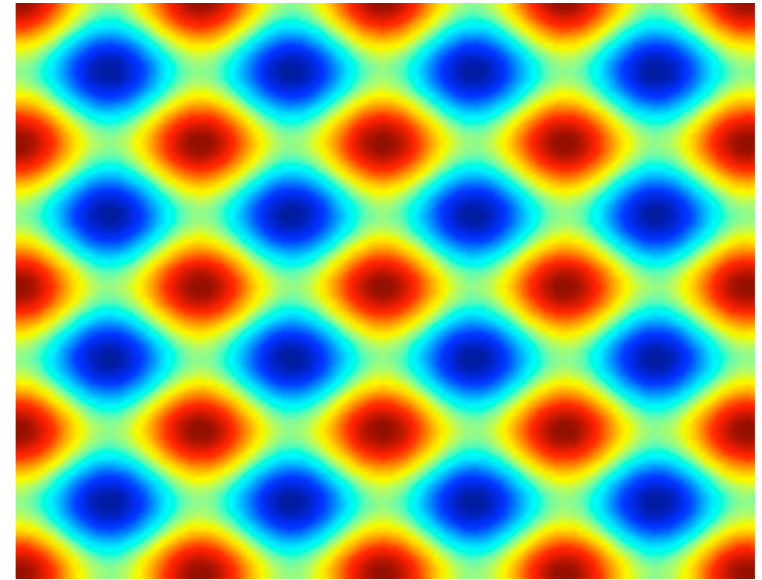
Linear stability analysis ... $e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$

same formal approach as before – λ now depends on the direction as well as the magnitude of \mathbf{k}

e.g. consider a square lattice with spacing d and choose

$$J_{ab}(\mathbf{r}) = [\cos(\mathbf{k}_1 \cdot \mathbf{r}) + \cos(\mathbf{k}_2 \cdot \mathbf{r})]/2$$

$$\mathbf{k}_1 = 2\pi/d(1, 0) \text{ and } \mathbf{k}_2 = 2\pi/d(0, 1)$$



$$J_{ab}^{\mathbf{q}} = [\delta(\mathbf{q} - \mathbf{k}_1) + \delta(\mathbf{q} + \mathbf{k}_1) + \delta(\mathbf{q} - \mathbf{k}_2) + \delta(\mathbf{q} + \mathbf{k}_2)]/4$$

...and we need only consider the modes indexed by $\mathbf{k}_{1,2}$

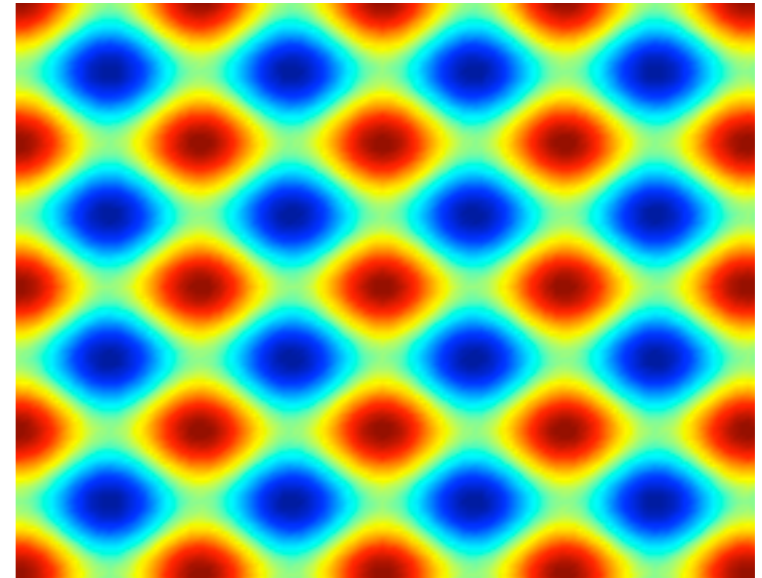
Linear stability analysis ... $e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$

same formal approach as before – λ now depends on the direction as well as the magnitude of \mathbf{k}

e.g. consider a square lattice with spacing d and choose

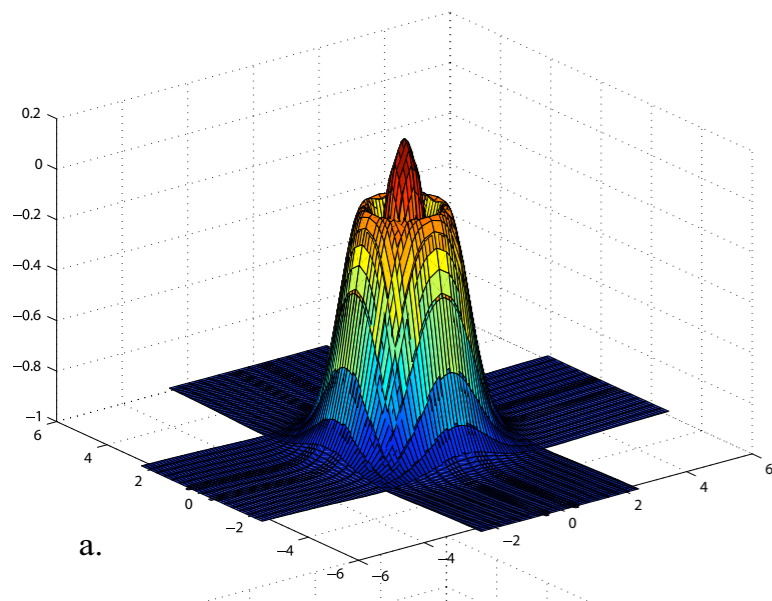
$$J_{ab}(\mathbf{r}) = [\cos(\mathbf{k}_1 \cdot \mathbf{r}) + \cos(\mathbf{k}_2 \cdot \mathbf{r})]/2$$

$$\mathbf{k}_1 = 2\pi/d(1, 0) \text{ and } \mathbf{k}_2 = 2\pi/d(0, 1)$$



$$J_{ab}^{\mathbf{q}} = [\delta(\mathbf{q} - \mathbf{k}_1) + \delta(\mathbf{q} + \mathbf{k}_1) + \delta(\mathbf{q} - \mathbf{k}_2) + \delta(\mathbf{q} + \mathbf{k}_2)]/4$$

... and we need only consider the modes indexed by $\mathbf{k}_{1,2}$



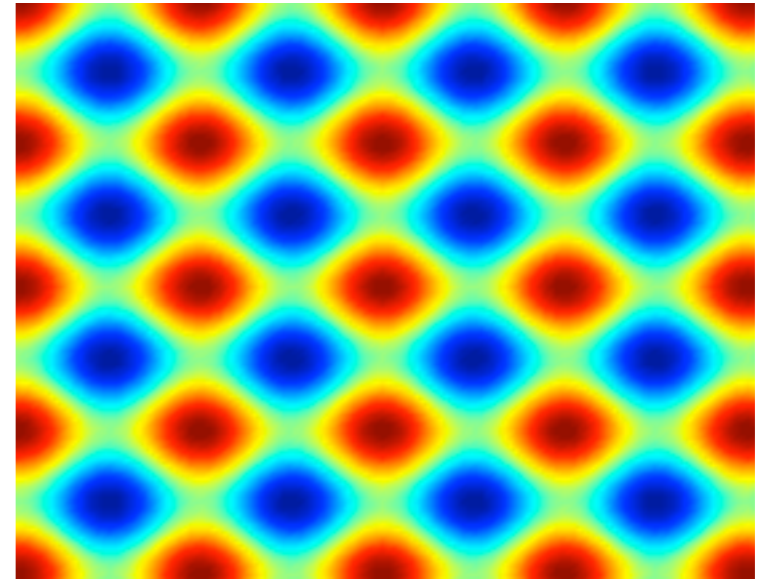
Linear stability analysis ... $e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$

same formal approach as before – λ now depends on the direction as well as the magnitude of \mathbf{k}

e.g. consider a square lattice with spacing d and choose

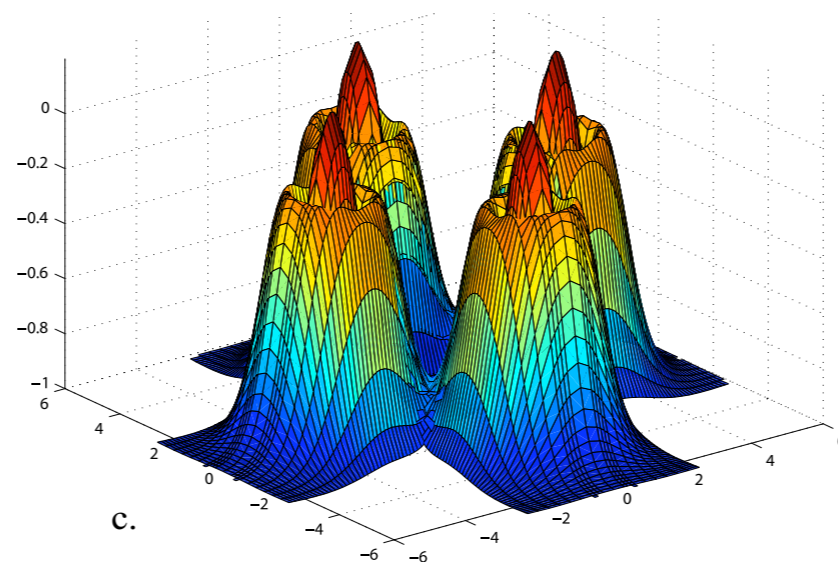
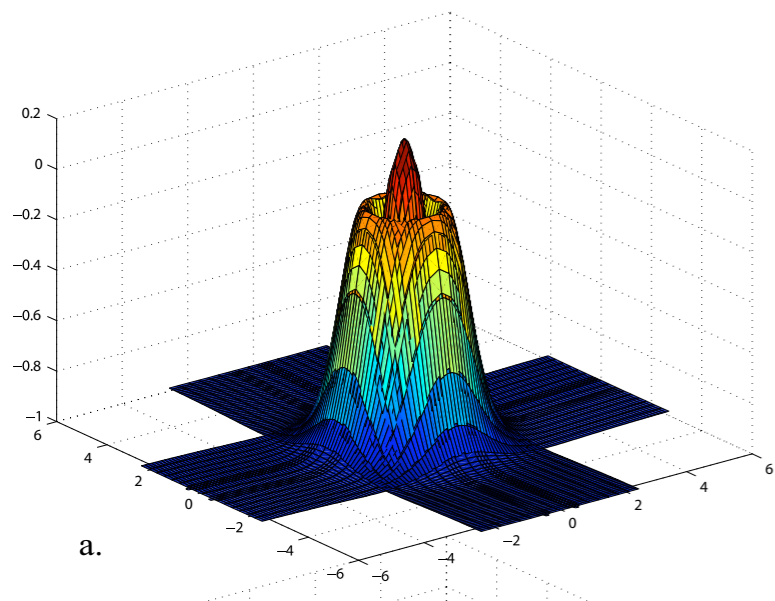
$$J_{ab}(\mathbf{r}) = [\cos(\mathbf{k}_1 \cdot \mathbf{r}) + \cos(\mathbf{k}_2 \cdot \mathbf{r})]/2$$

$$\mathbf{k}_1 = 2\pi/d(1, 0) \text{ and } \mathbf{k}_2 = 2\pi/d(0, 1)$$



$$J_{ab}^{\mathbf{q}} = [\delta(\mathbf{q} - \mathbf{k}_1) + \delta(\mathbf{q} + \mathbf{k}_1) + \delta(\mathbf{q} - \mathbf{k}_2) + \delta(\mathbf{q} + \mathbf{k}_2)]/4$$

... and we need only consider the modes indexed by $\mathbf{k}_{1,2}$



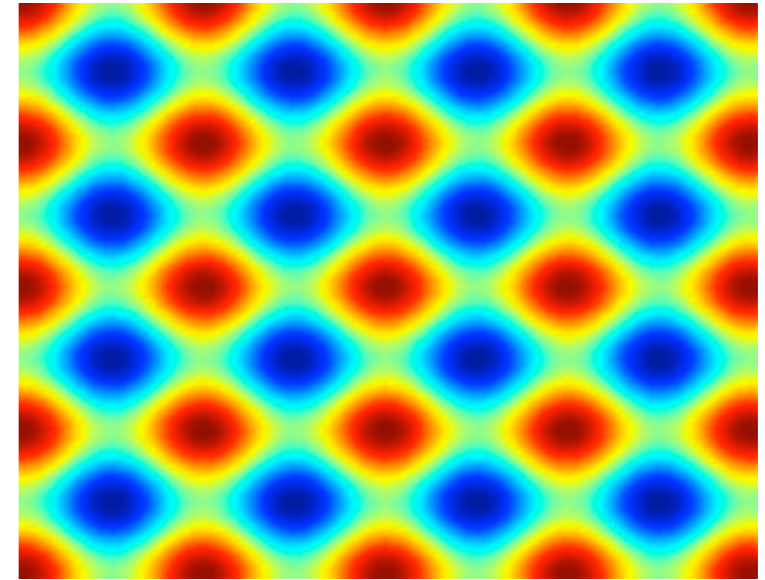
Linear stability analysis ... $e^{\lambda t} e^{i\mathbf{k}\cdot\mathbf{r}}$

same formal approach as before – λ now depends on the direction as well as the magnitude of \mathbf{k}

e.g. consider a square lattice with spacing d and choose

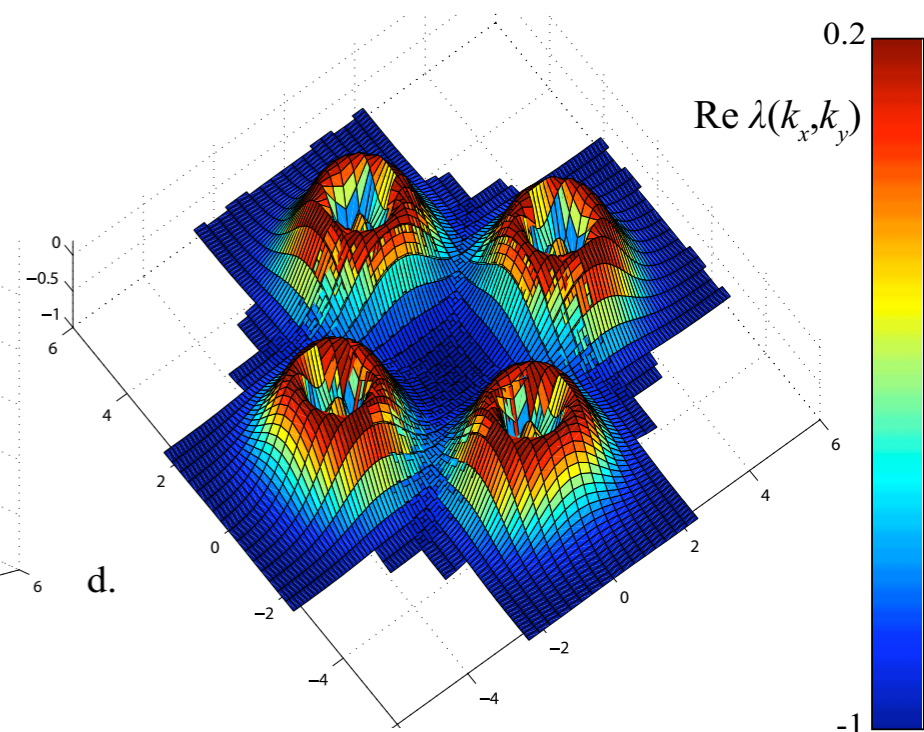
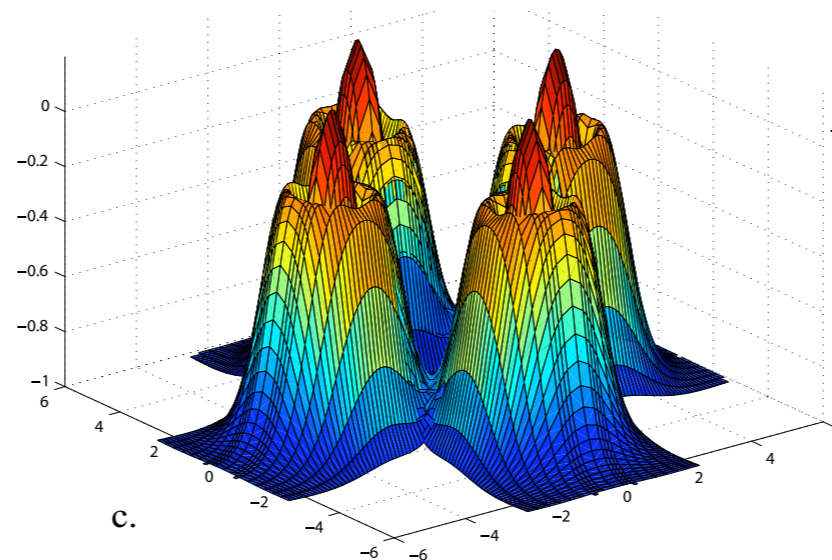
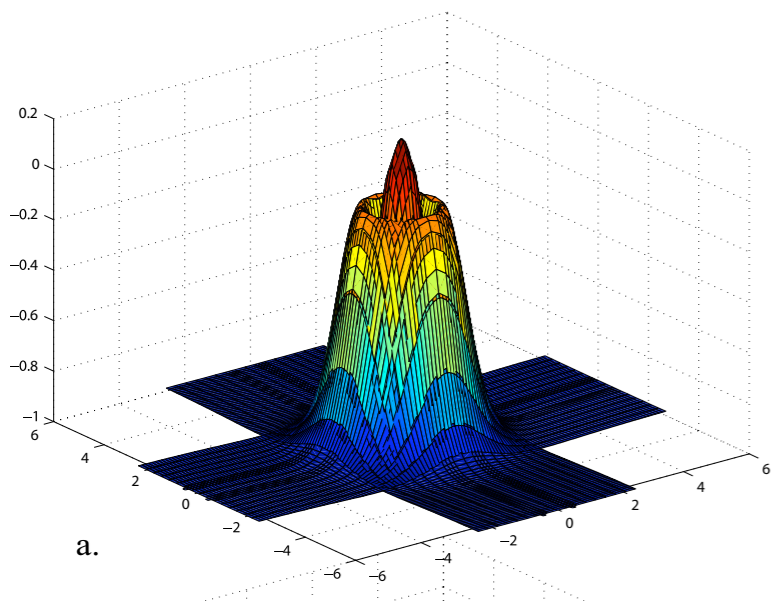
$$J_{ab}(\mathbf{r}) = [\cos(\mathbf{k}_1 \cdot \mathbf{r}) + \cos(\mathbf{k}_2 \cdot \mathbf{r})]/2$$

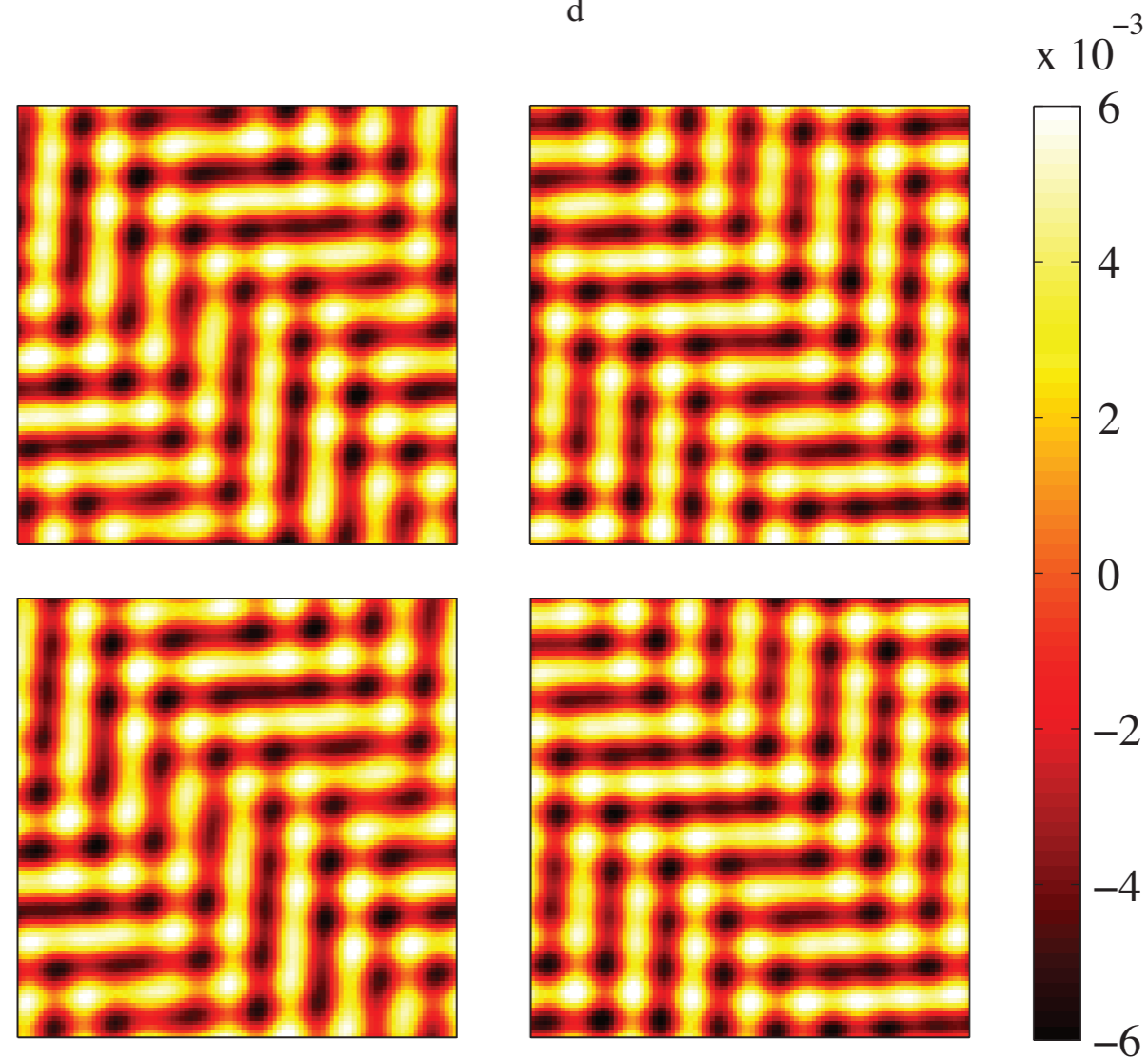
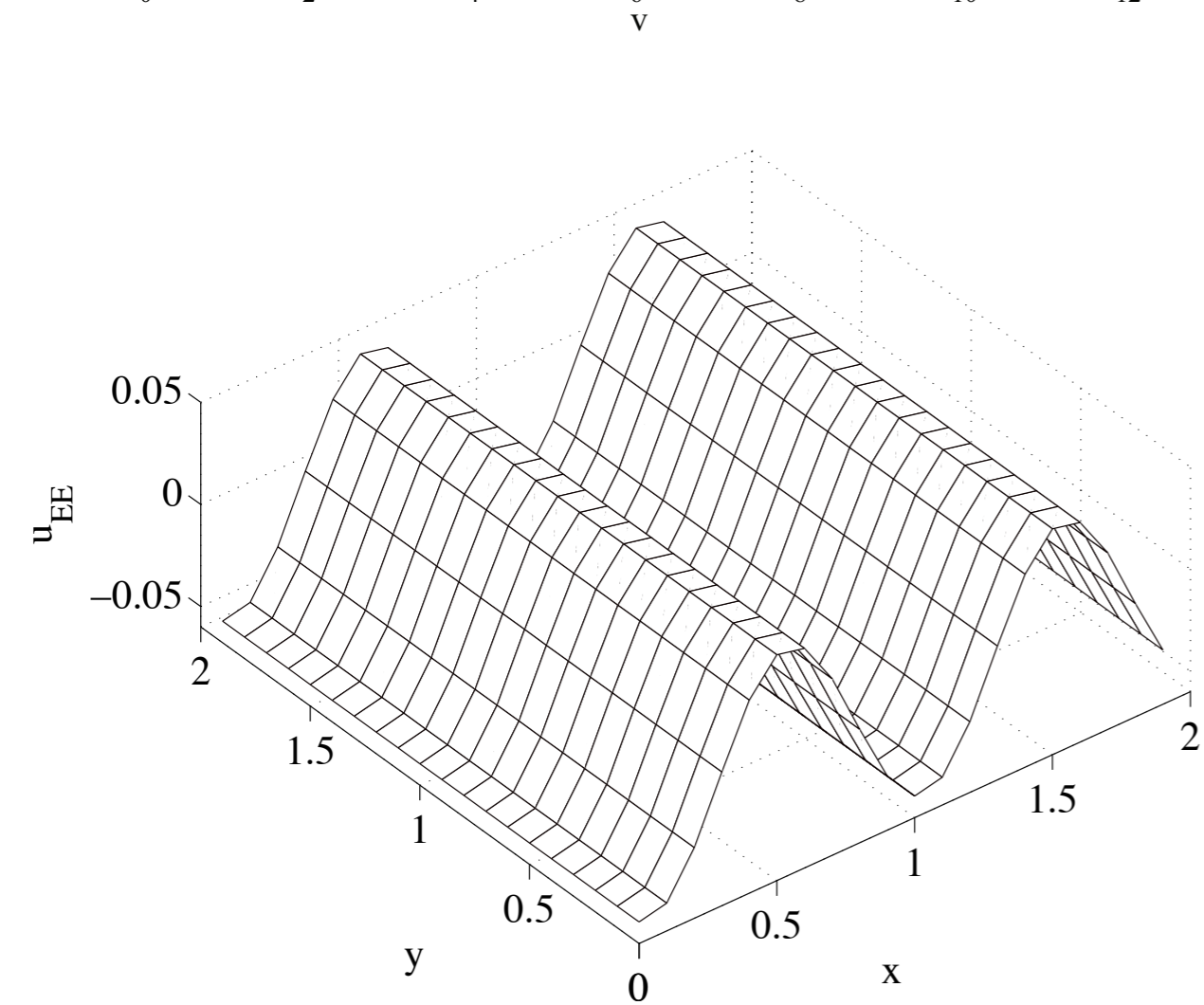
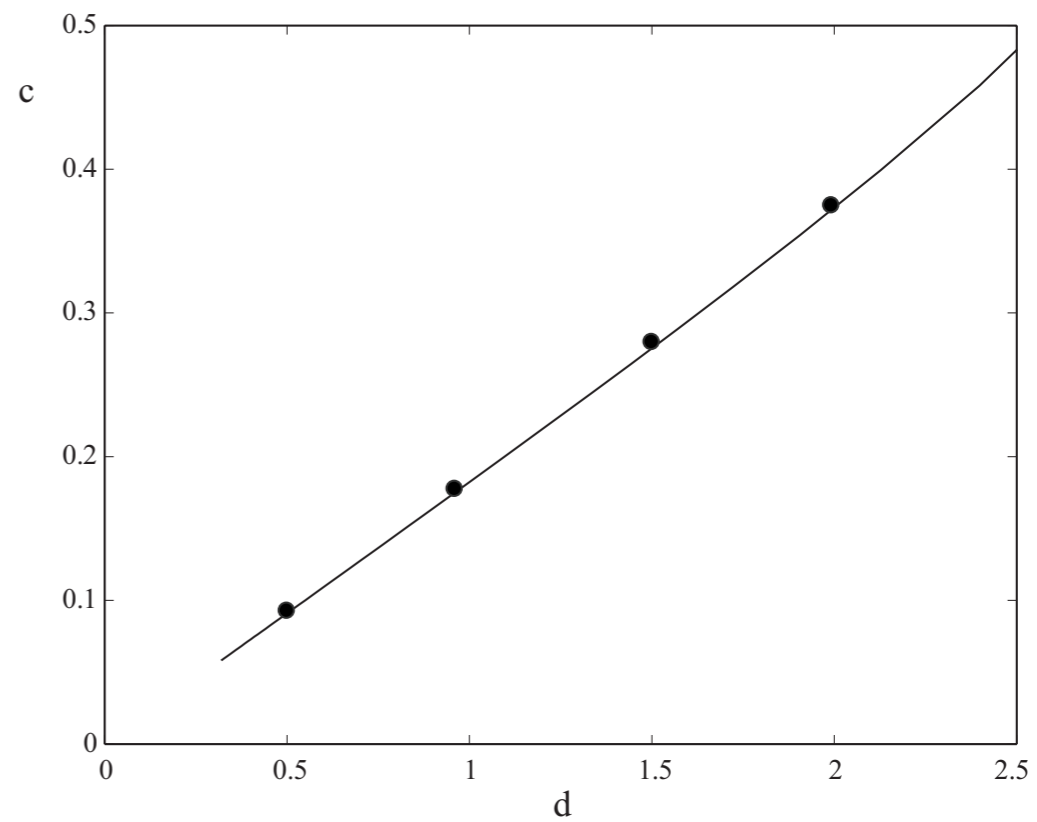
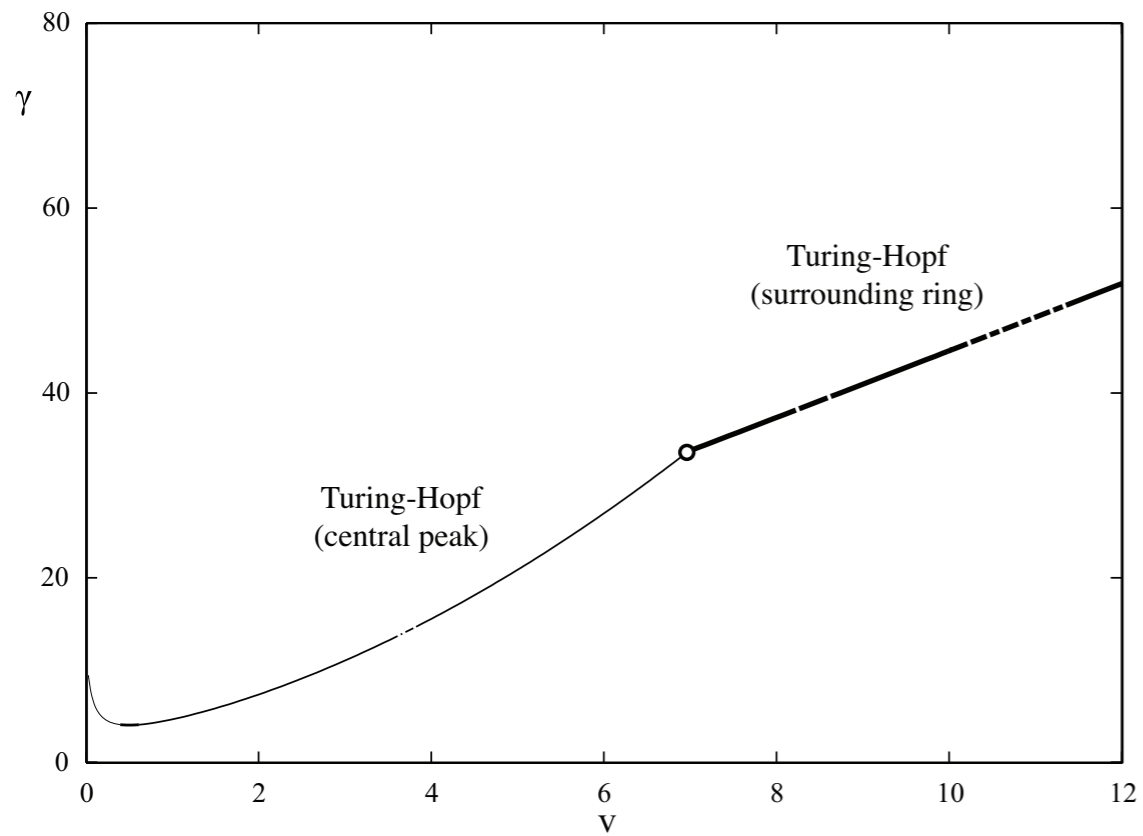
$$\mathbf{k}_1 = 2\pi/d(1, 0) \text{ and } \mathbf{k}_2 = 2\pi/d(0, 1)$$

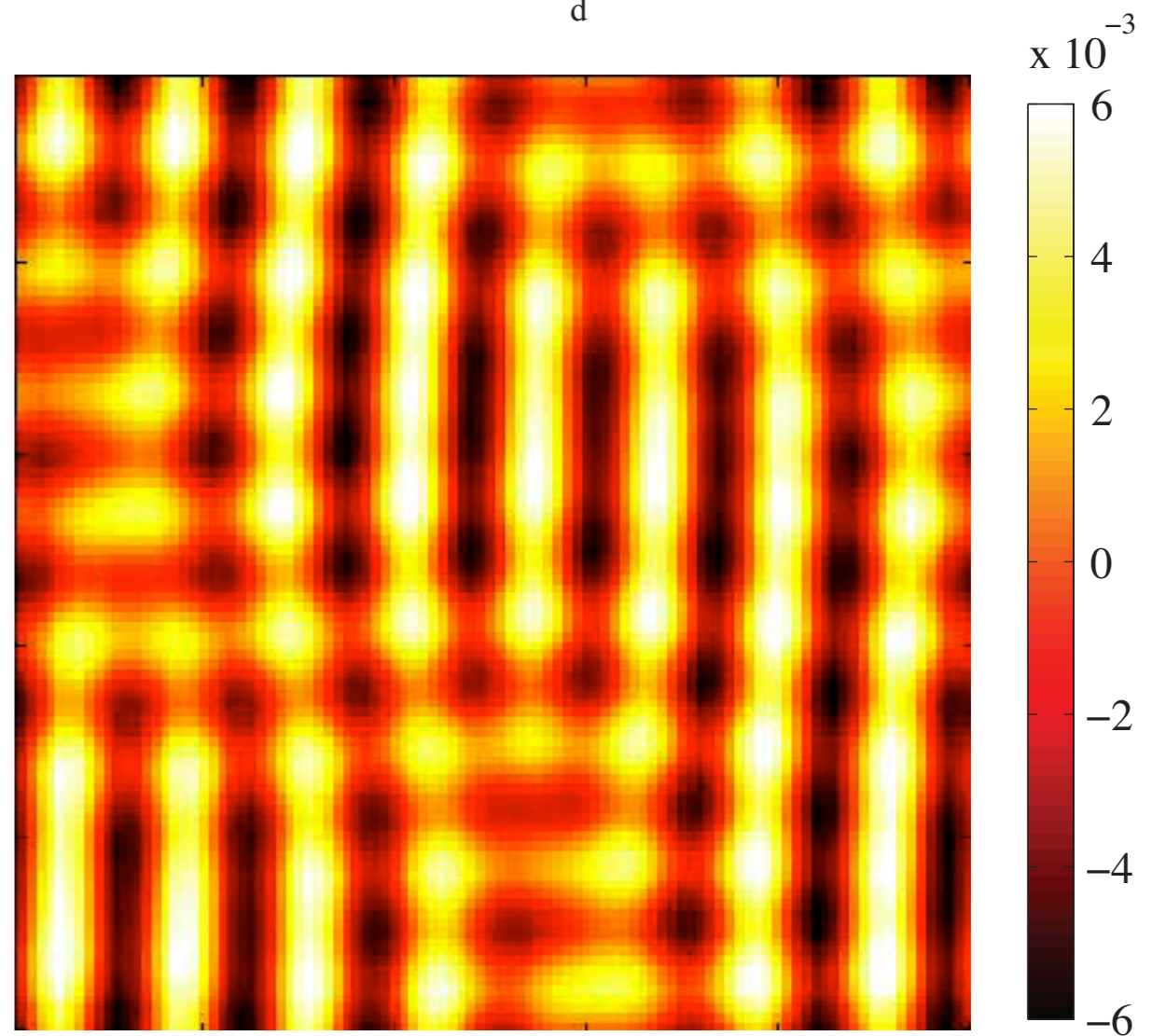
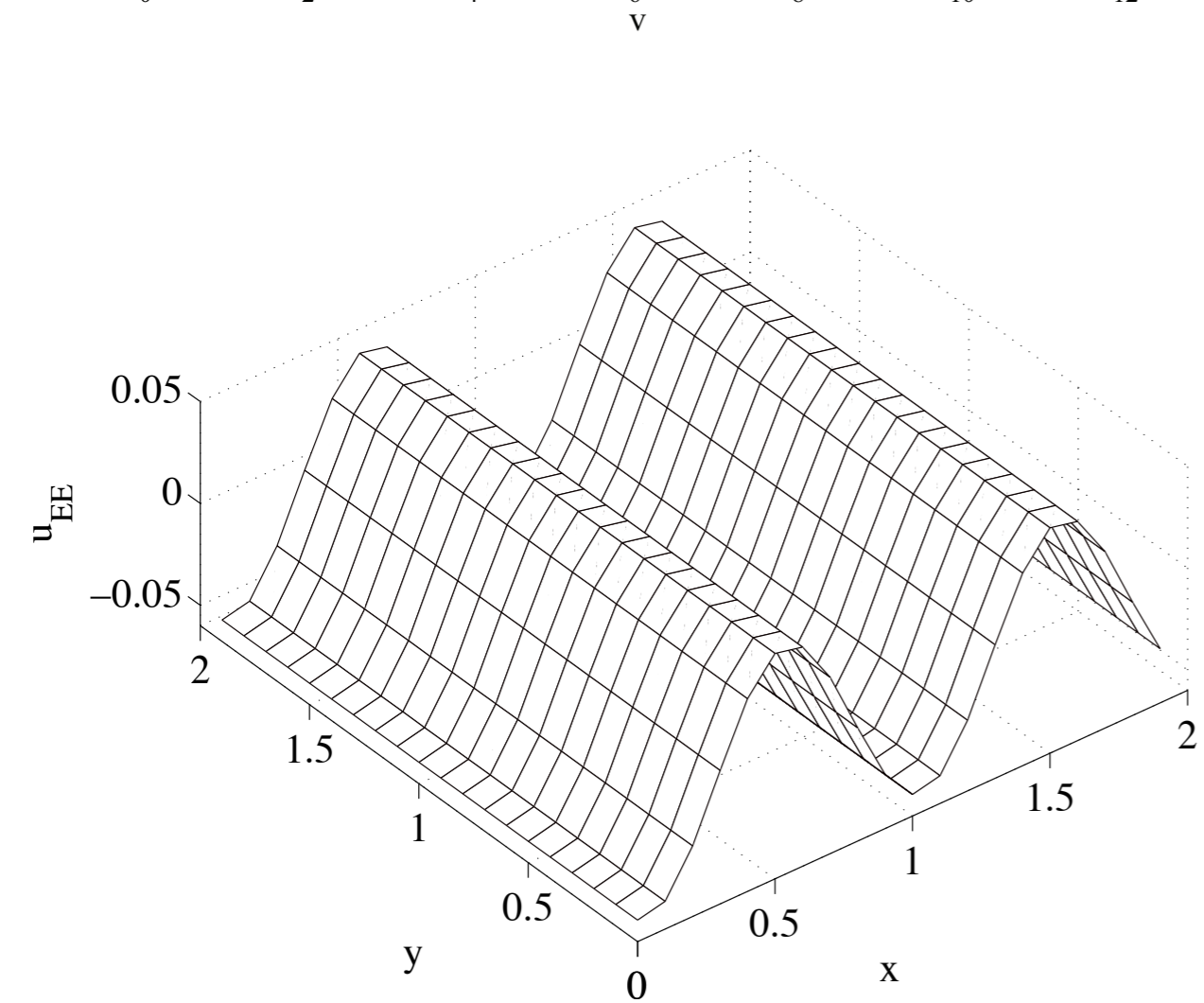
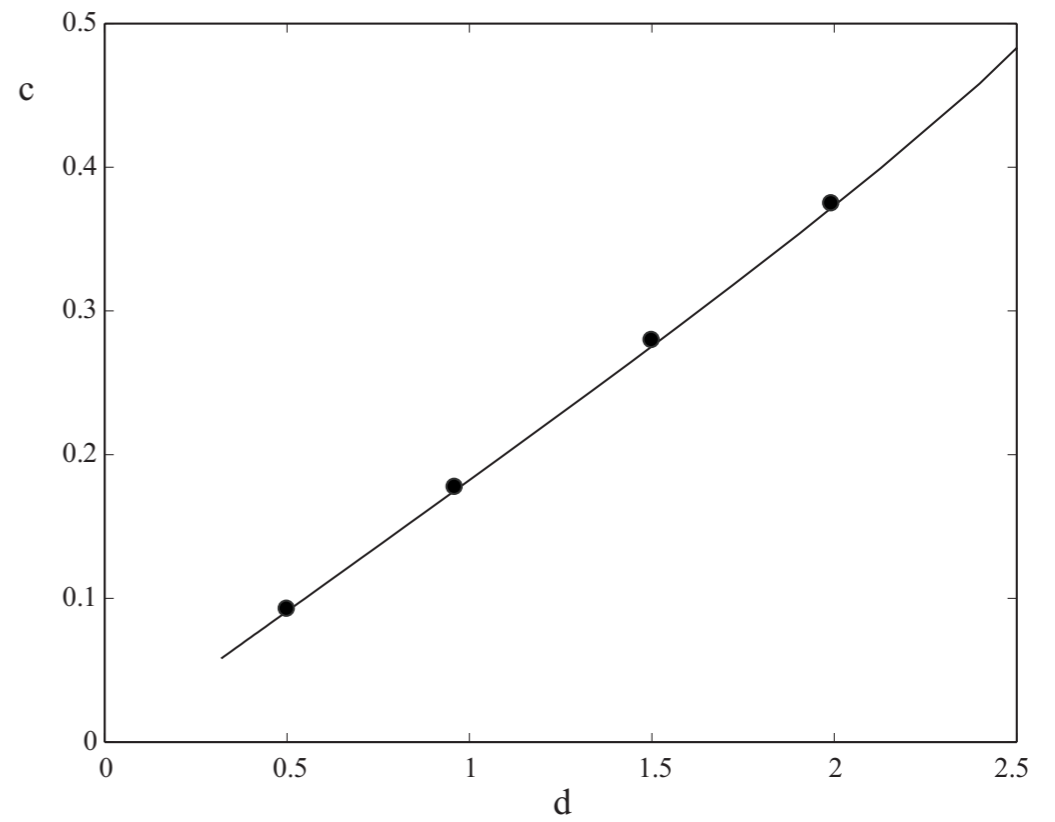
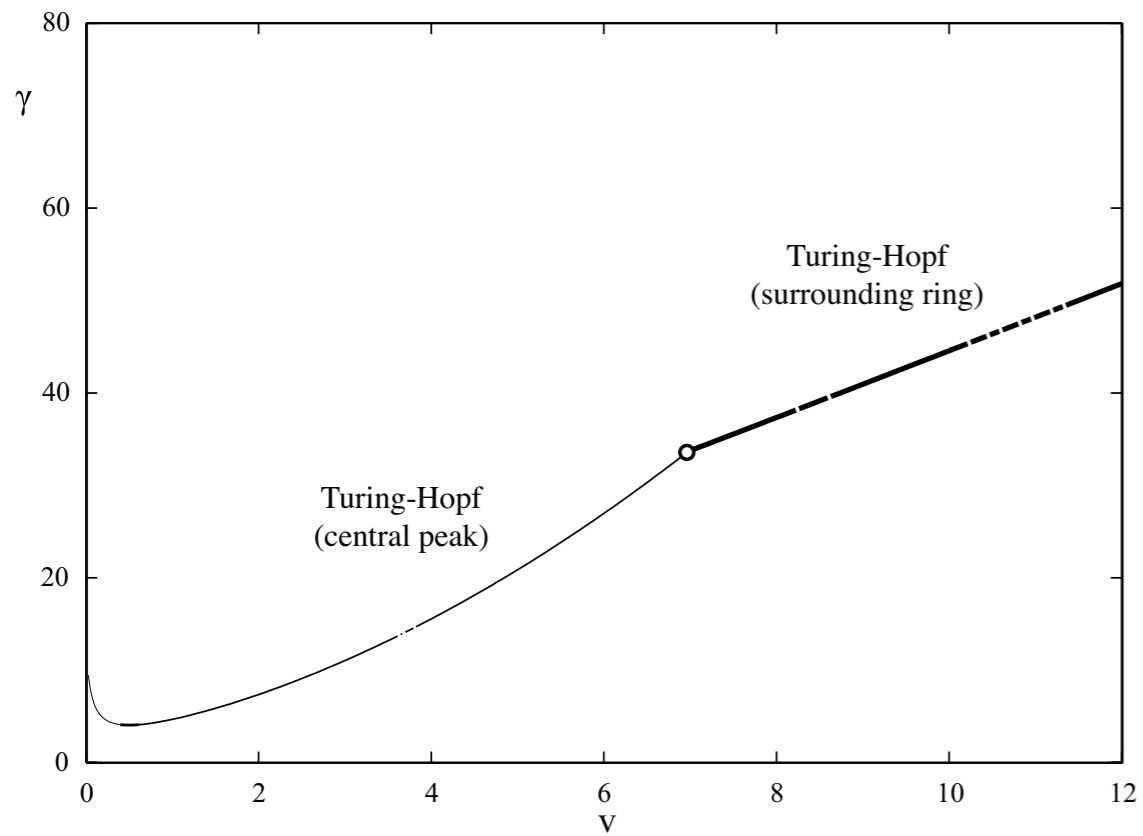


$$J_{ab}^{\mathbf{q}} = [\delta(\mathbf{q} - \mathbf{k}_1) + \delta(\mathbf{q} + \mathbf{k}_1) + \delta(\mathbf{q} - \mathbf{k}_2) + \delta(\mathbf{q} + \mathbf{k}_2)]/4$$

... and we need only consider the modes indexed by $\mathbf{k}_{1,2}$







Threshold Accommodation

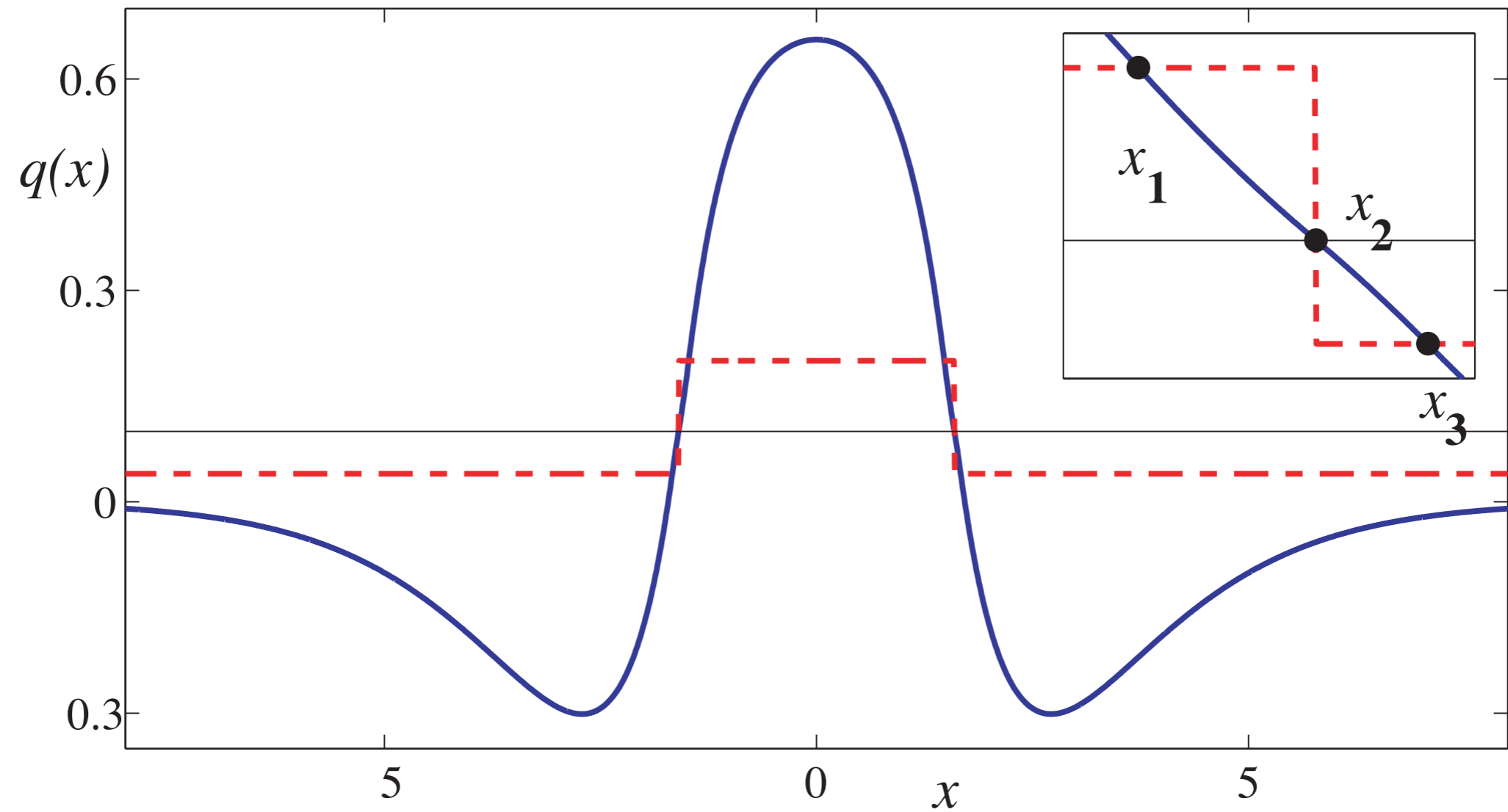
$$u(x, t) = \int_{-\infty}^{\infty} w(x-y) \int_{-\infty}^t \eta(t-s) f(u(y, s) - h) ds dy$$

Hill (1936), "... the threshold rises when the *local potential* is maintained ... and reverts gradually to its original value when the nerve is allowed to rest."

$$h_t = -(h - h_0) + \kappa H(u - \theta)$$

Time-independent solutions $(u, h) = (q(x), p(x))$

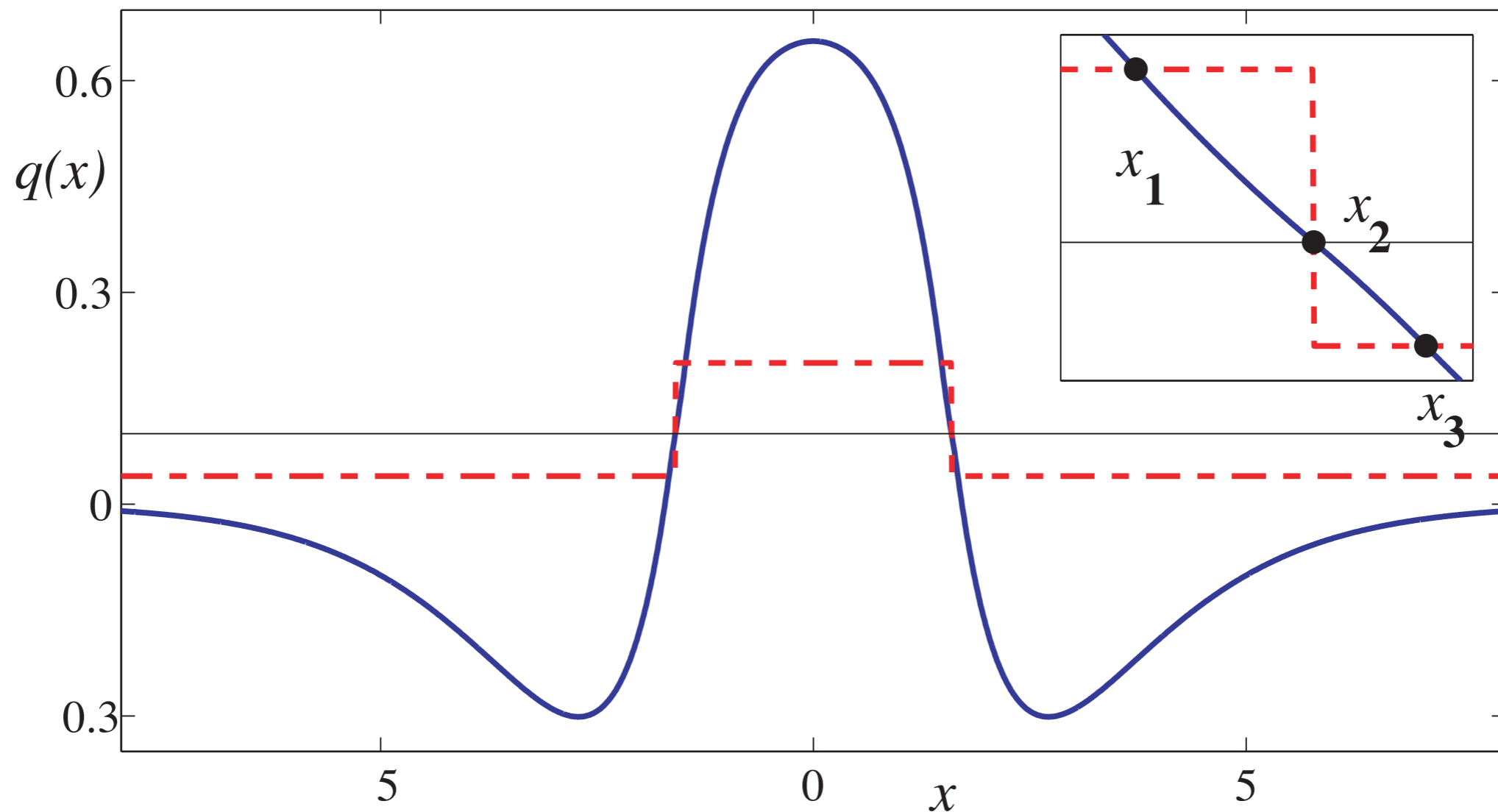
Time-independent solutions $(u, h) = (q(x), p(x))$



Time-independent solutions $(u, h) = (q(x), p(x))$

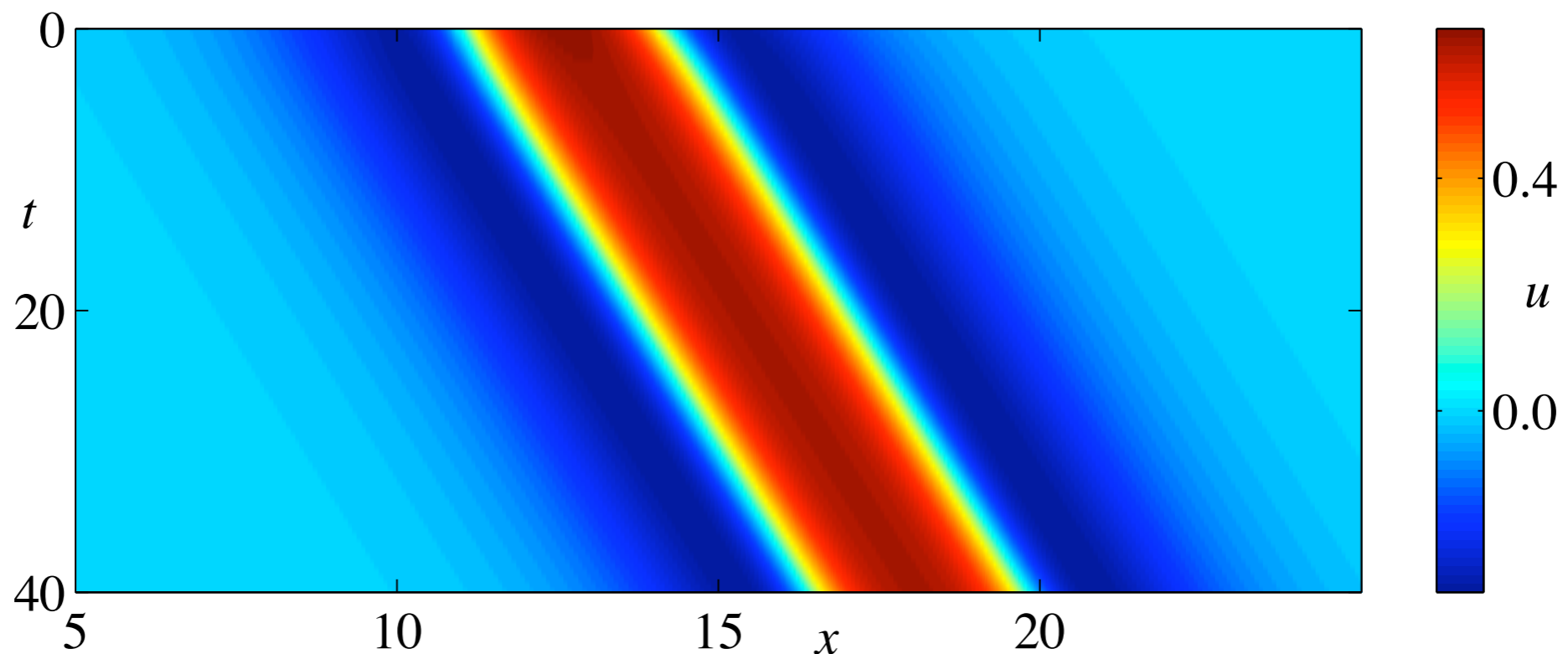
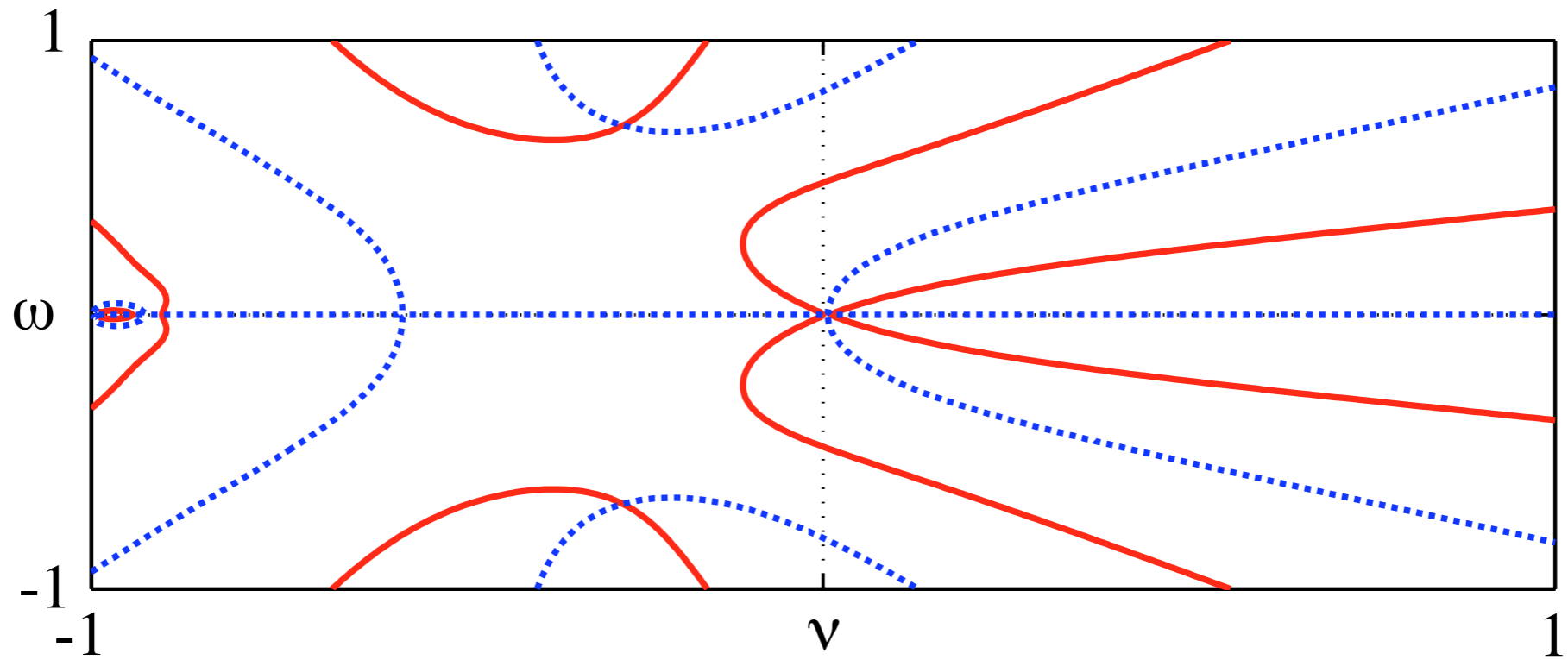
$$(w \otimes f)(x, t) = \int_{-\infty}^{\infty} w(y) f(x - y, t) dy$$

$$q = w \otimes H(q - p), \quad p = \begin{cases} h_0 & q < \theta \\ h_0 + \kappa & q \geq \theta \end{cases}$$



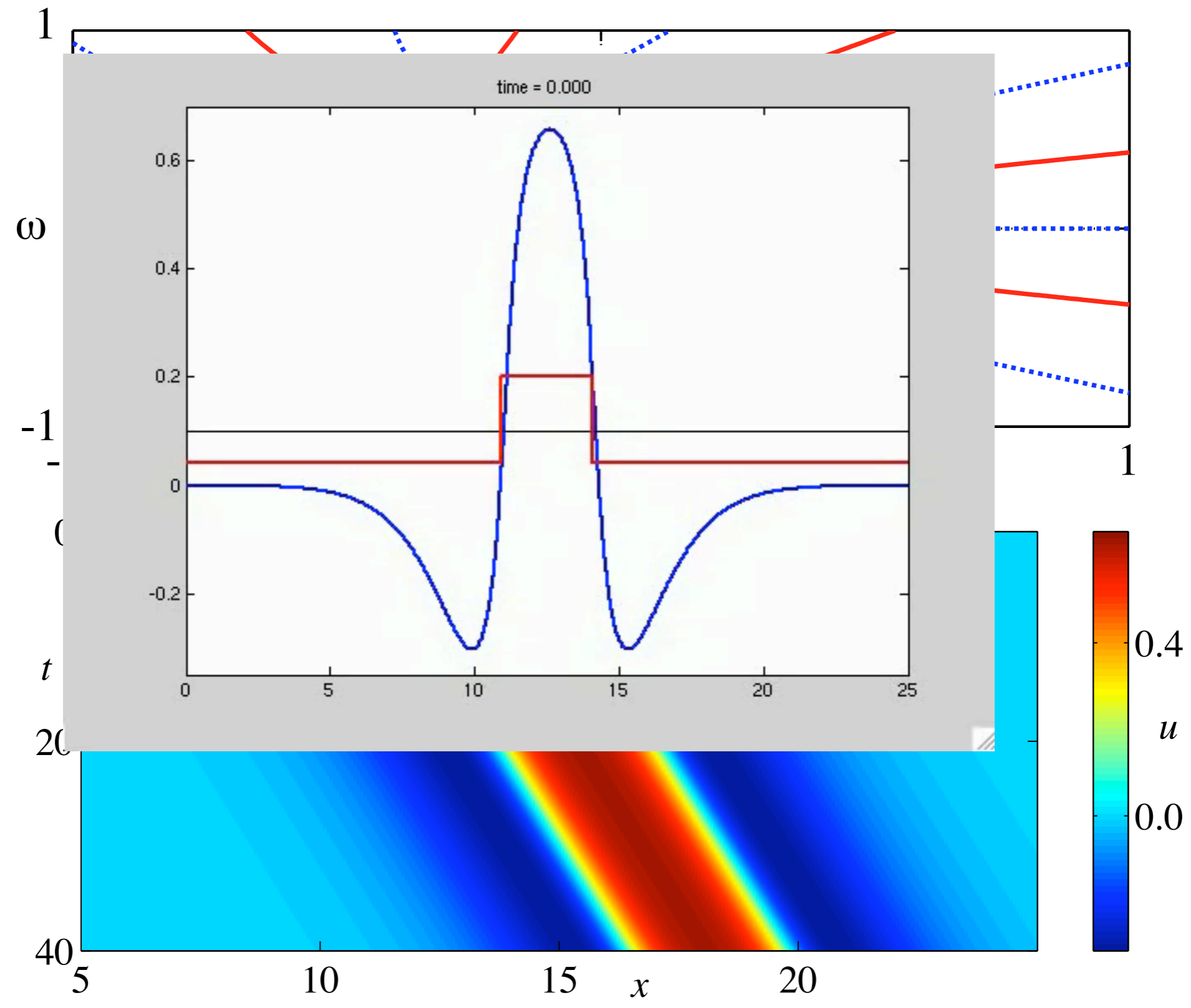
Bump Stability III: $\eta(t) = \alpha e^{-\alpha t}$

Low κ instability on Re axis (increasing α)



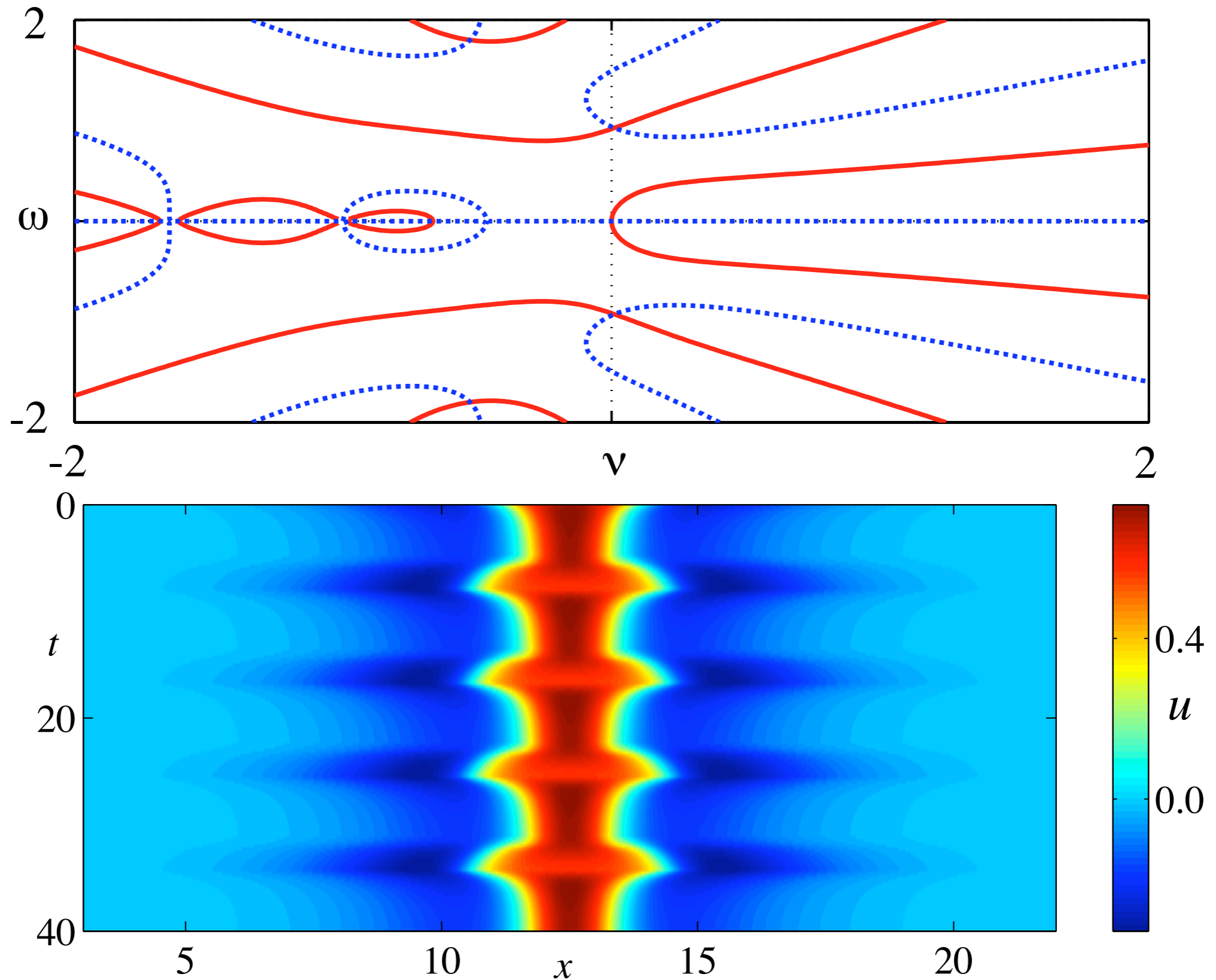
Bump Stability III: $\eta(t) = \alpha e^{-\alpha t}$

Low κ instability on Re axis (increasing α)



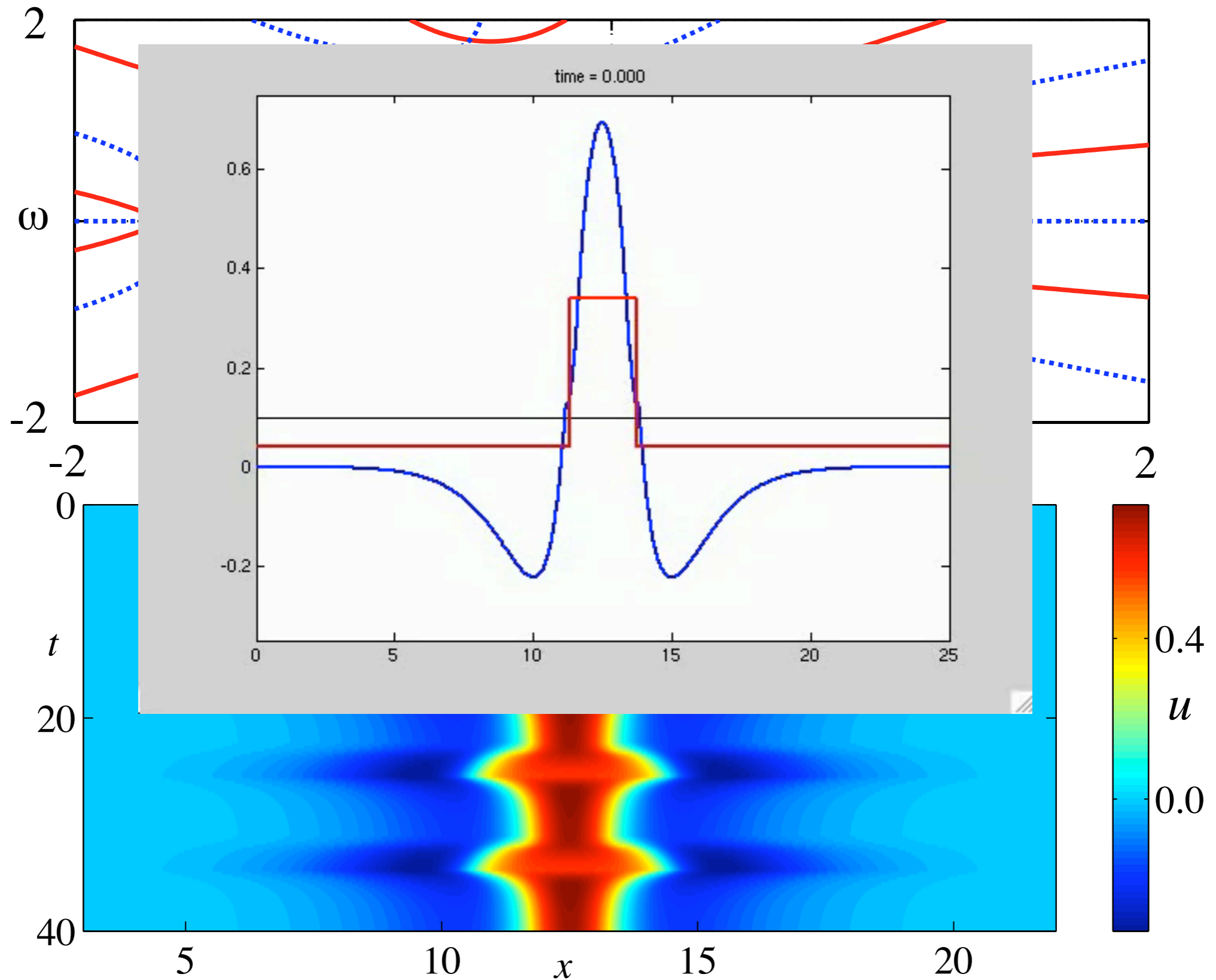
Bump Stability IV

High κ instability on Im axis (increasing α) gives a breather

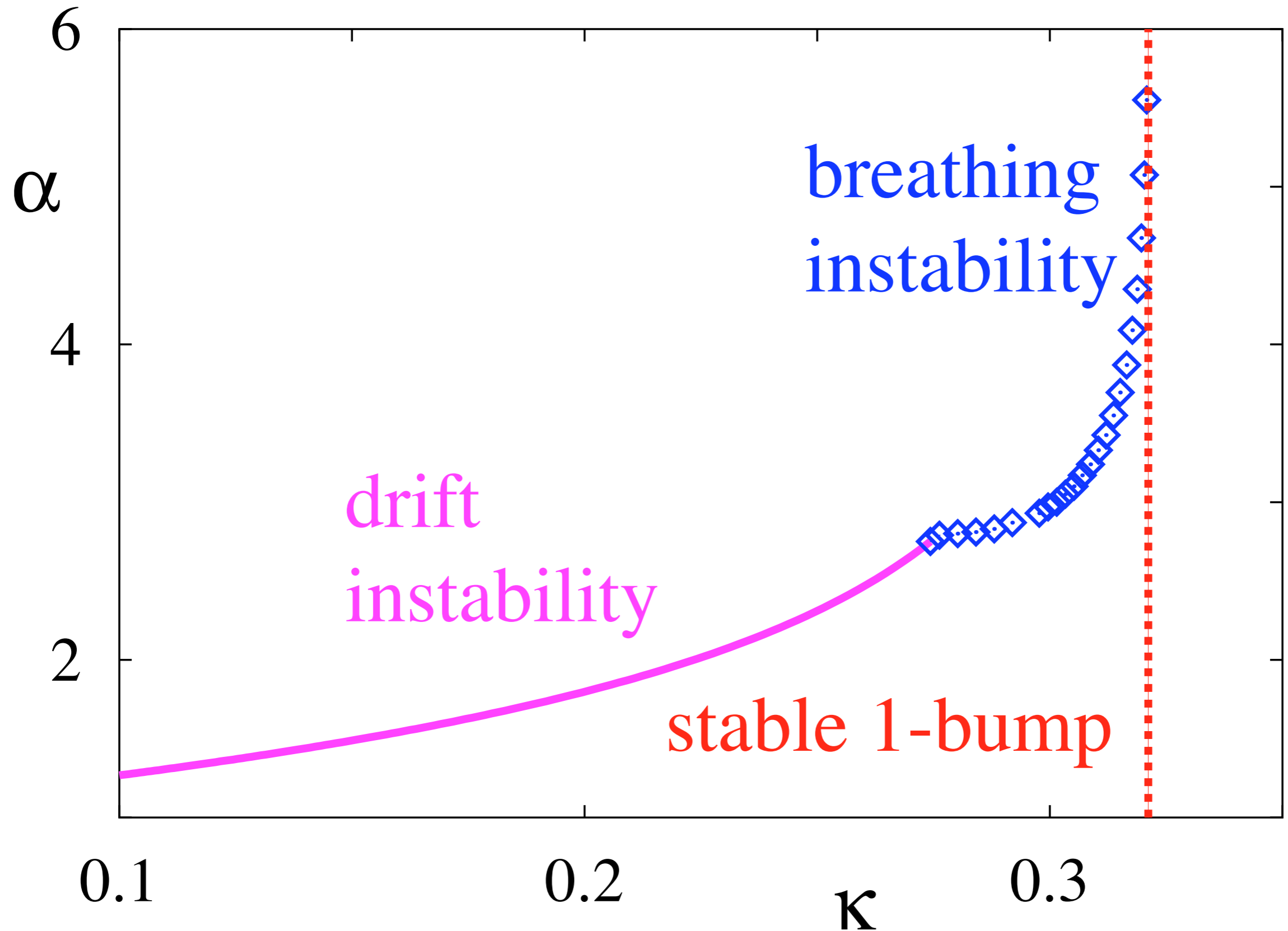


Bump Stability IV

High κ instability on Im axis (increasing α) gives a breather

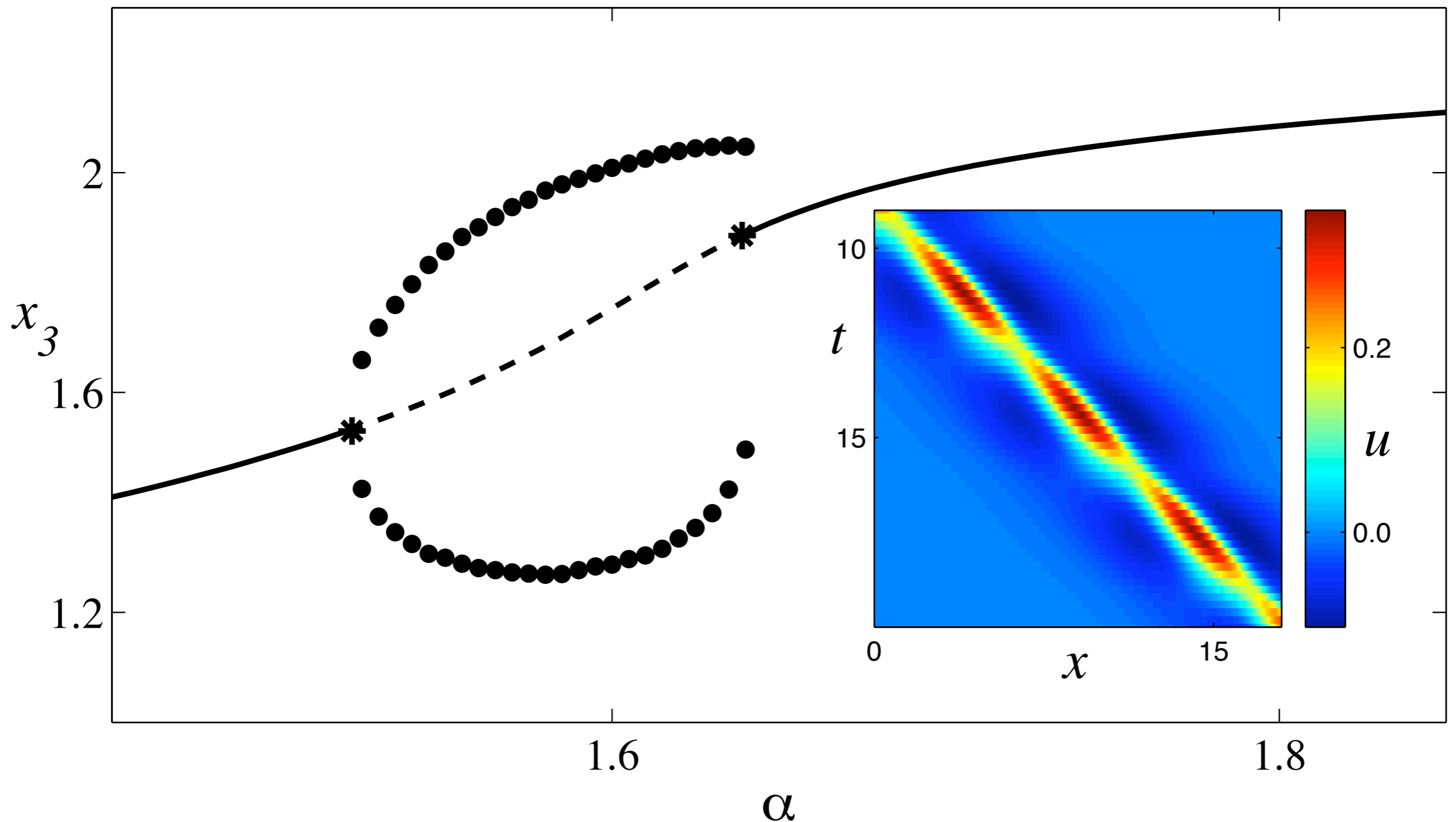


Summary of Bump instabilities



Dynamic instability of pulses

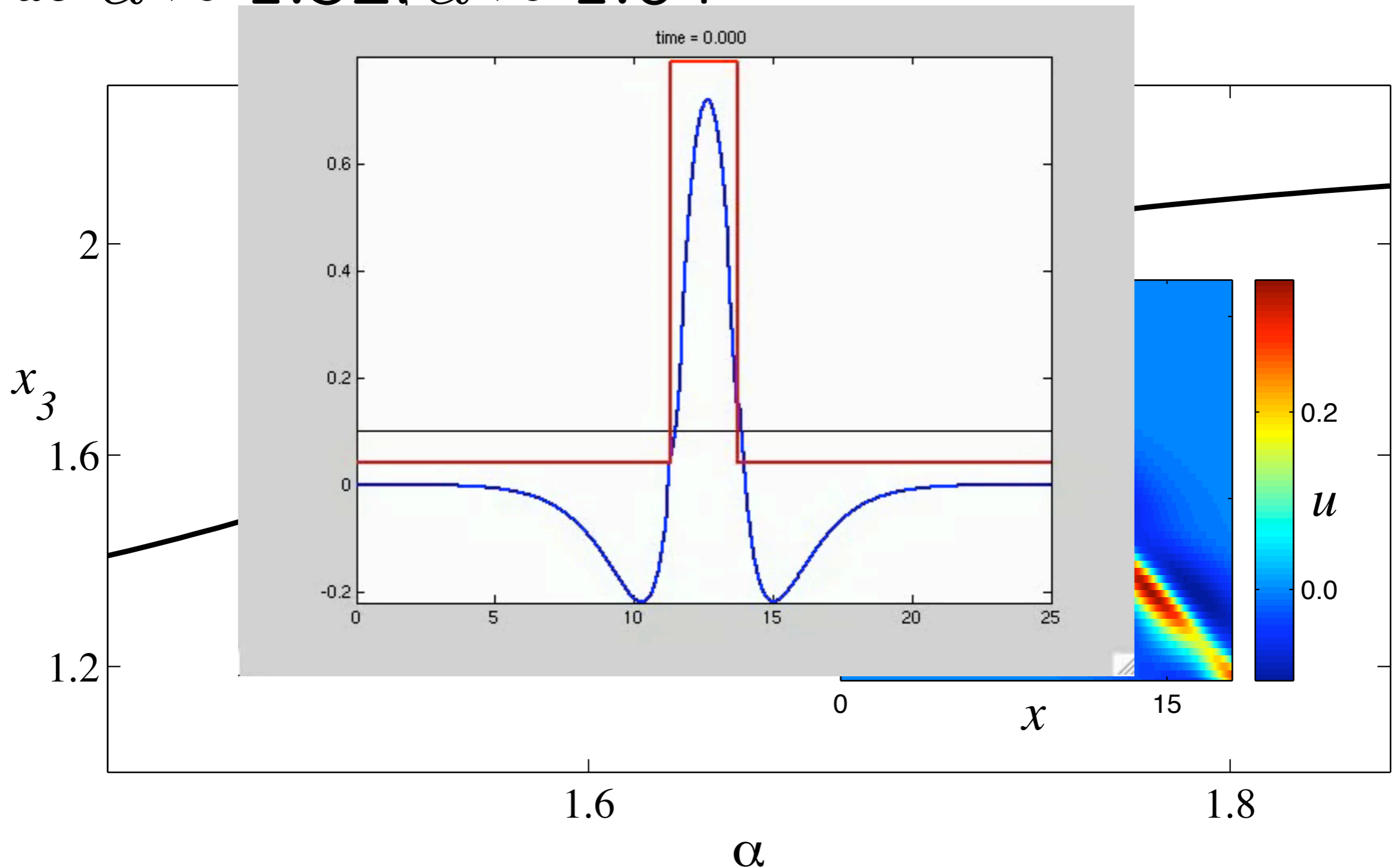
A pair of complex conjugate eigenvalues crosses the Im axis at $\alpha \approx 1.52, \alpha \approx 1.64$



Exact solution (curve); Hopf bifurcations (asterisks); Numerics (points).

Dynamic instability of pulses

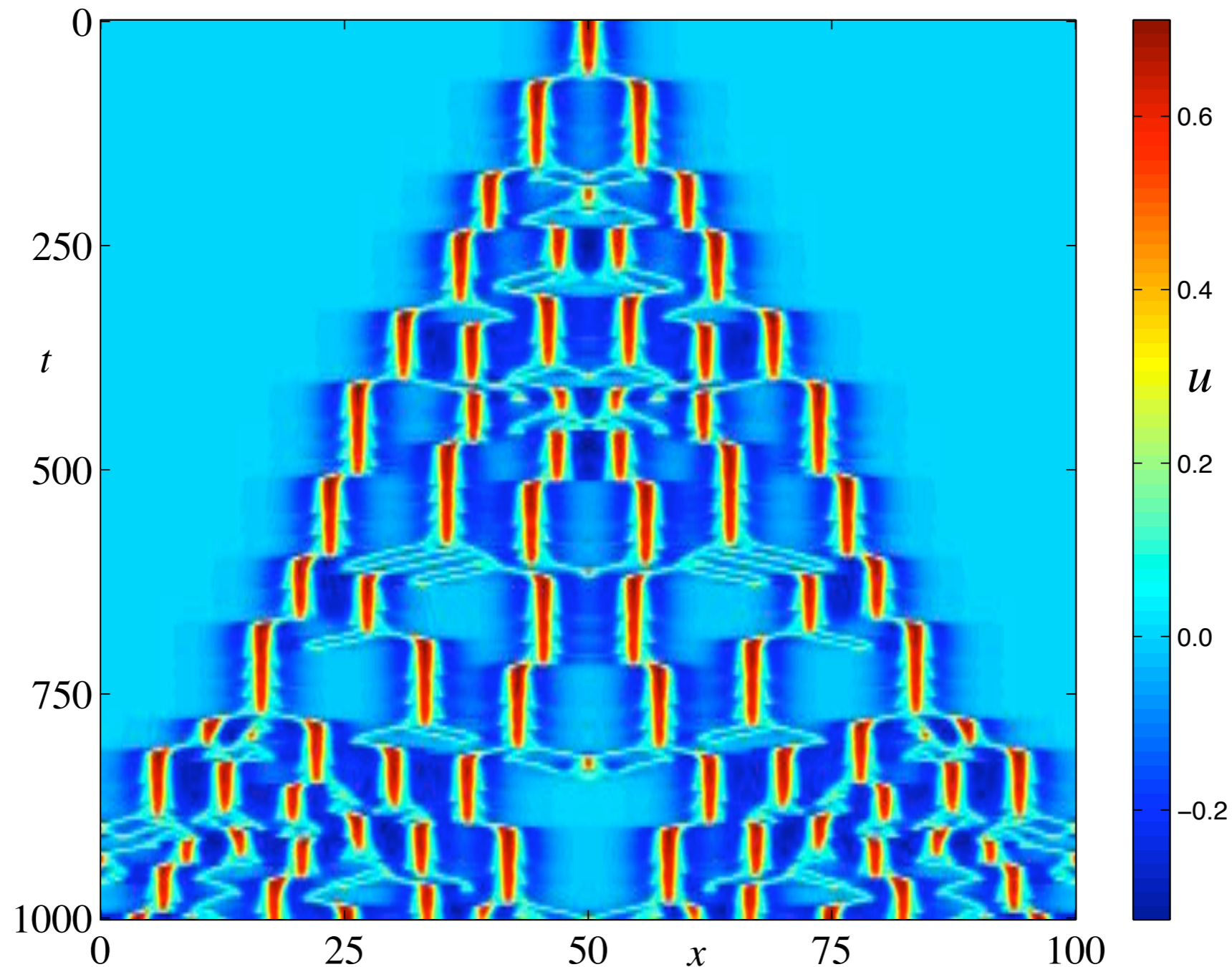
A pair of complex conjugate eigenvalues crosses the Im axis at $\alpha \approx 1.52, \alpha \approx 1.64$



Exact solution (curve); Hopf bifurcations (asterisks); Numerics (points).

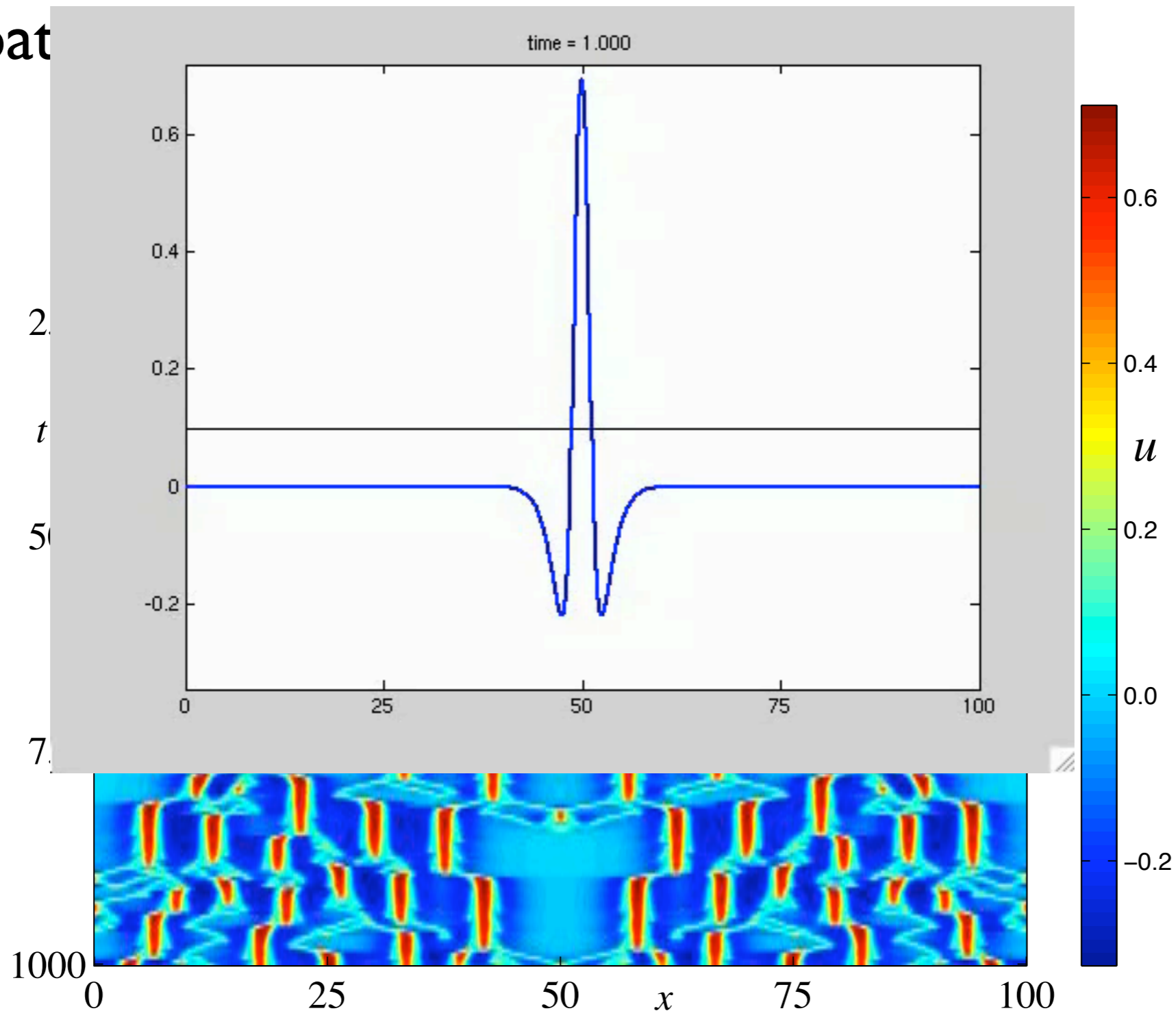
Exotic Dynamics

... including asymmetric breathers, multiple bumps, multiple pulses, periodic traveling waves, and bump-splitting instabilities that appear to lead to spatio-temporal chaos.



Exotic Dynamics

... including asymmetric breathers, multiple bumps, multiple pulses, periodic traveling waves, and bump-splitting instabilities that appear to lead to spat



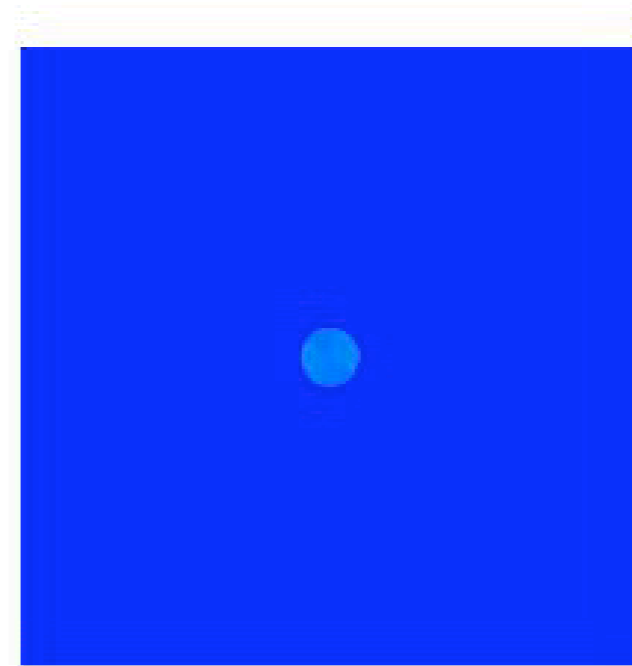
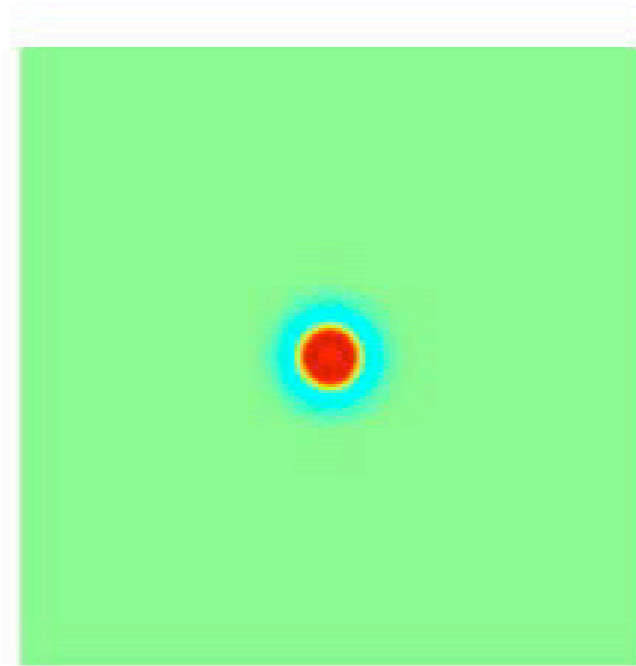
2D: Dimple Bumps

Complex splitting

2D: Dimple Bumps

Complex splitting

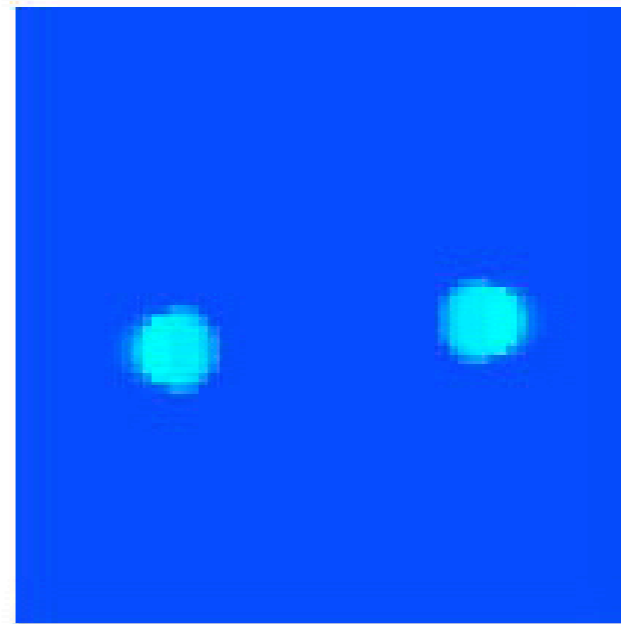
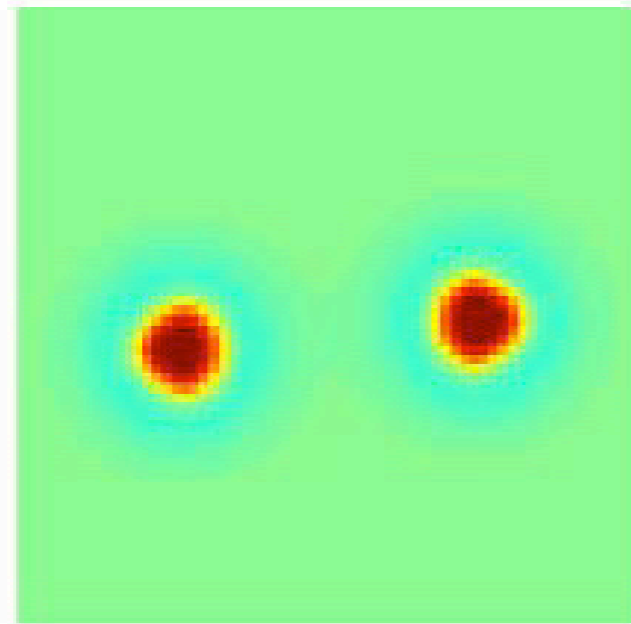
time = 2.000



2D: Collisions

2D: Collisions

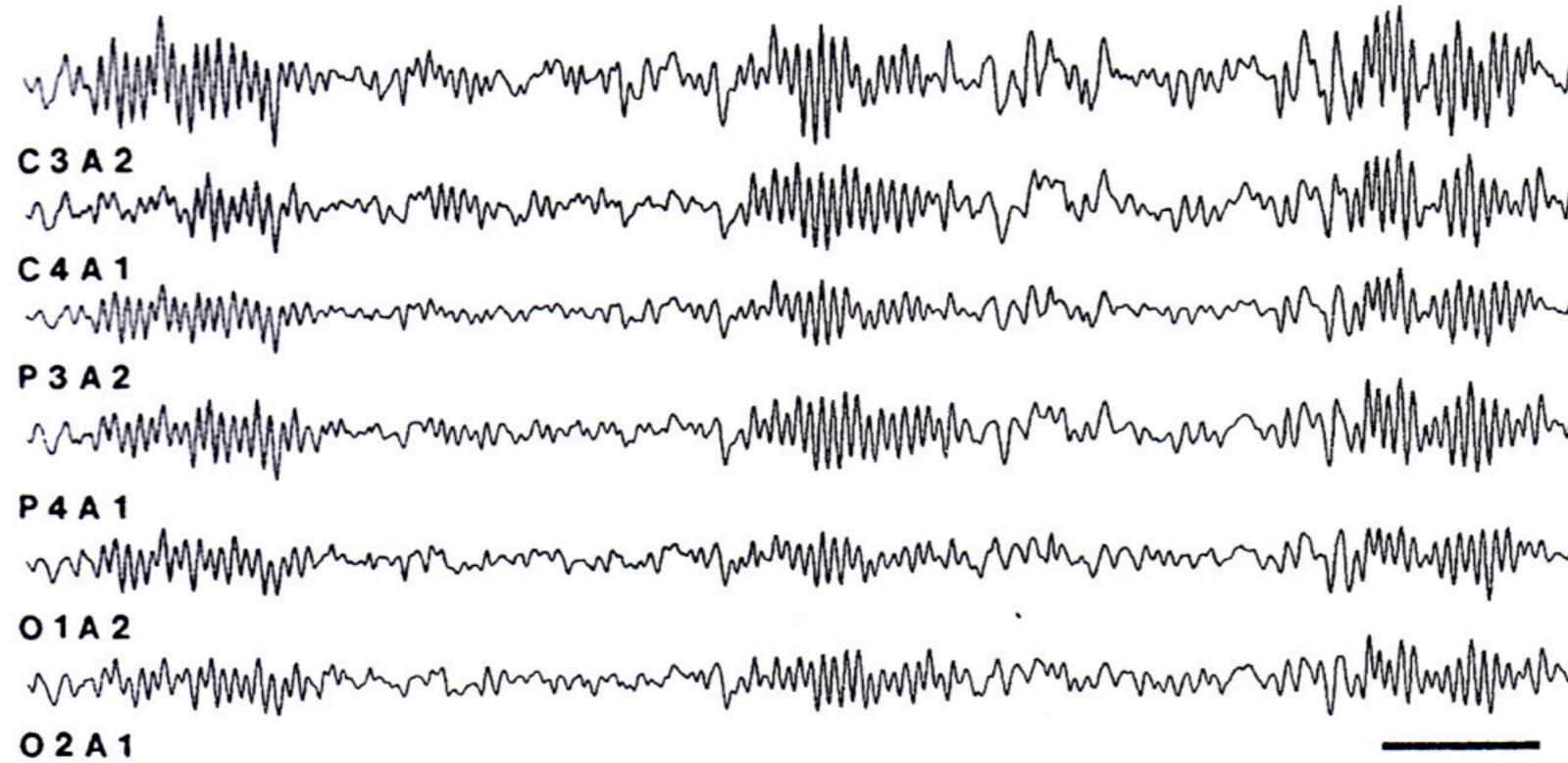
time = 1.000



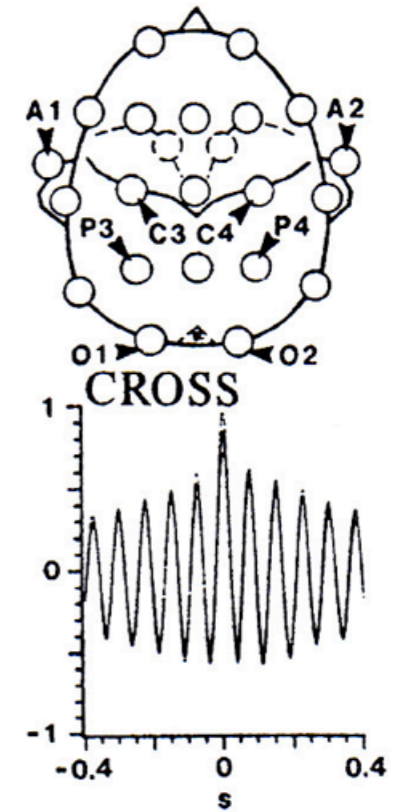
Rhythms and slow ionic currents

NATURAL SLEEP

HUMAN



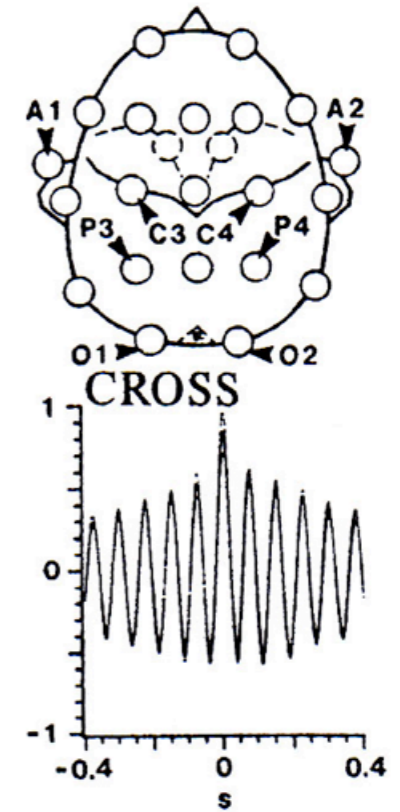
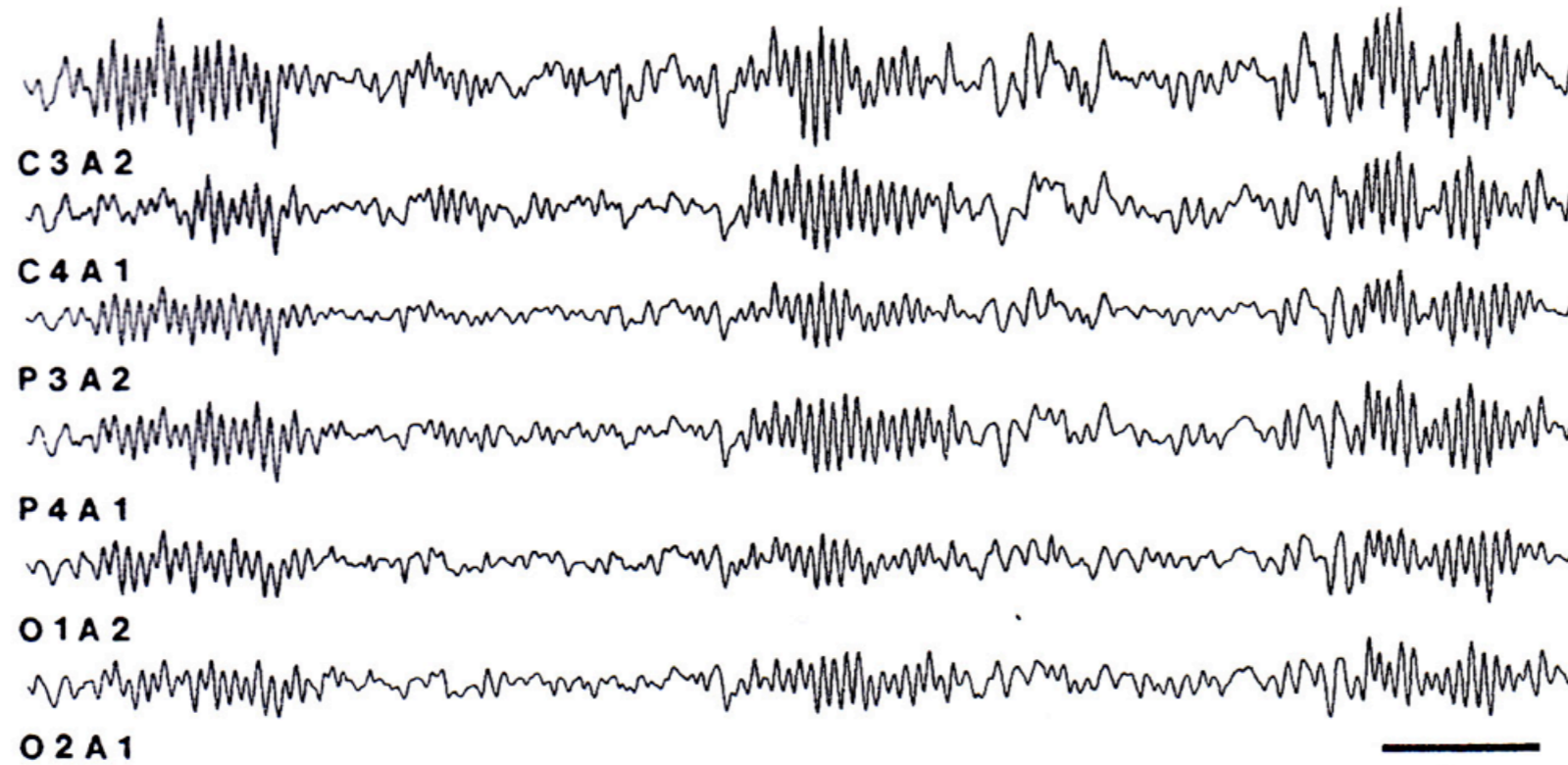
Synchronised (14 Hz) oscillations at onset of sleep



Rhythms and slow ionic currents

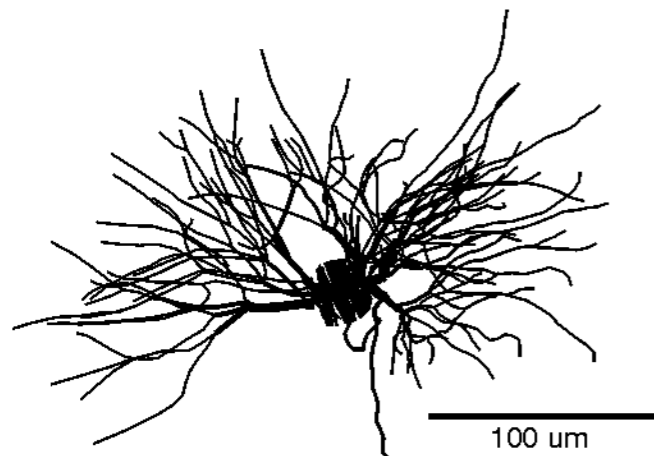
NATURAL SLEEP

HUMAN

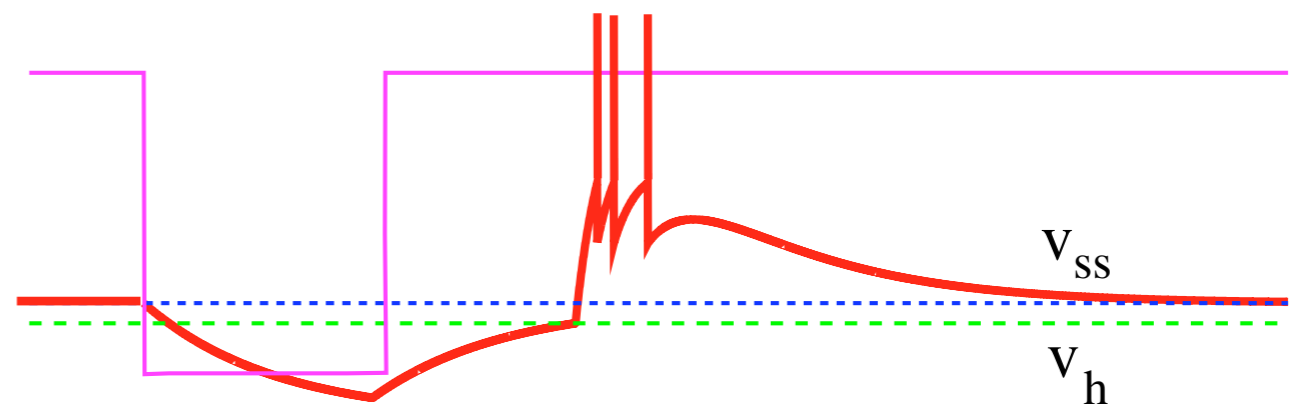


Synchronised (14 Hz) oscillations at onset of sleep

Thalamocortical (TC)



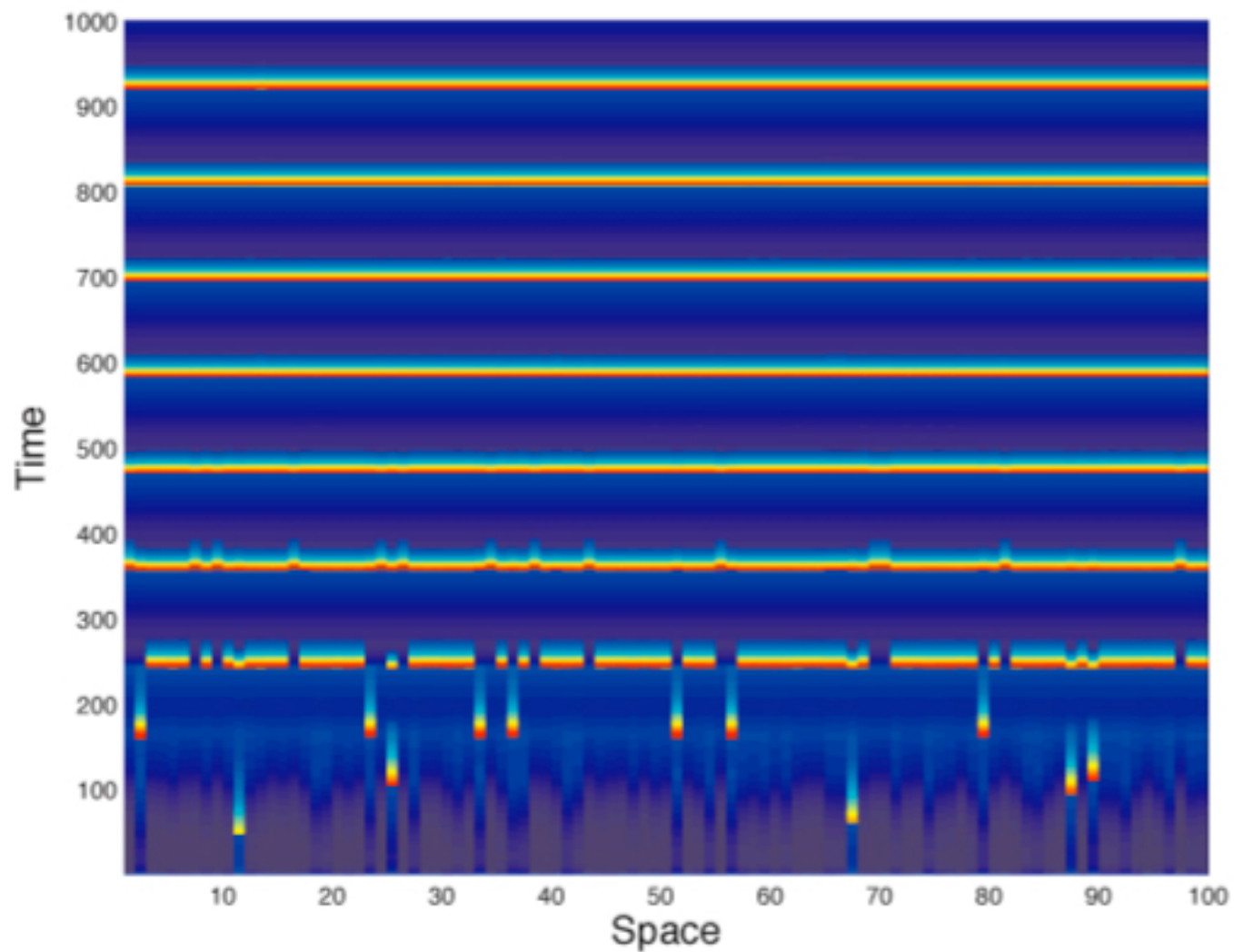
Post-Inhibitory Rebound



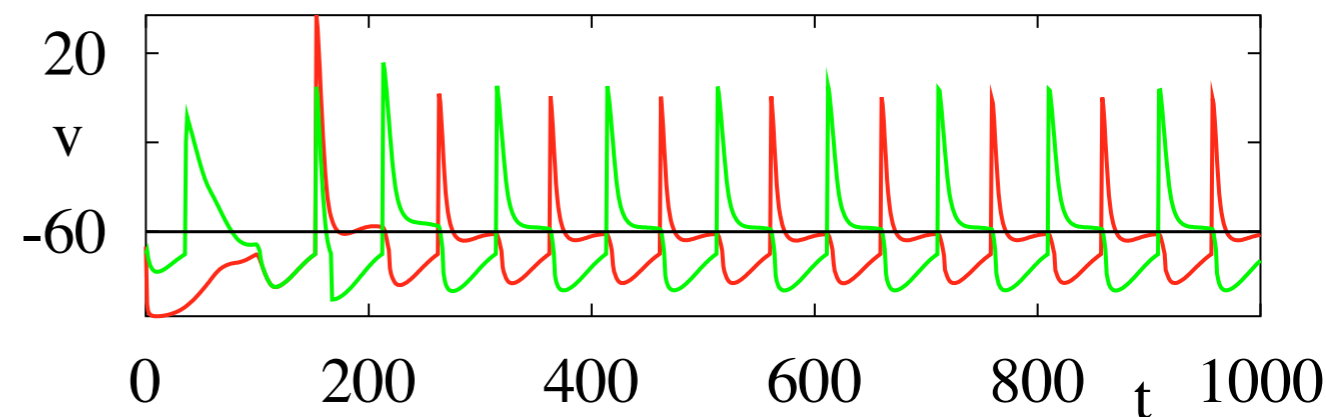
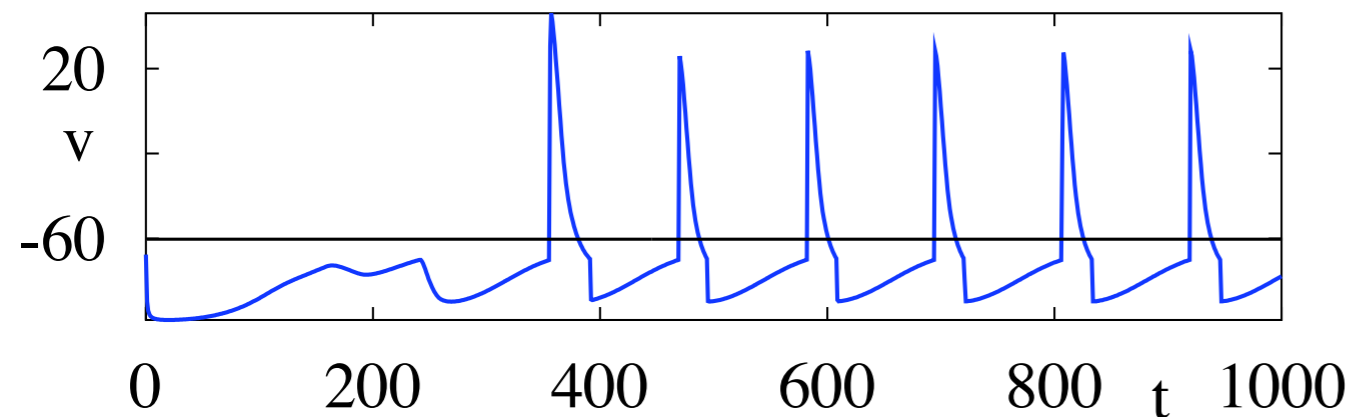
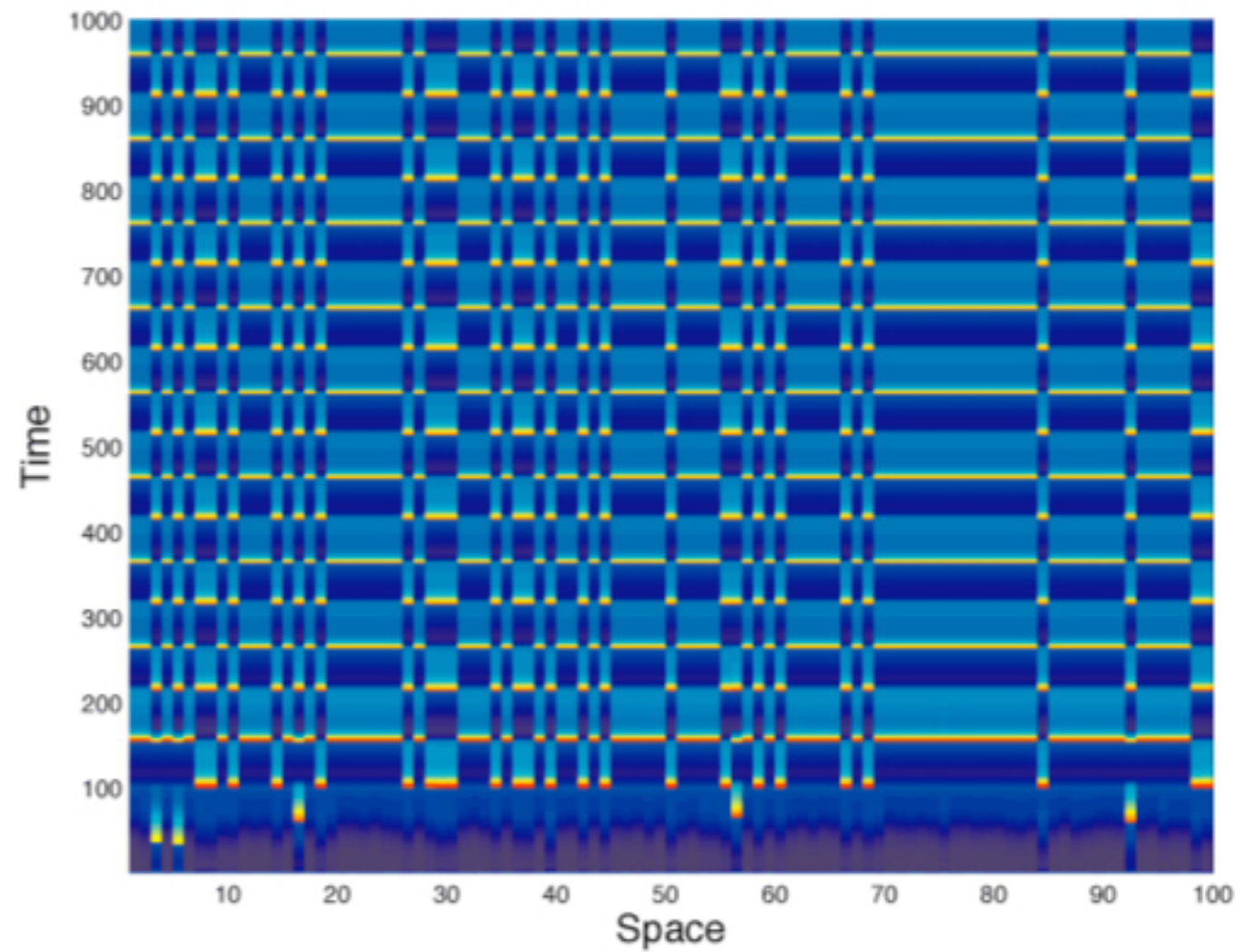
Dynamic response due to a low-threshold Calcium conductance.

Global inhibitory TC Network

Synchrony for slow synapses

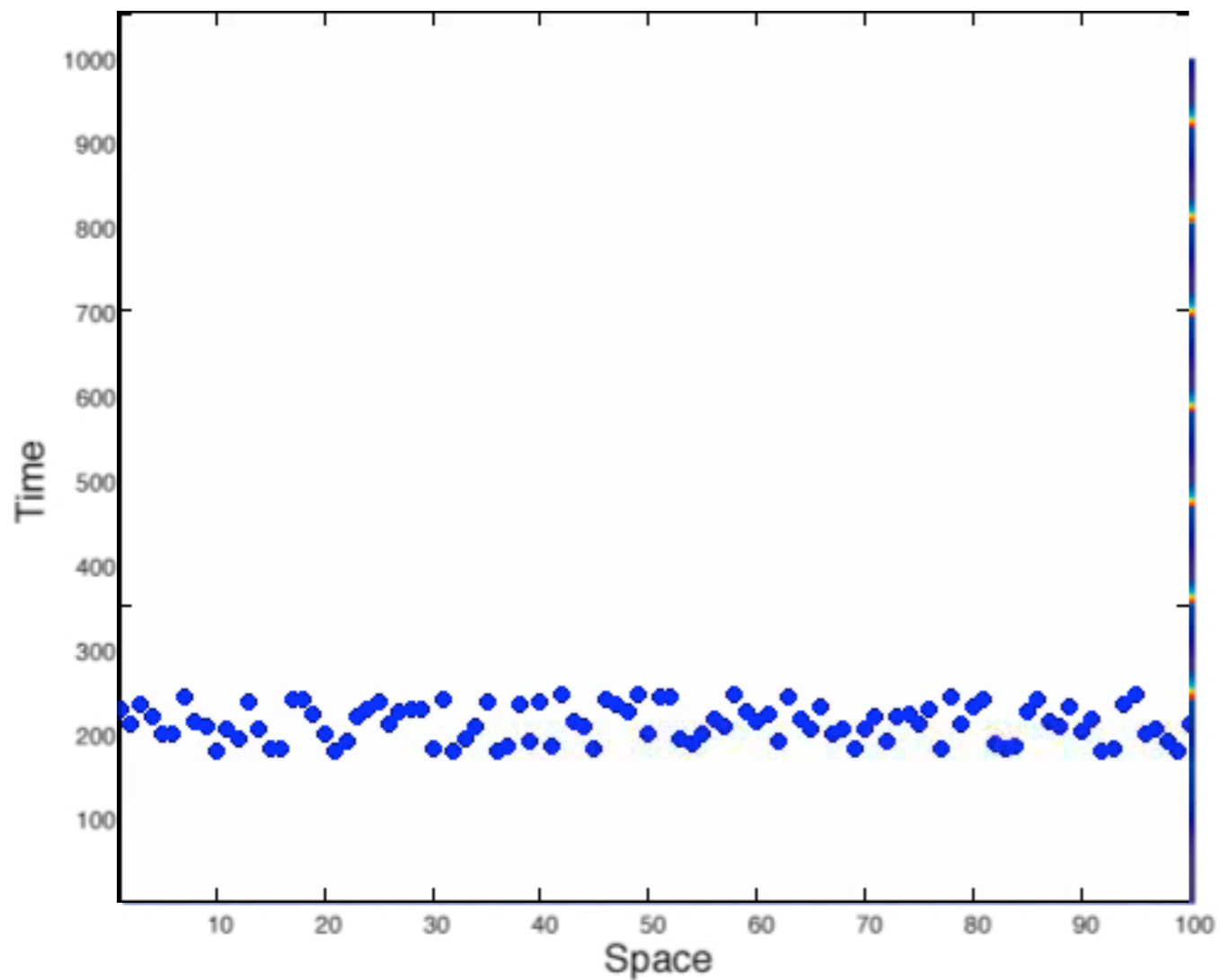


Clustering for faster synapses

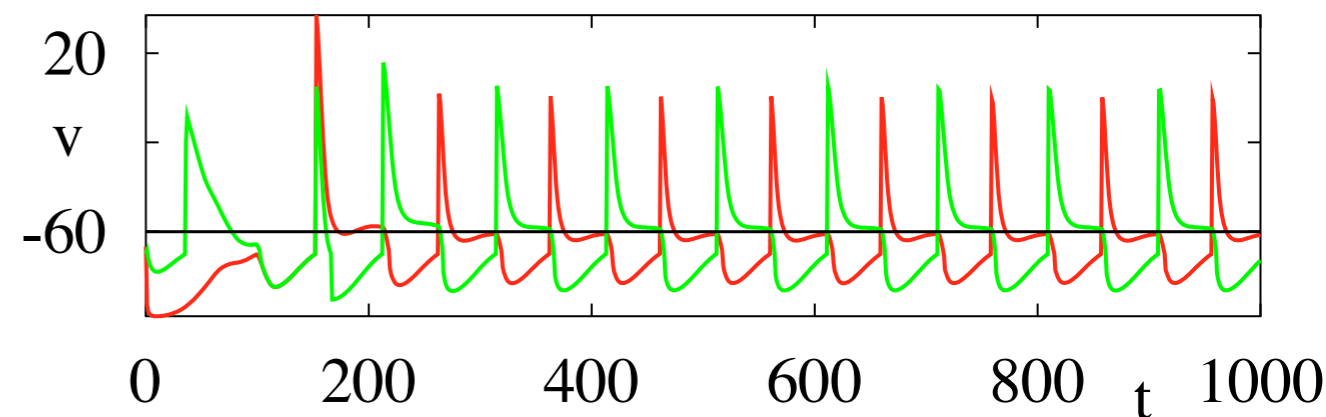
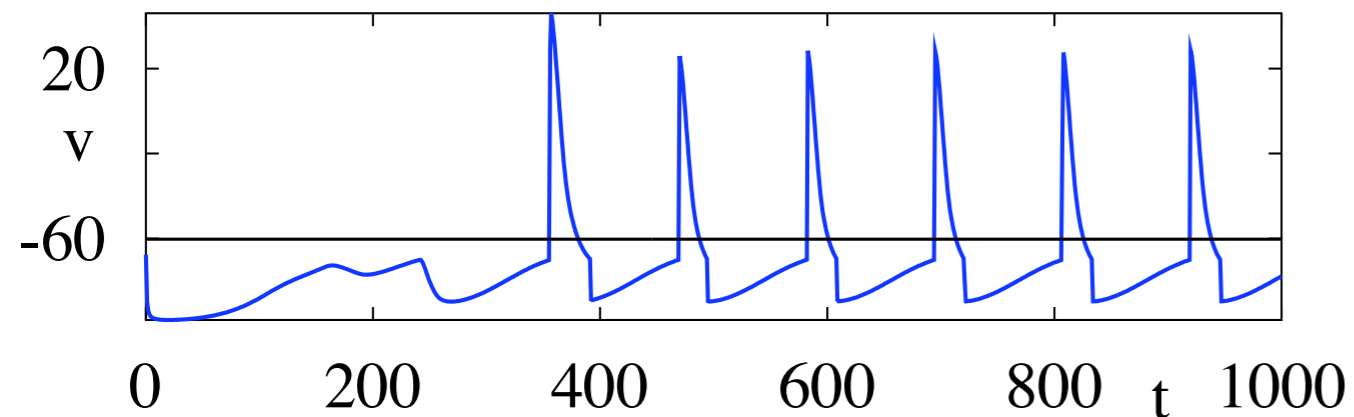
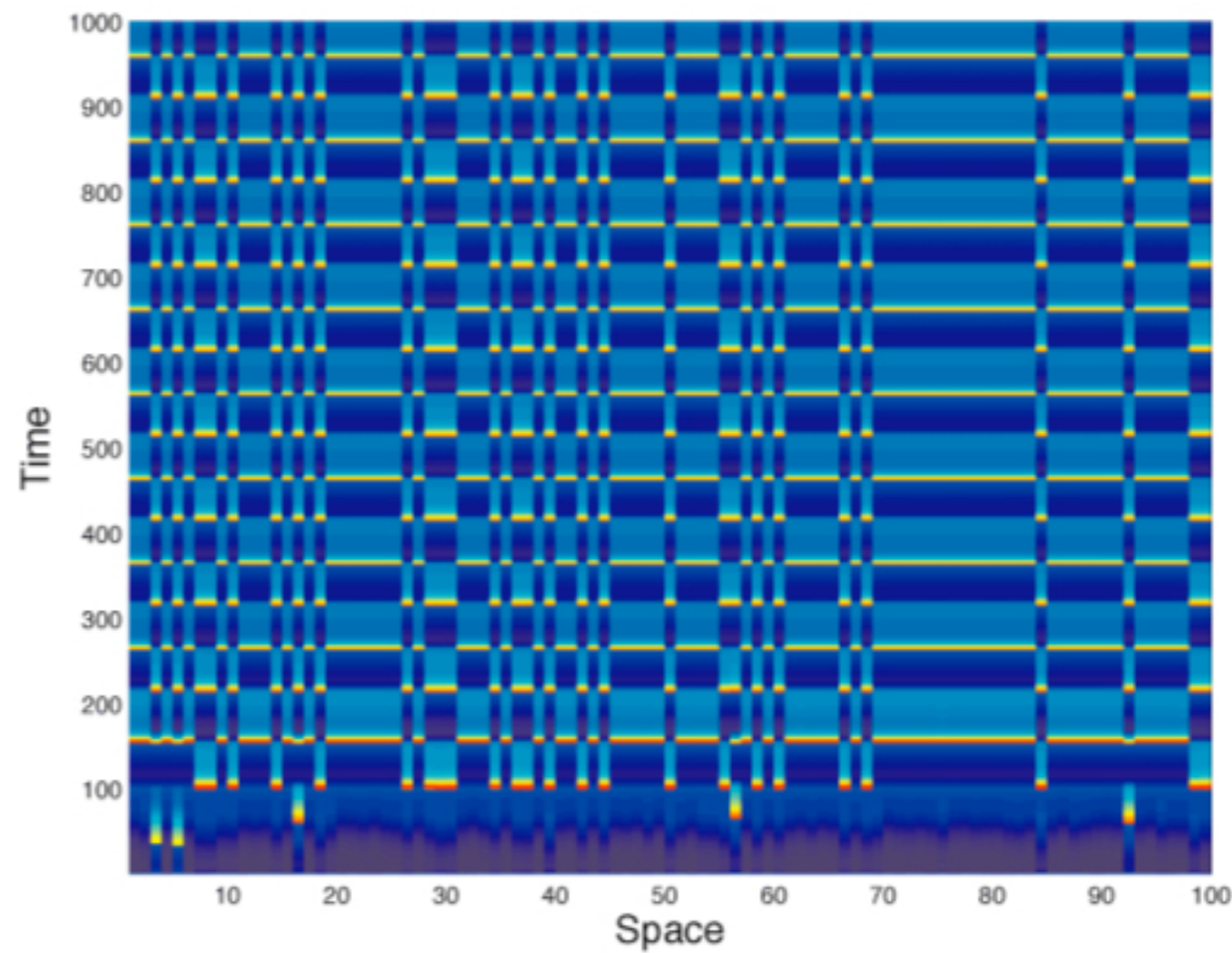


Global inhibitory TC Network

Synchrony for slow synapses

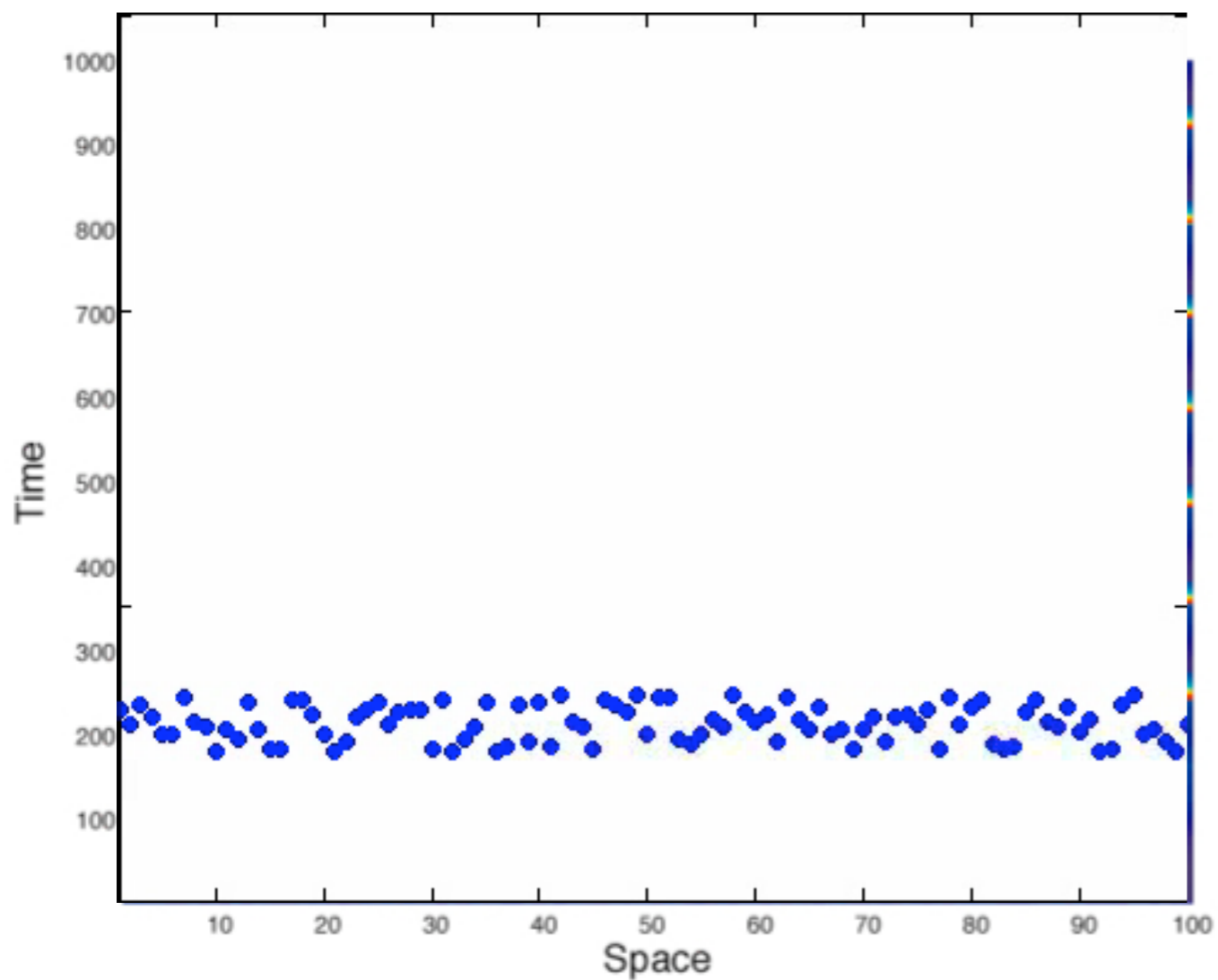


Clustering for faster synapses

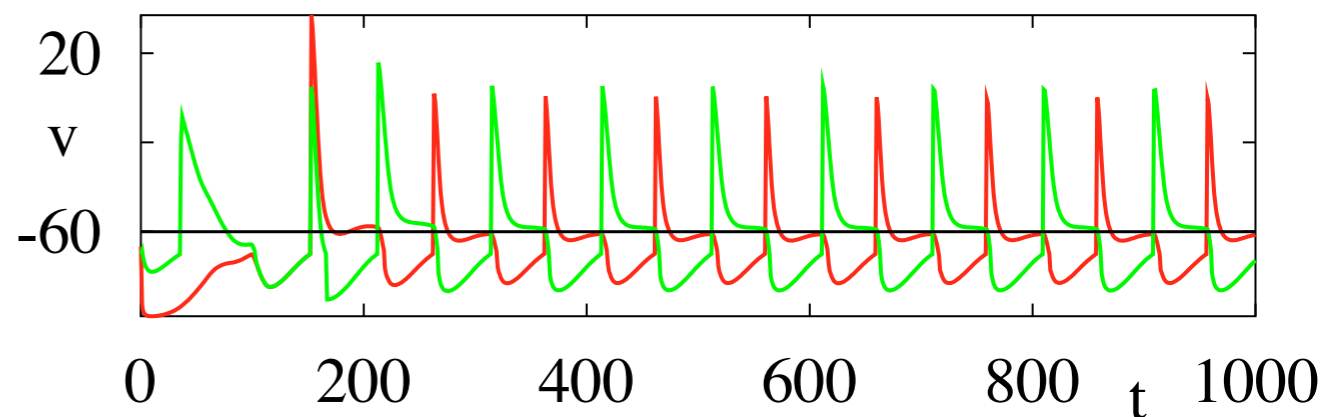
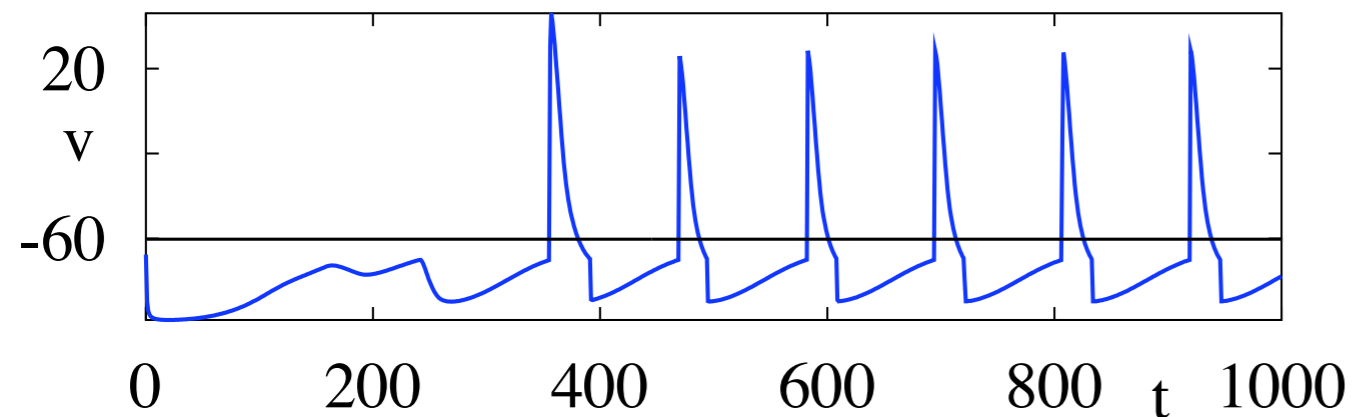
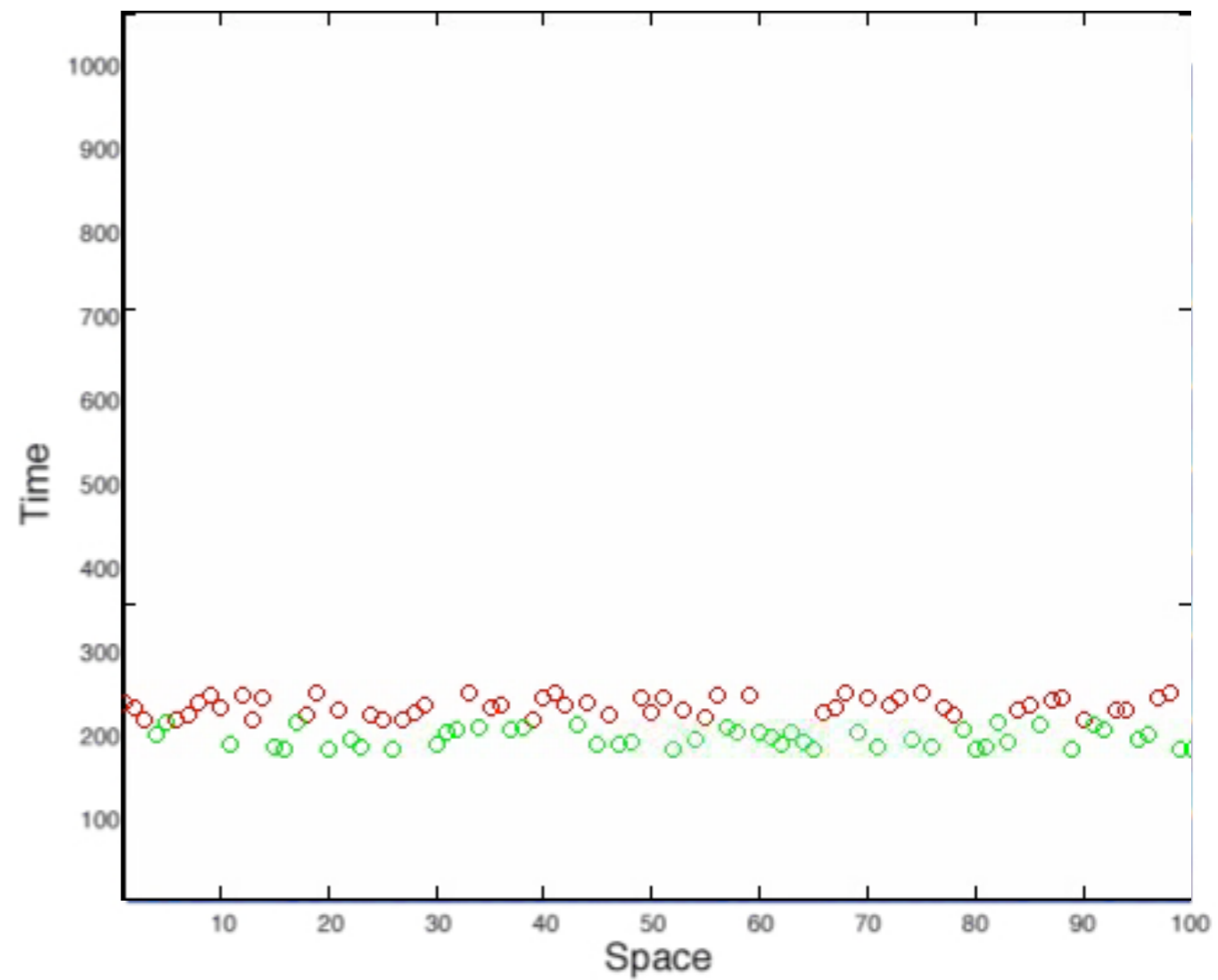


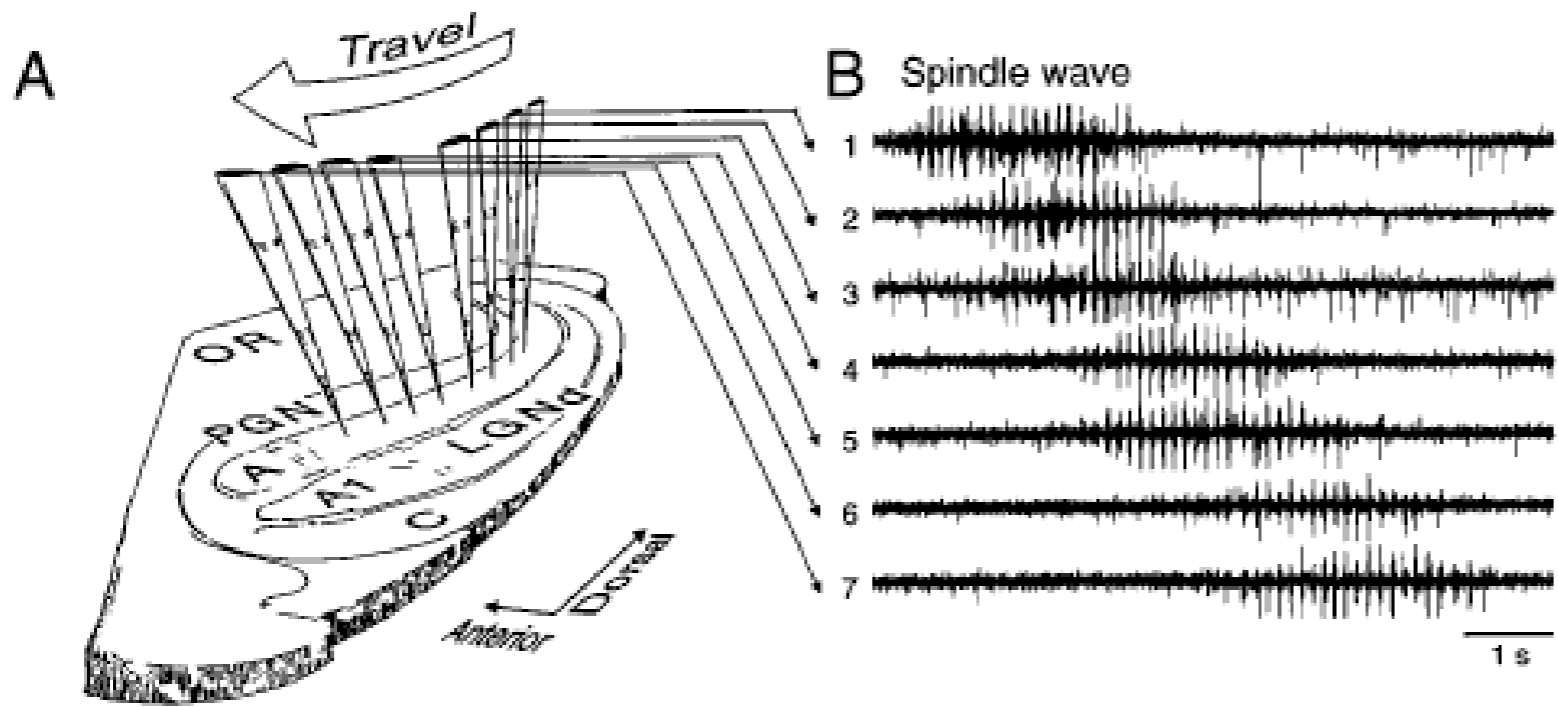
Global inhibitory TC Network

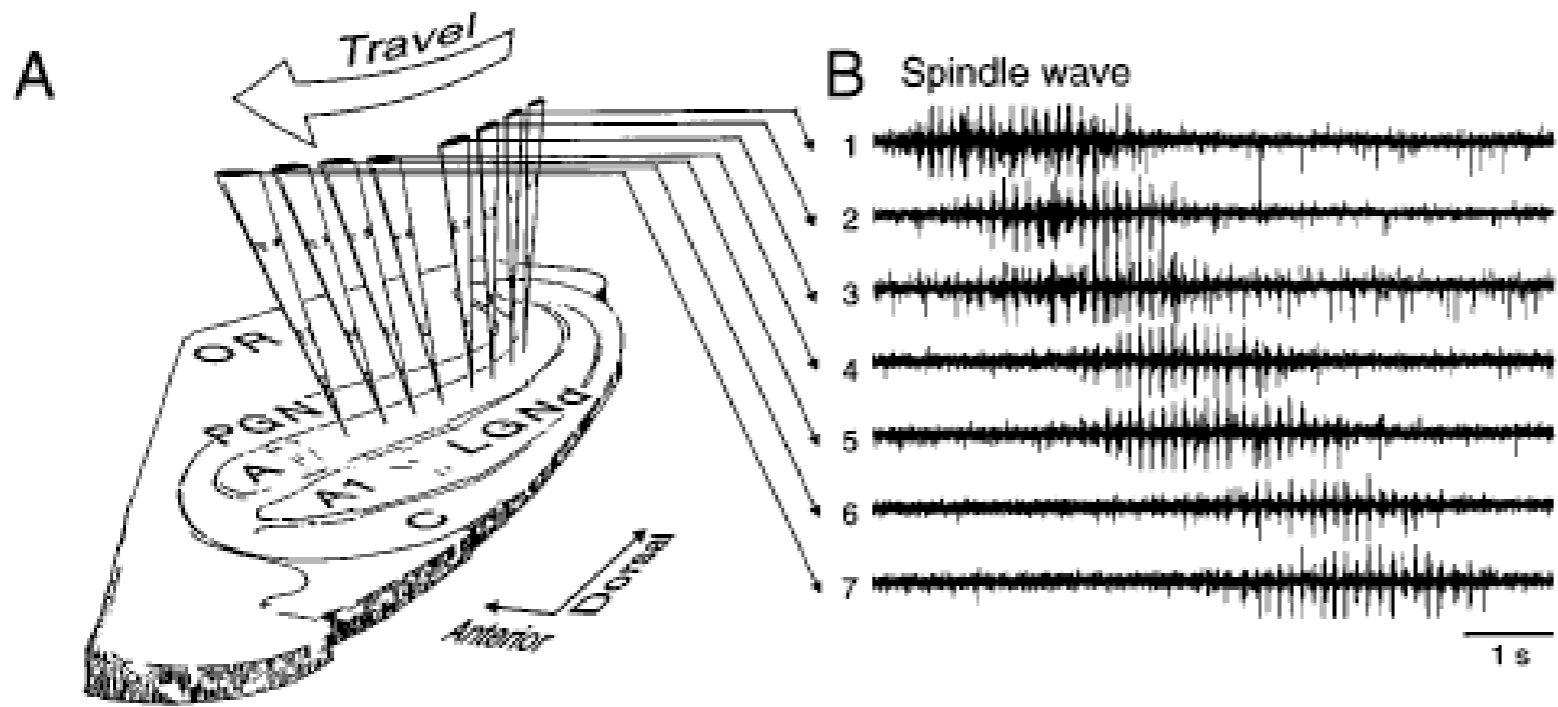
Synchrony for slow synapses



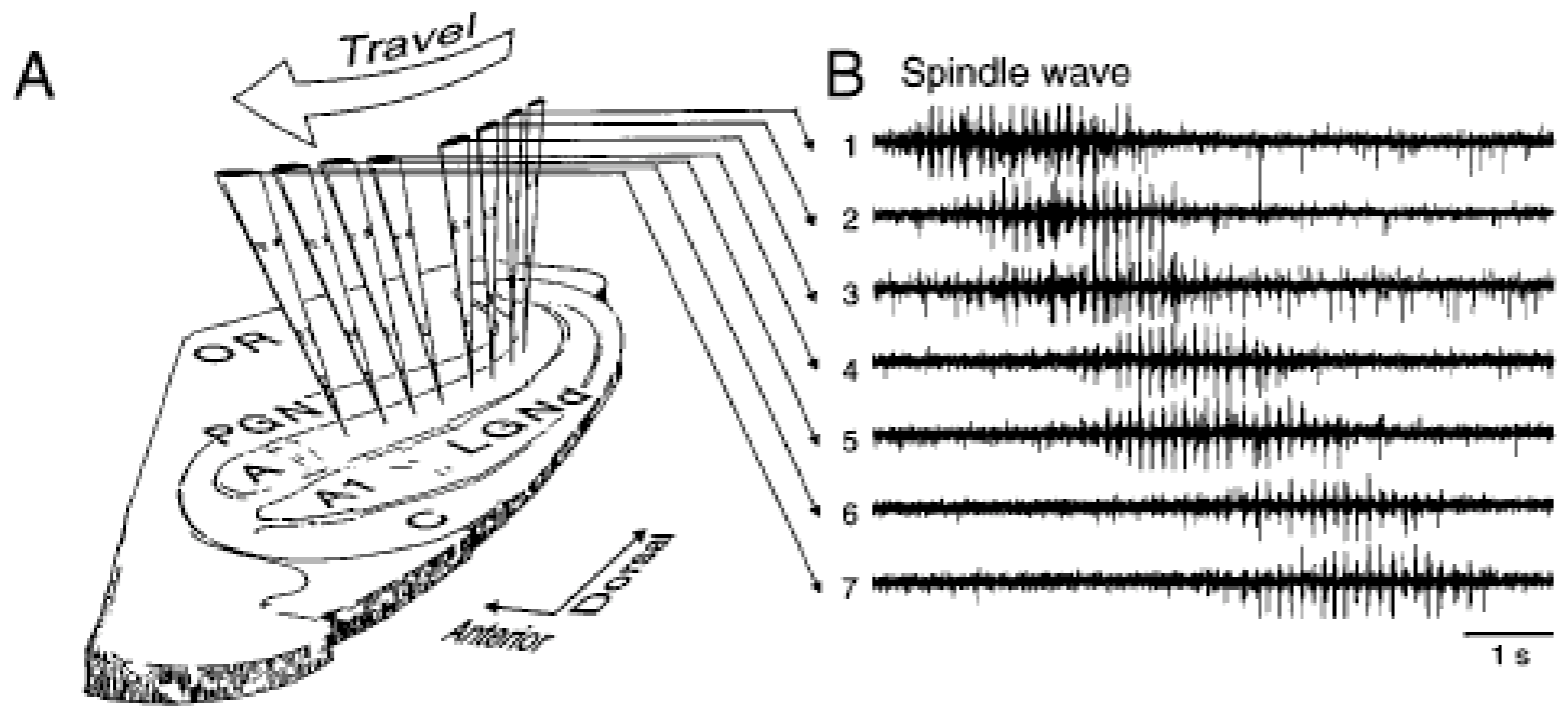
Clustering for faster synapses



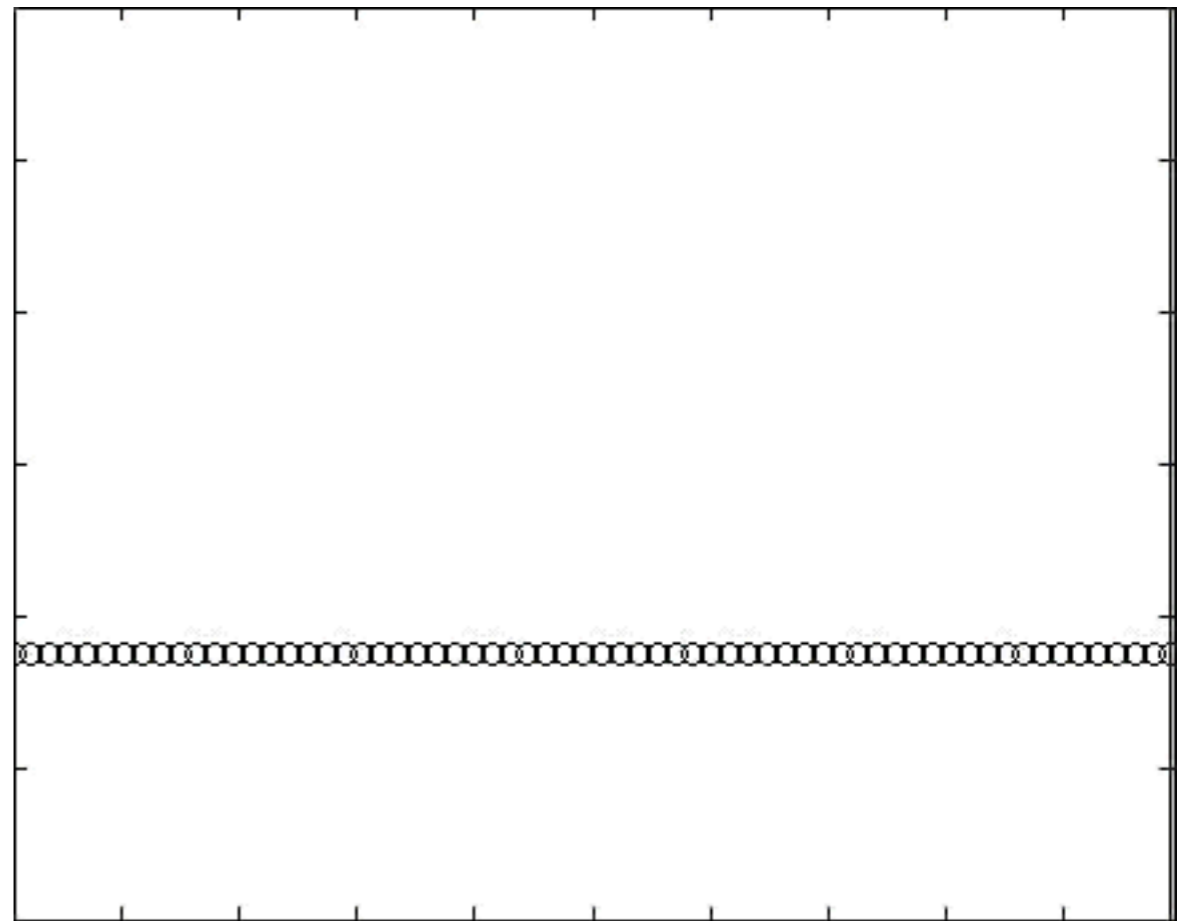


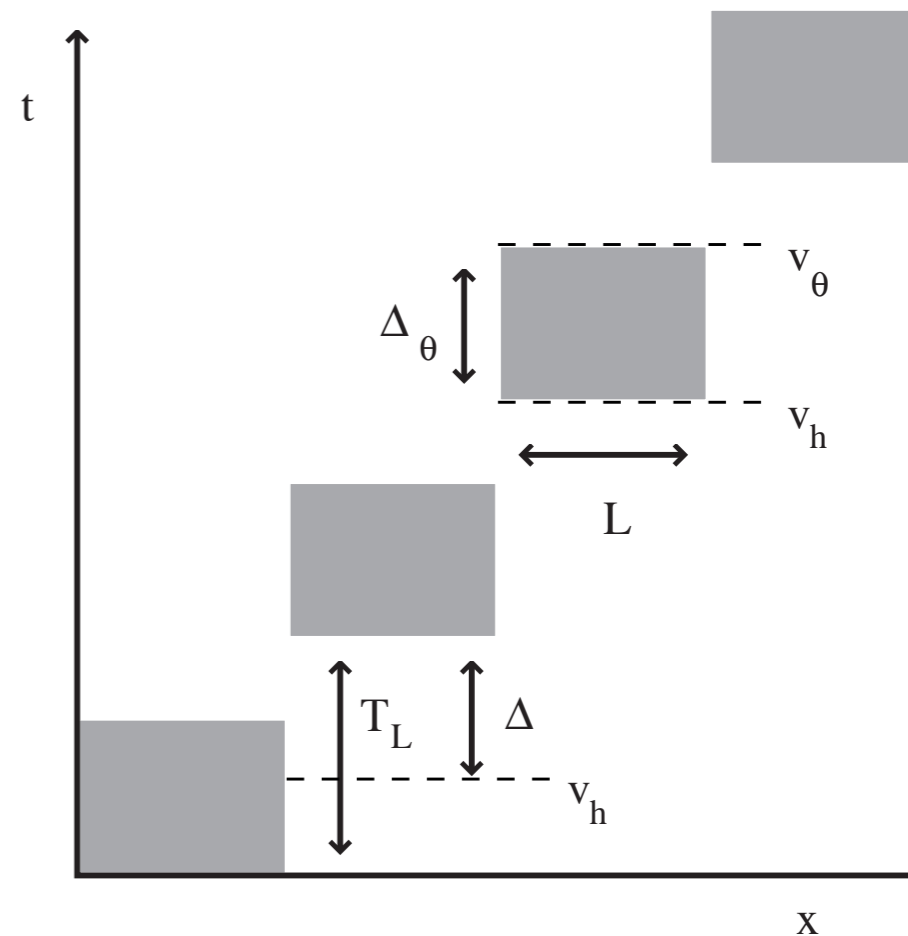
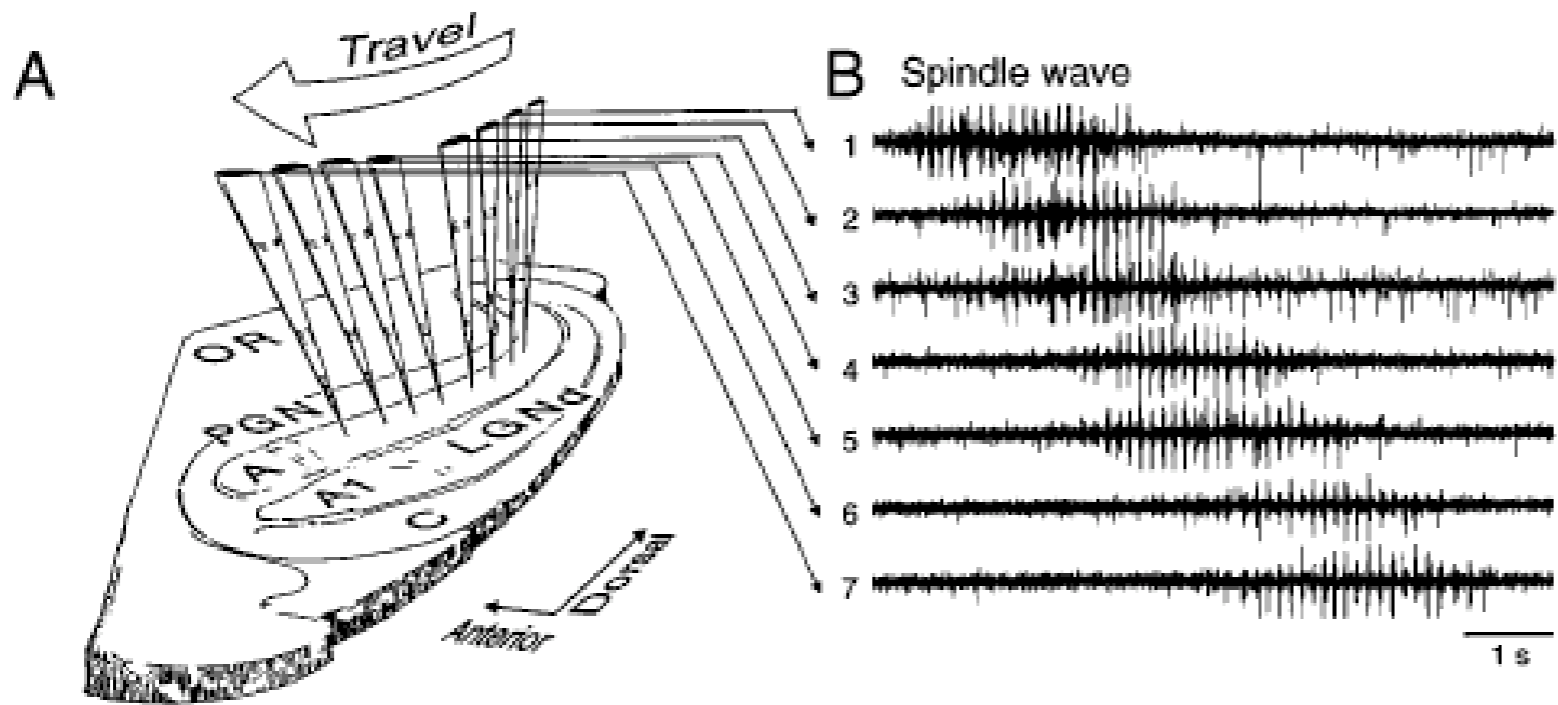


short
range

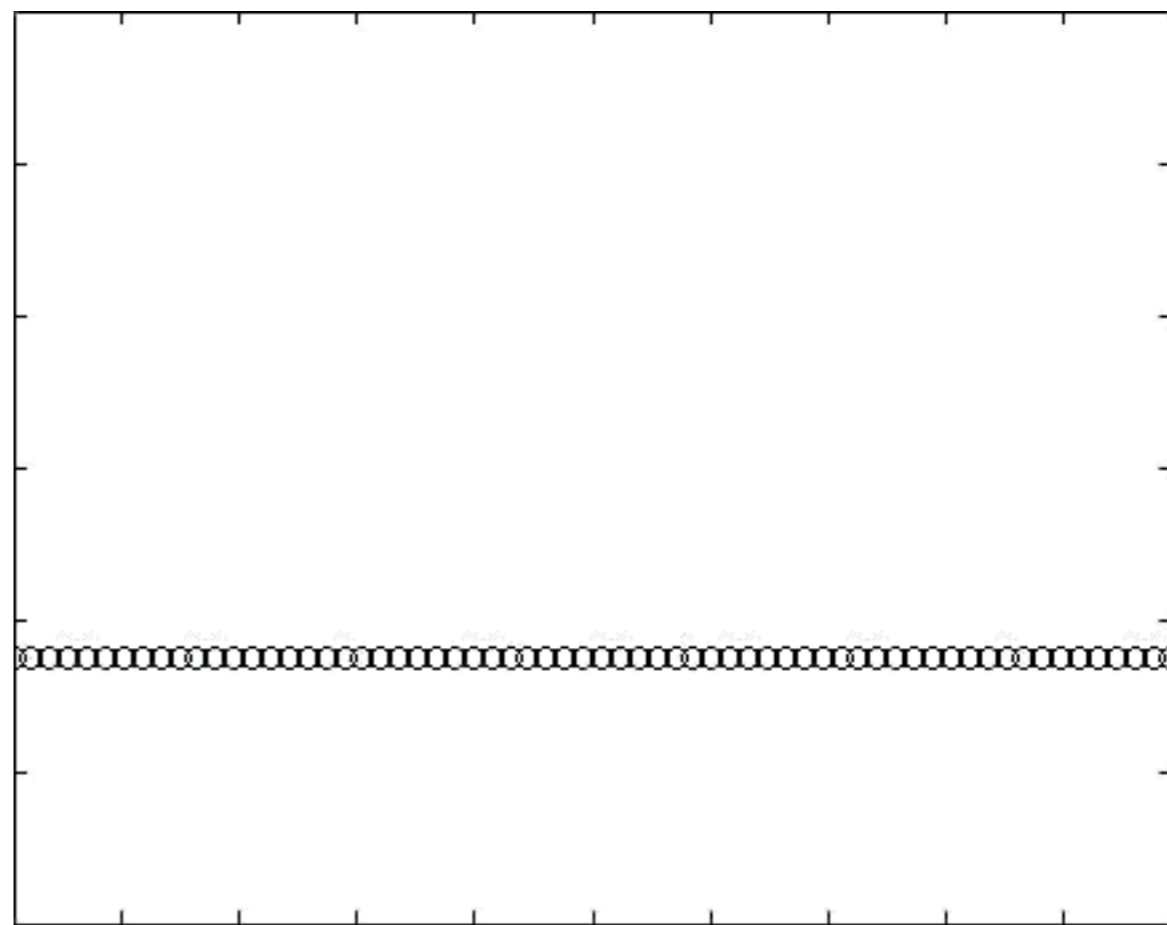
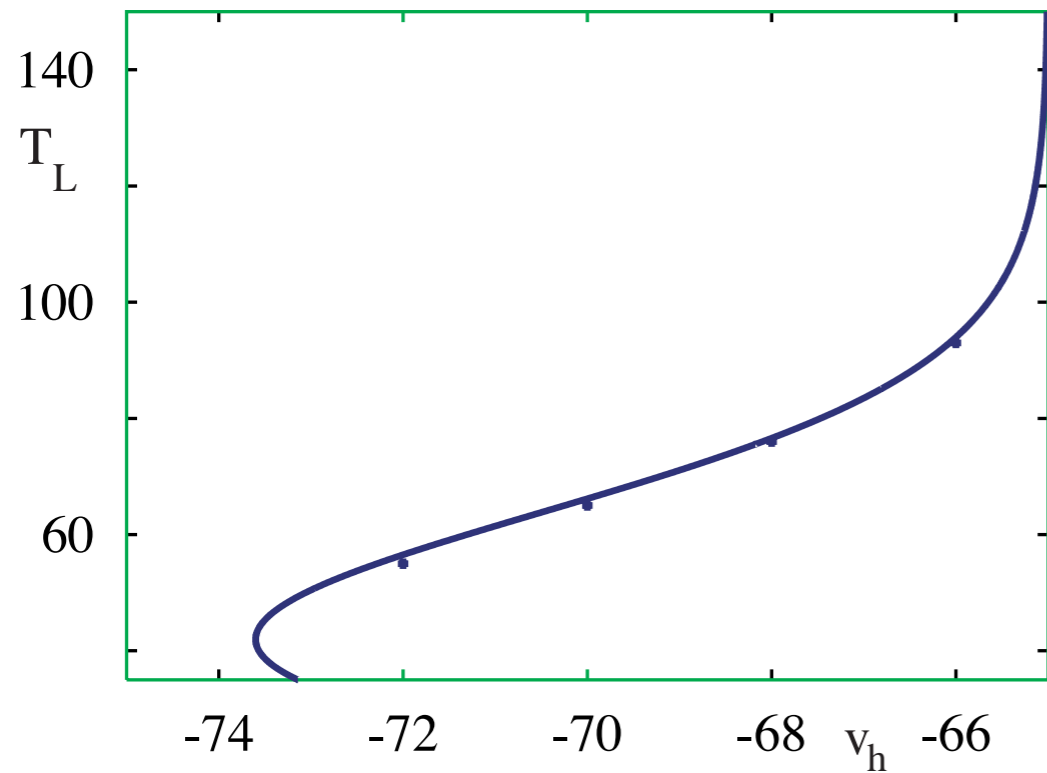


short
range

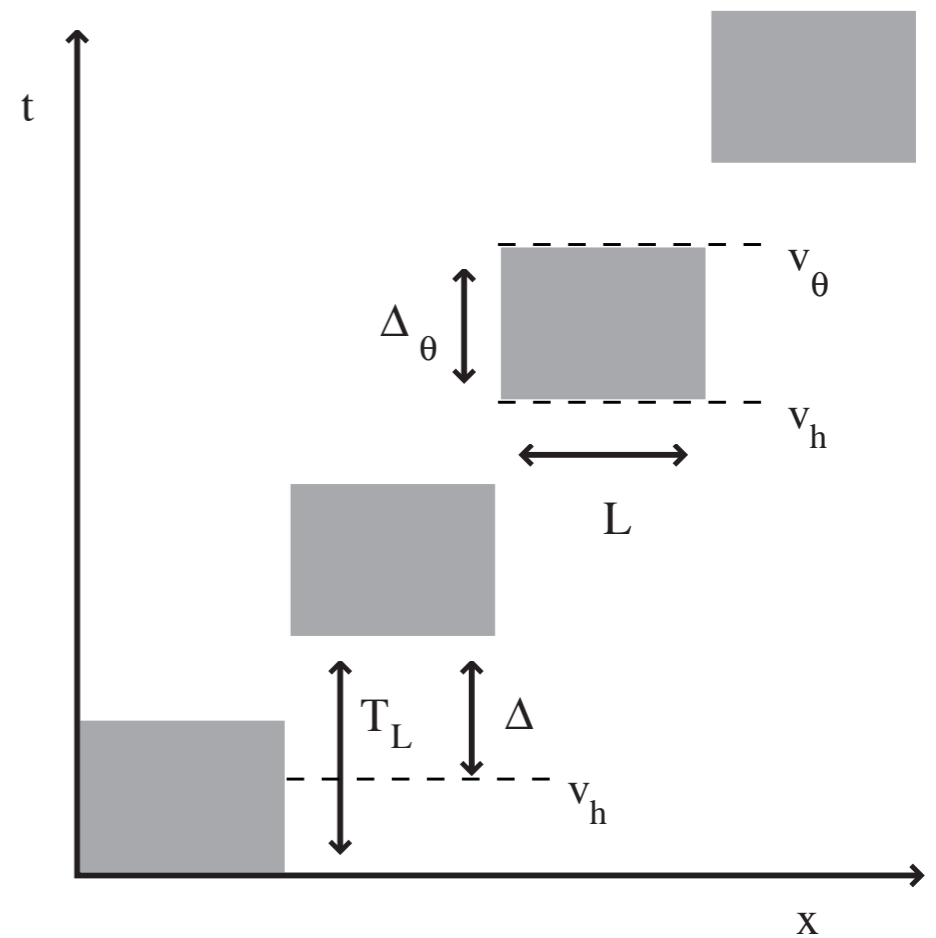
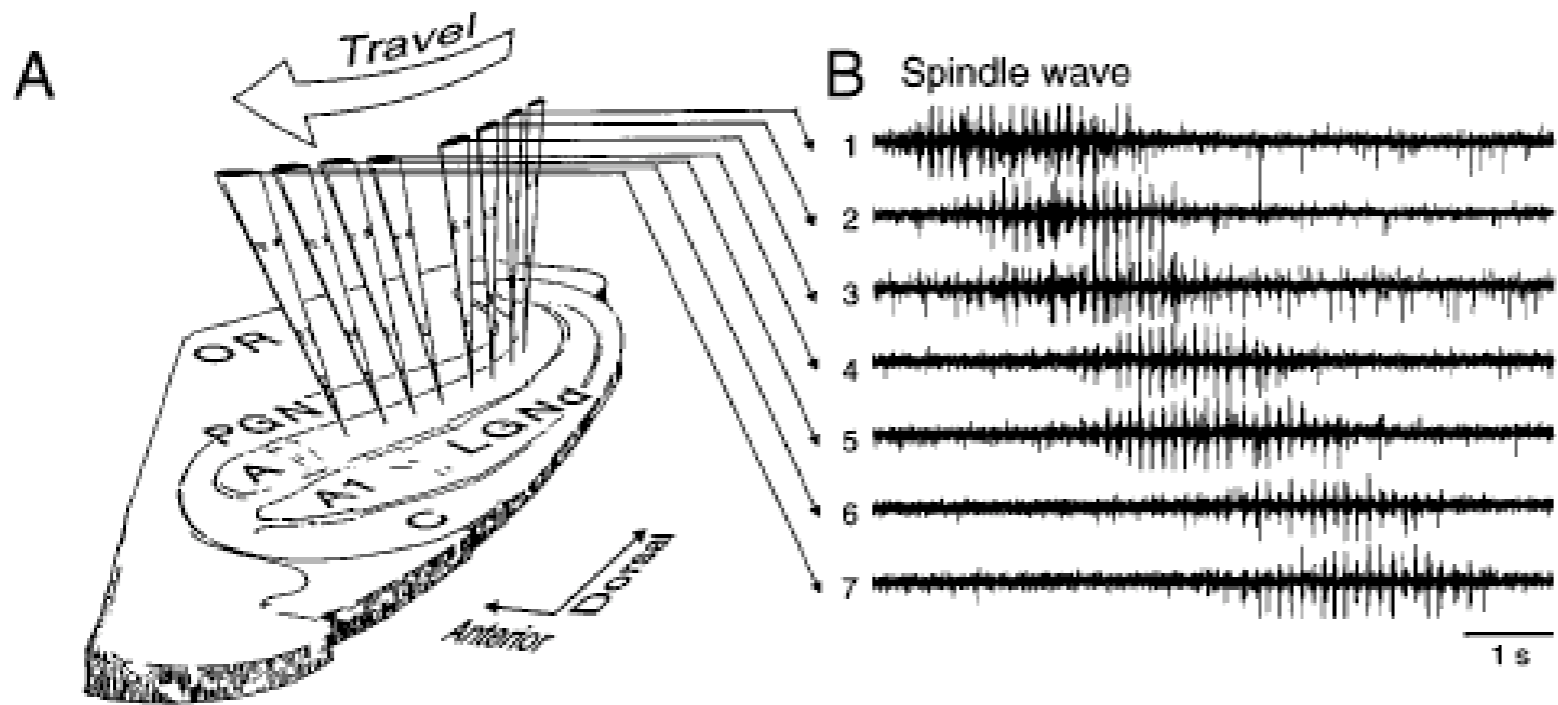




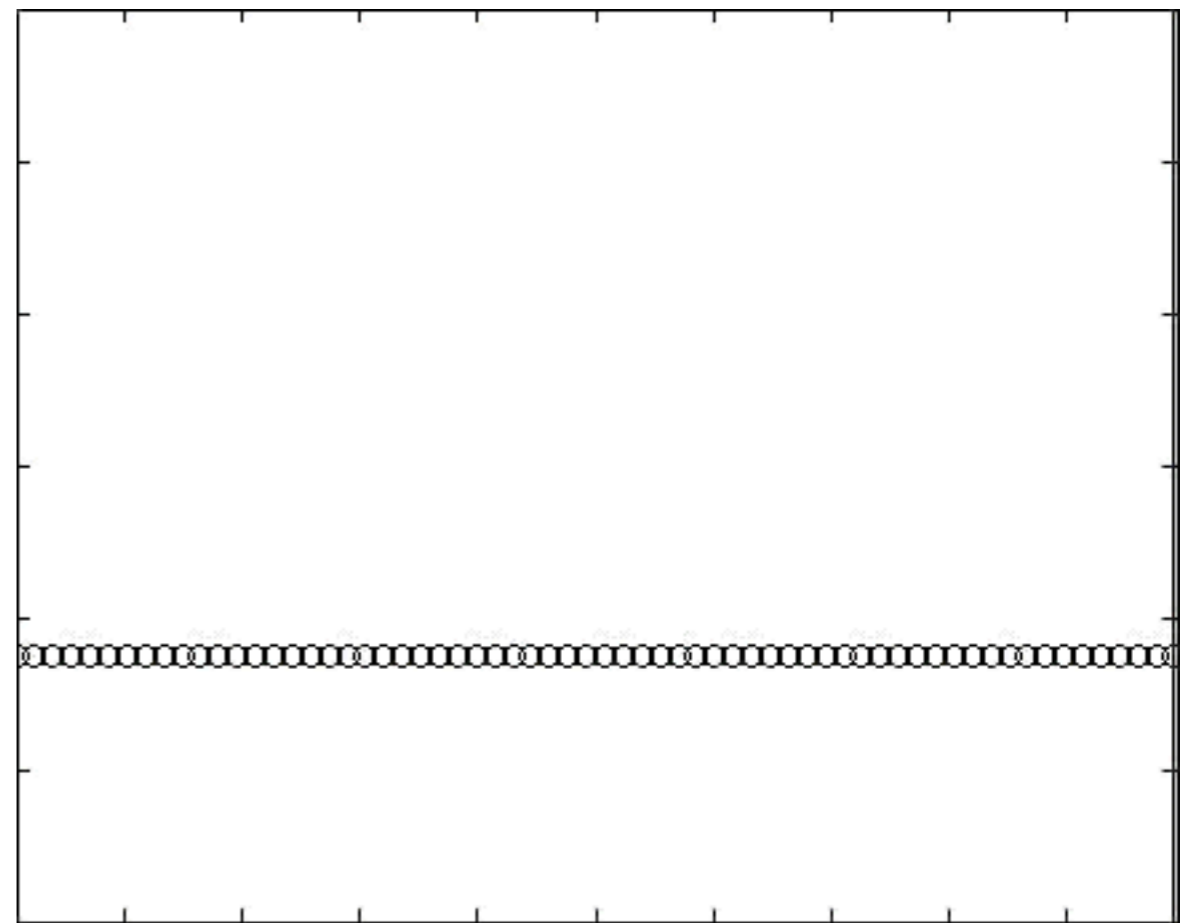
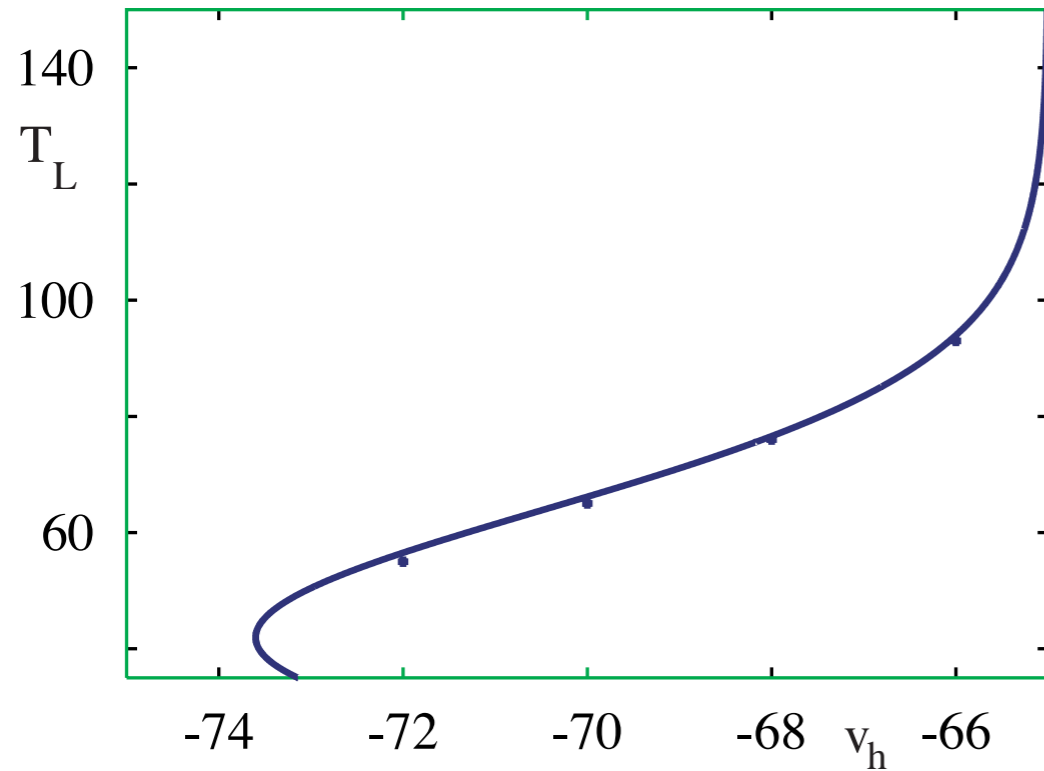
short
range



S Coombes, Dynamics of synaptically coupled integrate-and-fire-or-burst neurons, *Physical Review E* **67**, 041910 (2003)



short
range



S Coombes, Dynamics of synaptically coupled integrate-and-fire-or-burst neurons, *Physical Review E* **67**, 041910 (2003)

Slow lurching waves in a purely inhibitory network

In collaboration with

Nikola Venkov
(Notts)



Gabriel Lord
(Heriot-Watt)



Liejune Shiau
(Houston)



David Liley
(Melbourne)

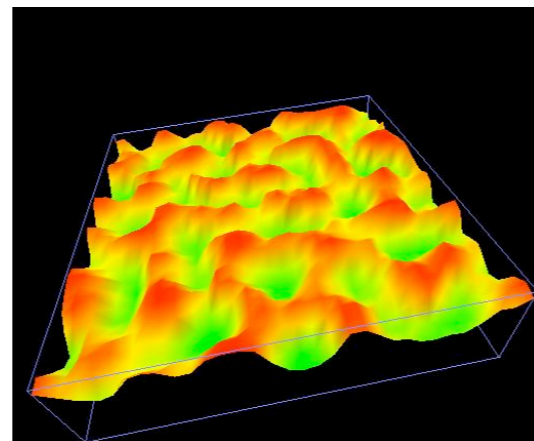


Engineering and Physical Sciences
Research Council

Markus Owen (Notts)



Ingo Bojak (Nijmegen)



Carlo Laing (Massey, NZ)

